A fundamental evaluation of the electrical properties and function of cardiac transverse tubules

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CONTENTS

1. Introduction – 3
2. T-tubules – the basics – 5
  2.1 T-tubules in cardiac excitability – 5
  2.2 Structure – 7
  2.3 Regulation and signaling – 8
3. Electrical properties of T-tubules – 8
  3.1 Passive properties – 8
  3.2 Delay of T-tubular membrane depolarization – 12
  3.3 Influence of T-tubules on macroscopic conduction velocity – 13
3.4 Action potentials in T-tubules? – 17
  3.4.1 Calcium-handling proteins – 19
  3.4.2 Sodium-handling proteins – 20
  3.4.3 Potassium-handling proteins – 22
  3.4.4 Other ion-handling proteins – 23
4. T-tubules in atrial cardiomyocytes – 25
5. Discussion – 25

ABSTRACT

This work discusses active and passive electrical properties of transverse (T-)tubules in ventricular cardiomyocytes to understand the physiological roles of T-tubules. T-tubules are invaginations of the lateral membrane that provide a large surface for calcium-handling proteins to facilitate sarcomere shortening. Higher heart rates correlate with higher T-tubular densities in mammalian ventricular cardiomyocytes. We assess ion dynamics in T-tubules and the effects of sodium current in T-tubules on the extracellular potential, which leads to a partial reduction of the sodium current in deep segments of a T-tubule. We moreover reflect on the impact of T-tubules on macroscopic conduction velocity, integrating fundamental principles of action potential propagation and conduction. We also theoretically assess how the conduction velocity is affected by different T-tubular sodium current densities. Lastly, we critically assess literature on ion channel expression to determine whether action potentials can be initiated in T-tubules.

KEYWORDS

Myocytes, cardiac; transverse tubules; cardiac conduction; action potential; voltage-gated sodium channels
1. Introduction

Transverse tubules are deep membrane invaginations of the lateral sarcolemma found in all striated muscle cells. In cardiomyocytes from mammalian cardiac ventricles, T-tubules form a particularly complex network throughout the cell [3, 4] (see Fig. 1), and are rich in ion channels and many regulatory proteins (reviewed in [5-7]). In atrial cells, the T-tubular network is dense, half-dense, or absent [8]. From the classic dictum “form follows function”, it is easy to formulate hypotheses about the function of T-tubules without becoming teleologic. Given that the T-tubular network protrudes deep into the cell, the increase in sarcolemmal surface functions as a platform for signaling and calcium handling [5]. Through T-tubules, signals from extra- to intracellular domains – both electrical and biochemical – can be conducted relatively fast. Indeed, T-tubules are the stage for many different processes (reviewed in [6]).

This work is dedicated to the question what the functional roles of T-tubules are. We discuss the depolarization delay of T-tubular membrane, and the effects of sodium current on membrane potential and electrical potentials in T-tubules. We also discuss principles of physiology in the context of T-tubules, and experimental data in the context of these physiological principles. Since the remodeling of T-tubules is observed in many cardiac diseases [9, 10], it is crucial to understand the roles of T-tubules in healthy cells before we can address the consequences of T-tubular remodeling.

We will only concisely discuss T-tubular structure (Section 2.2), signaling (Section 2.3), and remodeling in disease (Section 3.6), since excellent reviews have already been published on these topics [5-7, 9-11]. Rather, our work will focus on the electrical properties of T-tubules. Firstly, we will discuss the passive properties of T-tubules (Section 3.1), that is, the process by which initial depolarization of the cell membrane propagates into...
T-tubules without considering voltage-gated ion channels. This includes the discussion of relevant aspects of cable theory [12-16]. Special attention will go out to the capacitive current, which charges the membrane during cell excitation, whereby a larger membrane surface requires more time to charge. To this end, we discuss the delay of T-tubular membrane depolarization. We then assess extracellular potential along the T-tubules and its effect on sodium current amplitude (Section 3.2).

Section 3.3 is dedicated to conduction, since conduction velocity is an important parameter in arrhythmogenesis [17] and is affected by the presence of T-tubular membrane and ion channels. Over the course of the upstroke of an action potential, we describe the capacitive current, transmembrane current, and membrane potential changes for a cell. We will then assess theoretically how conduction velocity may differ between tissues with three different densities of T-tubular voltage-gated sodium channels: (1) no or few channels in T-tubules; (2) same channel density in T-tubules as in surface membrane; and (3) higher channel density in T-tubules than in surface sarcolemma.

In Section 3.4, we critically review published data on the expression of voltage-gated ion channels and ion pumps, because the data quality is at times unsatisfactory, yet often cited without critical connotations. This will elucidate whether T-tubules contain the necessary collection of ion handling proteins to generate an action potential. Although Section 3.1 and 3.2 will illustrate that theoretically passive conduction suffices to depolarize the T-tubular membrane quickly enough, action potentials may still be evoked.

Ion dynamics in the T-tubular lumen will be concisely discussed in Section 3.5. The small and tortuous T-tubular lumen restricts diffusion and affects driving force of voltage-gated ion channels [18, 19]. Lastly, we will discuss T-tubule heterogeneity in atrial cardiomyocytes in Section 4. Atrial cells contract as often as ventricular cells, yet have less T-tubules than ventricular cells or none at all [20].

Taken together, this work will thoroughly assess the electrical properties and roles of T-tubules and a better fundamental understanding of T-tubules will contribute to a better understanding of the implications of T-tubular remodeling in disease.
# 2. T-tubules – the basics

## 2.1 T-tubules in cardiac excitability

The benchmark role of T-tubules is to facilitate excitation-contraction coupling [7], the process that couples the action potential with muscle contraction. The cardiac action potential is triggered when gap junctions at the intercalated disc, where two cardiomyocytes are connected [21], conduct a depolarizing current from an excited to a resting cell. Once the membrane potential reaches the activation threshold of voltage-gated sodium channels, these channels open quickly and conduct a large inward sodium current, further depolarizing the membrane (see Fig. 2). Sodium channels at the intercalated disc may also be activated by ephaptic interactions [22-24], a process relying on interactions between sodium current and extracellular potential. Besides the intercalated disc, sodium channels are also expressed in the lateral membrane [25]. Whether sodium channels are present in the T-tubules will be discussed in a later chapter. As soon as the membrane has depolarized further, the L-type voltage-gated calcium channels, which are mainly located in the T-tubules [26], also open and conduct an inward calcium current. Voltage-gated calcium channels are close to (~ 15 nm) ryanodine receptors (isoform RyR2 in the heart) in the sarcoplasmic reticulum (SR) [7], together forming dyads. Calcium influx through T-tubular channels therefore quickly leads to calcium efflux from the SR conducted by RyR2, a process called calcium-induced calcium release [7]. The increased intracellular calcium concentration then induces sarcomere shortening [7]. By bringing voltage-gated calcium channels close to sarcomeres at the periphery as well as in the core of the cell, T-tubules provide a large stage for calcium handling proteins, promoting synchronous contraction of the whole depth of the cell [27]. Later during the action potential, voltage-gated potassium channels also open. The opposing potassium and calcium currents create the plateau of the action potential, where the membrane stays depolarized for a while, after which calcium channels close and the potassium current repolarizes the sarcolemma. Then, the sodium-potassium ATPase and the sodium-calcium exchanger among others restore ion concentrations at both sides of the membrane, while SERCA (sarcoplasmic reticulum Ca\(^{2+}\) ATPase) takes up calcium from the myoplasm back into the SR.

T-tubules also serve as hubs for signaling molecules. Catecholamines for instance bind to β-adrenergic receptors in the T-tubules, modifying cardiac contraction and relaxation [6].
As for membrane area, the accommodation of more calcium is efficient excitation of cardiac contraction coupling must occur, and the more membrane surface is required to accommodate more calcium-handling structures such as dyads.

As for membrane area, the measurements differ depending on the applied method, even within one species. For rats, T-tubular membrane area estimations range from 21 to 64% of total sarcolemma [30]. Causes and implications of these differences will be discussed extensively in Section 3.3.
On a nanoscale, T-tubules connect with the flat ends of SR, known as terminal cisternae, where dyads are located. Connections of T-tubules with the nuclear envelope, endoplasmic reticulum (ER), and mitochondria have also been described [5, 31]. Nuclear envelope-T-tubule connections are likely to facilitate a process called excitation-translation coupling [32]. Moreover, the T-tubular membrane forms microfolds. Their folding is mediated by the BAR-domain protein Bin1 [33].

T-tubules also contain caveolae, 50-100 nm-wide flask-shaped invaginations of the membrane. Their coating of caveolin and cavin sets them apart from other lipid rafts. The density of caveolae seems to be equally high in the T-tubular and other sarcolemma [34, 35].

Considering that T-tubules associate with many proteins and have tortuous membrane folds, T-tubules are complex structures that still can be dynamic under certain circumstances, for instance in disease [9, 36].

2.3 Regulation and signaling

Many factors regulate T-tubules biochemically or biophysically, a few of which will be mentioned in this section (reviewed in [5, 6, 27]). Caveolae for instance compartmentalize specific signaling proteins, such as the β2-adrenergic receptor (β2-AR), which forms signaling clusters with the G protein Gs and Gi. Gs activates the adenylyl cyclases AC5 and AC6 and PKA. AC5 and AC6 in turn produce cAMP, which has a wide range of downstream effects. Targets of PKA include calcium-handling proteins such as the voltage-gated calcium channel and RyR2, and the sarcomeric proteins tropinin I and C. β1-AR on the other hand is not confined to caveolae and couples exclusively to the stimulatory Gs pathway (reviewed in [6]). The L-type voltage-gated calcium channel is expressed both inside and outside of caveolae. The caveolar channels appear to be cAMP-activated and affect hypertrophy instead of excitation-contraction coupling [37]. Other noteworthy regulators of T-tubules are: (1) Tcap, which regulates the formation of T-tubules under influence of load and stretch of the cell [38]; (2) junctophilin-2, which stabilizes dyads [39, 40]; and (3) Bin1, which folds the T-tubular membrane and recruits calcium channels [41, 42].

3. Electrical properties of T-tubules

3.1 Passive properties

In this section, we will address the electrical properties of T-tubules, and how they affect cardiomyocyte excitability.

During the propagation of the depolarizing wavefront, the voltage difference between neighboring cells is the driving force for intercellular current. This current is conducted by gap junctions and depolarizes the
membrane of the neighboring cell. Once a membrane depolarizes at the surface of a myocyte, the change in voltage induces passive electrotonic responses of the T-tubular membrane for which ion channel openings are not required. To assess whether passive conduction suffices to depolarize the entire T-tubular system, cable theory comes into play [43, 44]. If we consider a passive T-tubular membrane, when the cell is excited by an action potential, the T-tubular transmembrane voltage drops due to current leaking across the membrane resistance \((r_m)\) and flowing outwards down the T-tubules through their longitudinal resistance \((r_{TT})\). These effects limit the extent of electrotonic responses. To assess the decay of action potential propagation, the length constant for guinea pig T-tubules was theoretically estimated at \(\lambda = 240 \mu m\) [10, 45]. Kong et al. also calculated the length constant based on optical measurements in mouse cardiomyocytes at \(\lambda \approx 240 \mu m\) [46].

To assess passive electrical properties experimentally, Scardigli et al. [47] were inspired by an experimental geology study, in which conductivity of a porous rock was determined to assess the diffusive properties. In isolated cardiomyocytes, the apparent diffusion coefficient \(D'\) of fluorescent dextran was determined by fluorescence recovery after photobleaching (FRAP) microscopy. The cardiomyocytes were modeled as porous cylinders and Fick’s second law of diffusion was applied to calculate \(D'\) from the time to recovery after photobleaching and the cellular radius \((D' \approx 1.4 \mu m^2\cdot s^{-1})\). Based on the analogies between Fick’s first law of diffusion and the electrical current density law, the electrical conductivity of T-tubules was determined at \(K' = (5.3 \pm 0.5) \times 10^{-4} \Omega^{-1}\cdot cm^{-1}\). Lastly, the length constant was calculated from the diffusion rate and the previously measured T-tubular surface area per unit cell volume: \(\lambda = 290 \pm 90 \mu m\). This is very similar to the aforementioned values.

Scardigli et al. however did not properly consider that the analogy between conductivity and diffusion depends on the detailed porous structure, which in geology is described by the so-called “formation factor”. This factor describes the ratio of the resistivity of a porous rock filled with water to the resistivity of water and must be empirically measured or defined for different rock types. For cardiac myocytes, the actual factor is unknown. Scardigli et al.’s estimate of the diffusion coefficient may therefore be subject to uncertainty and thus differ from values reported by for instance Kong et al. \((D' \approx 23 \mu m^2\cdot s^{-1})\) [29]. The obtained length constant however resembled that found by Kong et al. because the formation factor offsets diffusion [29].

For a rat T-tubule with a mean length of 6.84 \(\mu m\) [1], Scardigli et al. calculated a voltage drop from the surface sarcolemma to the cell core of \(\sim -4 mV\) assuming an infinite cable, and thus an exponential decay of potential along the T-tubule. The voltage drop in an infinite cable can be described by [48, 49]

\[
\Delta V(x) = \Delta V_0 e^{-\frac{x}{\lambda}}, \quad (\text{Eq. 1})
\]
where $\Delta V_0 = 100$ mV and $\lambda$ is the length constant, which can be described by the following cable-theory-derived equation [49]

$$\lambda = \sqrt{\frac{r}{2g_m\rho_c}} \quad \text{(Eq. 2)}$$

where $r$ is the T-tubular lumen radius, $g_m$ the conductance of resting membrane per unit area, and $\rho_c$ the resistivity of extracellular space. A sealed-end cable however approaches a T-tubule more closely than an infinite cable does. Voltage drop in a sealed-end cable at depth $x$ can be described by [48, 49]

$$\Delta V(x) = \Delta V_0 \frac{\cosh \left( \frac{L-x}{L/\lambda} \right)}{\cosh \left( \frac{L}{L/\lambda} \right)} \quad \text{(Eq. 3)}$$

where $L$ is the T-tubule length [48]. Eq. 3 decreases less steeply with $x$ than Eq. 1, and if $L \ll \lambda$, $\Delta V(x)$ hardly decays at all as described by Eq. 3. For the T-tubular properties described by Scardigli, Eq. 3 yields a much smaller voltage drop of -0.028 mV [50].

Uchida and Lopatin [51] recently showed that a length constant of $\sim 290$ μm is probably a significant overestimation, because the previously discussed publications [45, 47] did not take the dilations and constrictions of T-tubules into account. In addition, T-tubules form branching networks, and can also form loops. In neuroscience, much effort was invested in deriving simplifying formulas to describe the decay of the electrotonic potential in neuritic trees with changing neurite diameters (e.g., flaring, tapering) and branches [49, 52-55]. For example, Rall has shown that a branching dendritic tree can be collapsed to a cylinder model under certain stringent constraints such as defined relationships between the diameters of parent and daughter branches [54-56]. While certain neurons may exhibit parent and daughter branch diameters compatible with such constraints, these constraints certainly do not pertain to T-tubules. The morphological heterogeneity of T-tubules disqualifies the straightforward application of classic cable theory.

Uchida and Lopatin determined the diffusion coefficient $D'_{\text{TT}} \approx 4$ μm$^2$·s$^{-1}$ in murine cardiomyocytes. This is comparable to the findings of Scardigli et al. ($D'_{\text{TT}} \approx 1.4$ μm$^2$·s$^{-1}$ as mentioned previously) [47], but not to the findings from Kong et al. who calculate a much higher diffusion coefficient in murine T-tubules of $\sim 23$ μm$^2$·s$^{-1}$ for 1 kDa solutes [29]. The diffusion coefficient found by Kong et al. corresponds to a $\sim 12$ times reduction of free diffusion rate, which is comparable to the $\sim 8$ times reduction of calcium ion diffusion estimated based on electrophysiological experiments [57]. This finding from Kong et al. likely approximates the in vivo diffusion coefficient better since their model is not confounded by cell and illumination beam geometries, which might have affected the results from Scardigli and Uchida [47, 51]. Interestingly, the diffusion coefficient in rabbit T-tubules is only $\sim 5$ times slower than free diffusion rate.
compared to the ~12 times reduction in mice [29]. This difference can be explained by the structural differences between murine and rabbit T-tubules: murine T-tubules are more tortuous close to the surface than rabbit T-tubules, and the luminal “ground substance” identified by electron microscopy – potentially representing the glycocalyx – forms a stronger diffusion barrier in murine T-tubules [29].

Uchida and Lopatin furthermore illustrate the importance of T-tubular structure in determining the diffusion time [51]. The in vivo diffusion time constant of 3.9 s could only be approached in an in silico 3D cell model of murine T-tubules by introducing 20-nm-wide constrictions and 600-nm-wide dilations in the model (values represent luminal diameters; $\tau_{\text{dilations}} = 2.2$ s and $\tau_{\text{constrictions}} = 2.5$ s) [51]. Extrapolating this to electrical properties under standard conditions for specific membrane capacitance, extracellular and membrane resistivity, the time constant for cylindrical T-tubules would be $\tau_{\text{Vm}} = 10.1$ μs and for T-tubules with dilations or constrictions $\tau_{\text{Vm}} = 102$ μs [51]. The latter time constant approaches the experimental finding $\tau_{\text{Vm}} \approx 200$ μs in murine T-tubules [58]. Uchida and Lopatin however do not consider microfolds [33]. This membrane tortuosity could be corrected for by increasing the specific capacitance from 1 to 2 μF/cm$^2$, which would double the time constant $\tau_{\text{Vm}}$ from ~100 μs to ~200 μs, which is even closer to the experimental finding [58].

Constrictions and dilations considerably affect the length constant. For a 200-nm-wide T-tubule with short 20-nm-wide constrictions every 2 μm, $\lambda_{\text{TT}} \approx 68$ μm [51]. This is considerably smaller than the $\lambda_{\text{TT}} \approx 240$-290 μm reported before [45, 47]. To transition from the steady state to the non-steady state, the length constant needs to be corrected for the characteristic frequencies ($f$) of the action potential (during the upstroke, $f \approx 150$ Hz) [46], and $\lambda_{\text{AP}}$ will be $\sim 34$ μm [51]. Since the length constant also depends on membrane resistance $R_m$, the length constant will decrease further when voltage-channels are open.

This implies that, in an average murine T-tubule of 9 μm long [42], the voltage drop will be ~13 mV from the surface to the end of the T-tubule when assuming an infinite cable as described by Eq. 1. When assuming a sealed-end cable, the applicable equation is Eq. 3, which gives a voltage drop of only 0.87 mV. In any case, this voltage drop is small enough to allow voltage-gated calcium channels to open, assuming that the surface sarcolemma depolarizes to +20 mV and the threshold of calcium channels lies near -40 mV. The question whether the entire T-tubular membrane depolarizes quickly enough to open calcium channels once a cell is excited is discussed in Section 3.2. To conclude, we have to keep in mind that passive electrical parameters will differ significantly between species. Whereas rabbit T-tubules openings are relatively accessible and easy to follow, openings of murine T-tubules are relatively tortuous and filled with glycocalyx, which considerably limits diffusion [29].
3.2 Delay of T-tubular membrane depolarization

T-tubules greatly increase membrane surface, so charging the capacitance of the membrane of a tubulated cell will take longer. The time it takes to charge the membrane as a capacitor and depolarize T-tubules was assumed to be negligible [45, 59], but recent results indicate this is incorrect [50, 60]. Quantifying the capacitance charging time of T-tubules would give important insight into the delay between the opening of voltage-gated channels at the sarcolemma and those deep in the T-tubules. This section addresses this depolarization delay and the effect of voltage-gated sodium channels on tubular depolarization.

We developed a simple computational model of a murine T-tubule, represented as a 9-µm-long cylinder with a diameter of 98.5 nm and a specific capacitance of 2 µF/cm² to simulate microfolds (normal specific capacitance is considered 1 µF/cm²), subdivided in 100 segments. This model revealed a depolarization delay of 10 µs in the innermost segment in the absence of inward currents in the T-tubules [50]. This value is expected to greatly increase in the presence of T-tubular branches and frequent 20-nm-diameter constrictions. Conversely, sodium currents may facilitate T-tubular depolarization. Although Uchida and Lopatin determined that these structural changes increase the time constant of depolarization from ~ 10 to ~ 100 µs for this model (see Section 3.1) [51], they did not determine the delay of membrane depolarization of a deep T-tubular segment. Of note, linear cable theory assuming a passive linear membrane has limited applicability to a T-tubule given their non-linear current-voltage relationships. It is expected that the depolarization delay of a deep segment of a T-tubule with constrictions and dilatation is considerably greater than the 10-fold increase by which the time constant increased, because each dilatation poses a large capacitive load, in addition to the depolarization delay due to the higher resistance in each constriction. Taken together, the question remains open how much later voltage-gated calcium channels open deep inside the T-tubules when the surface sarcolemma is excited.

The effect of a voltage-gated sodium current on the depolarization of T-tubules and the effects of intratubular potentials on the sodium current have been simulated by Hatano et al. [60] in an elaborate 3D model with structurally simple T-tubules without branches, constrictions and dilations. The T-tubular sodium current density was set at 30 mS/µF according to experimental results from [45] (see Section 3.4.2 for a critical evaluation of these results). In the deepest segments of the modelled T-tubule (5 µm from the surface), the extracellular potential was slightly negative (-1 mV) because of the flow of current along the narrow tubule. Due to this negative potential, the sodium current was ~ 8% smaller in the deep segment than the segment at the cell surface [60]. A similar result was obtained with our simpler model [50] with a slightly smaller T-tubular sodium current density of 23 mS/µF. The mechanism underlying this sodium current reduction in the deep portions of the T-tubules resides in the notion that an ion current \( I_X \) carried by ion species X can be described by \( I_X = g_X \cdot P_O \cdot (V_m - E_X) \), where \( g_X \) is the maximal conductance of that
current, $P_o$ is the open probability of the corresponding channels, $V_m$ is membrane potential, and $E_X$ is the reversal potential of $I_x$. The term $V_m - E_X$ is called the driving force. For the case of the sodium current, the driving force becomes smaller (in absolute value) when the extracellular potential becomes negative, thus decreasing the current (in absolute value). A comparable mechanism was called “self-attenuation”, and has been proposed to occur in intercalated discs due to ephaptic interactions [22-24]. Thus, the sodium current is slightly self-attenuating deep inside the T-tubules. Additionally, the tubular lumen can be transiently depleted from sodium if the T-tubular sodium current is substantial. Sodium depletion decreases $E_{Na}$, the Nernst potential of sodium and thus the driving force, representing an additional mechanism reducing the sodium current. This was specifically addressed in a modeling study by Mori et al. who showed that sodium depletion in the intercalated disc contributes to self-attenuation [61]. In addition, a recent study by Greer-Short et al. indicates that sodium depletion in intercalated discs and the resulting decreased driving force for sodium ions are also important in shaping repolarization as they attenuate the late sodium current. Interestingly, this phenomenon can mask the phenotype of long-QT syndrome type 3 [62].

3.3 Influence of T-tubules on macroscopic conduction velocity

In this chapter, we will assess theoretically how macroscopic conduction velocity (i.e., conduction along the surface membrane and throughout cardiac tissue) is influenced by T-tubular ion channels and membrane. To do so, one has to take a closer look at the propagation of action potentials through cardiac tissue. Much insight has been gained from the ground-breaking work of Hodgkin and Huxley and others [12-16, 63].

Firstly, one has to understand the dynamics of current flow into, out of, and within a cell during the propagation of the depolarizing wavefront as detailed and illustrated in Fig. 2. The dynamics of various currents during the upstroke of the action potential are crucial determinants of conduction.

During the passage of a wavefront, electrotonic current enters the cell from upstream ($I_{in}$) and exits the cell downstream ($I_{out}$) (Fig. 2B and 2D, in black). According to Kirchhoff’s current law [64], the difference $I_{in} - I_{out}$ must be transferred to the membrane in the form of capacitive current ($I_{cap}$; Fig. 2B and 2D, in green) or flow through membrane ion channels and transporters ($I_{ion}$; Fig. 2B and 2D, in red).

The capacitive current $I_{cap}$ directly follows the upstroke velocity of the action potential (Fig. 2C), i.e., the rate of change of membrane potential $V$, expressed as $\partial V/\partial t$. Assuming that cell capacitance $C$ is constant, $I_{cap}$ is indeed defined as:
\[ I_{\text{cap}} = C \cdot \frac{\partial v}{\partial t} \]  
(Eq. 4)

Thus,

\[ I_{\text{in}} - I_{\text{out}} = C \cdot \frac{\partial v}{\partial t} + i_{\text{ion}} \]  
(Eq. 5)

Assuming that the cell is much smaller than the length constant of the tissue and the spatial extent of the upstroke (which is the case for well-coupled cardiac tissue), then the tissue can be considered as a continuous syncytium \[17, 64\]. In one dimension, Eqs. 4 and 5 can be written in a spatially continuous form as:

\[ \sigma \frac{\partial^2 v}{\partial x^2} = c \cdot \frac{\partial v}{\partial t} + i_{\text{ion}} \]  
(Eq. 6)

where \( x \) is position, \( \sigma \) is the lumped conductivity (normalized by surface-to-volume ratio) of the gap junctions, the intracellular space, and the extracellular space; \( c \) is the membrane capacitance per unit area; and \( i_{\text{ion}} \) is the total ion current per unit area. This equation is known as the cable equation \[12, 13\], or in the cardiac community as the monodomain equation, which is valid for cardiac tissue when assuming that the anisotropy ratios of intra- and extracellular resistivities are equal \[64\].

When conductivity (\( \sigma \)) changes, conduction velocity (\( \theta \)) follows a square-root relationship \[13, 49, 64\]:

\[ \theta \propto \sigma^{1/2} \]  
(Eq. 7)

Now, we will consider three possible scenarios to assess the influence of T-tubular membrane and channels on conduction (Fig. 3). In the first (Fig. 3A), T-tubules contain a low density of channels, especially of voltage-gated sodium (Na\(_v\)) channels (or an absence of Na\(_v\) channels in the extreme case). In the second scenario (Fig. 3B), we will consider that the composition of Na\(_v\) channels in the T-tubular membrane is the same as in the surface sarcolemma. In the third (Fig. 3C), we will consider that the density of Na\(_v\) channels is higher in the T-tubular membrane than in surface sarcolemma. We initially assume that the axial resistance of the T-tubules and the intratubular potential gradient are negligible. This is equivalent to assuming a homogeneous membrane potential throughout the cell surface and T-tubular membranes. Thus, we assume that the T-tubular membrane potential quickly follows the lateral membrane potential, and that extracellular potentials in the T-tubular lumen are small \[50, 60\]. These three scenarios therefore correspond to changing the total capacitance and/or the number of channels in the membrane.

The first scenario (few or no Na\(_v\) channels in the T-tubules, Fig. 3A) corresponds to scaling the specific membrane capacitance \( c \) in the cable equation (Eq. 6). Intuitively, adding capacitance would increase the load to excite the tissue and slow conduction. This scenario is however difficult to assess theoretically in a
quantitative manner. It is reminiscent of myofibroblasts connected to myocytes with a high level of gap junctional coupling, although myofibroblasts will render the resting membrane potential of the myocytes less negative. The effects of such an additional capacitive load have been studied by Jacquemet and Henriquez [65].

Huxley [66] as well as Jack, Noble and Tsien [49] have addressed our first scenario theoretically in detail, and showed that the kinetic properties of the voltage-gated ion currents also determine the exact relationship between conduction velocity ($\theta$) and capacitance ($c$). Nevertheless, $\theta$ is expected to decrease when $c$ increases. The relationship between $\theta$ and $c$ lies between

$$\theta \propto c^{-1/2}. \quad \text{(Eq. 8)}$$

and

$$\theta \propto c^{-1}. \quad \text{(Eq. 9)}.$$
These relationships are illustrated in Fig. 3D. Thus, in the extreme case of this first scenario and assuming that T-tubules contribute 50% to cell capacitance, \( \theta \) may be decreased by an amount between 29% (Eq. 8) and 50% (Eq. 9).

The second scenario (same ion channel composition of the T-tubular membrane as the surface sarcolemma; Fig. 3B) corresponds to changing \( I_{\text{ion}} \) and \( I_{\text{cap}} \) by the same factor. If the two summands of the right hand side of the cable equation (Eq. 6) are scaled by a factor \( n \), we can observe that the cable equation remains the same if \( \sigma \) is also scaled by \( n \). Thus, in terms of effects on \( \theta \), scaling \( I_{\text{ion}} \) and \( I_{\text{cap}} \) by \( n \) is equivalent to scaling \( \sigma \) by \( 1/n \) and we can expect that \( \theta \) will behave as

\[
\theta \propto n^{-1/2}. \quad \text{(Eq. 10)}
\]

Therefore, if T-tubules contribute 50% to cell capacitance, \( \theta \) may be decreased by up to 29% (Fig. 3E).

Finally, the third scenario corresponds to adding relatively more Na\(_v\) channels than membrane (Fig. 3C). In the extreme situation of adding Na\(_v\) channels only, \( \theta \) is expected to increase [63]. In a computer model of conduction, King and Fraser [44] have shown empirically that \( \theta \) is linearly related to the logarithm of \( g_{\text{Na, max}} \), i.e.,

\[
\frac{\theta}{\theta_{\text{ref}}} = 1 + k \log \left( \frac{g_{\text{Na, max}}}{g_{\text{Na, max, ref}}} \right) \quad \text{(Eq. 11)}
\]

where \( \theta_{\text{ref}} \) is the reference control velocity, \( g_{\text{Na, max}} \) is the maximal conductance of Na\(_v\) channels, \( g_{\text{Na, max, ref}} \) is the reference control \( g_{\text{Na, max}} \), and \( k \) is a model-dependent constant that must be determined empirically. For conduction in the Luo-Rudy dynamic model as investigated by Shaw and Rudy [63], we estimated \( k \) to be 0.308. The function given by Eq. 11 is illustrated in Fig. 3F. Thus, \( \theta \) may increase due to the larger Na\(_v\) channel density, but this effect may be opposed by the increase of \( c \). The net effect will depend on the exact change of \( g_{\text{Na, max}} \) and \( c \). This scenario however appears unlikely in the light of recent experimental data, since studies addressing Na\(_v\) expression in the T-tubules have suggested that at most 29% of Na\(_v\) channels are present in the T-tubules [67], which may be an overestimation (see Section 3.4.2 for a detailed discussion).

Next, we conjecture about the additional effects of T-tubular exit resistance and associated large T-tubular potential gradients on macroscopic conduction velocity (along the surface membrane and throughout cardiac tissue). On the one hand, a substantial exit resistance would tend to decouple the capacitive load of the T-tubules, thereby accelerating macroscopic conduction. In the limit of a very large resistance, the load of the T-tubules would be fully decoupled, corresponding to a detubulated situation. However, the experimentally observed time constants in the range of 100-200 \( \mu \)s [51, 58] and the small potential gradients
reported in modeling studies [50, 60] suggest that the exit resistance is small. The reported time constants of 100-200 μs are moreover much shorter than the typical timescale of the upstroke of the action potential (~ 1 ms), suggesting that the capacitive load only has small effects on the action potential upstroke. These considerations are in line with those of Jacquemet and Henriquez [65], who studied the effects of the capacitive load of fibroblasts in various myocyte-fibroblast coupling regimes. Of note, a large exit resistance would slow conduction along the T-tubules.

On the other hand, we note that the exit resistance of the T-tubules will directly influence the level of self-attenuation. If the exit resistance is small, self-attenuation will be negligible, and we expect no influence on conduction velocity. If the exit resistance is sufficiently large, self-attenuation of the sodium current will appear because of the large negative tubular potential, which would tend to slow conduction. Current results of computer simulations suggest that self-attenuation is in the range of a few percent [50, 60], and the resulting conduction slowing would also be in the percent range. An increase in the tubular resistance would however also decrease the capacitive load (see above), which may compensate the conduction slowing. In the extreme case of a very large exit resistance, we would retrieve once more the detubulated situation. Since there is no experimental proof of sodium current self-attenuation in the T-tubules for the moment, we must limit ourselves to speculation.

In summary, the reality most likely lies somewhere between the first two scenarios, so the presence of T-tubules will decrease θ. Conversely, loss of T-tubules may accelerate conduction. To explicitly investigate all these factors, including T-tubule resistance, intratubular potential gradients, and self-attenuation of the sodium current, specific modeling approaches need to be developed.

3.4 Action potentials in T-tubules?

A logical follow-up question to the passive conductive properties of T-tubules is: are action potentials generated in T-tubules? To generate an action potential, the right collection of depolarizing currents (conducted by voltage-gated sodium and calcium channels) and repolarizing current (conducted by potassium channels) must be present. This section discusses evidence regarding the presence or absence of ion channels in the T-tubules.

Firstly, it must be noted that data regarding T-tubular action potentials and ion channel expression require careful interpretation. In functional experiments, the cell capacitance is used to calculate membrane surface area, often assuming that 1 μF corresponds to 1 cm² membrane area [35]; thus, dividing whole-cell current by capacitance gives current density. Any difference in current density between normal and detubulated cells should indicate that the density of ion channels differs between T-tubular and non-T-tubular
sarcolemma [26, 67]. A caveat of this assumption is that measuring the capacitance underestimates the T-tubular membrane surface, since optical measurements give much higher fractions of T-tubular membrane (65% versus 32% from capacitance-based measurements [30]). One partial explanation for this discrepancy is the fact that capacitance is often determined too straightforwardly when the current transient evoked by a voltage-clamp step is simply integrated and divided by the commanded voltage step [68]. This method underestimates the capacitance if the opening of the cell at the pipette is imperfect, yielding a higher access resistance and a concomitant increase in capacitance charging time, and if the cell cannot be approximated as a single resistance in parallel with a single capacitor, which is the case for structurally complex cells such as cardiomyocytes. As a result, the apparent time constant of the current transient charging the capacitor increases, and from a certain time constant value, the “tail” of the current transient is usually excluded from the integration. The resistance that certain structural obstacles pose to the capacitive current may moreover be so high that the membrane depolarizes too late for the ion channels in this region to be activated. The ability to depolarize the entire cell membrane on a satisfactorily short time scale is called space clamp. In other words, in a voltage-clamped cell, the transmembrane potentials will differ between different portions of the membrane due to electric field interactions with geometric barriers.

Taken together, these pitfalls lead to an underestimation of membrane area of normal (non-detubulated) cells and an overestimation of current density. Therefore, the fraction of T-tubular current density must be lower than the capacitance-based measurements implicate, and ion channel density in T-tubules will be overestimated.

Indeed, when ion channel densities are determined from functional experiments, only rarely the results are corrected for cell surface [30]. Pásek et al. [30] reconcile the discrepancy between membrane area estimations by assuming a lower specific capacitance of T-tubules due to the relatively high cholesterol concentration in T-tubular membranes, and correcting for incomplete detubulation, setting T-tubular membrane surface at 49% instead of 32%. This reconciliation should lead to a decrease in current attributed to the T-tubules, even after correcting for the fact that formamide treatment leaves about 7% of T-tubules unaffected. A recent study however invalidated the hypothesis that cholesterol reduces the specific capacitance [69]. The discrepancy between electrophysiological and optical surface measurements may also be sought in difficulties in obtaining a complete space clamp in cardiomyocytes.

Comparing currents between whole-cell and detubulated cells may be complicated by more factors. Firstly, applying an osmotic shock to a cell may elicit unwanted effects which have not yet been investigated. The sudden internalization of membrane proteins may result in other acute membrane remodeling responses; the remodeling of structural proteins may affect protein expression at the membrane. Secondly, the sodium current in deep T-tubular segments may self-attenuate (see Section 3.2) [50, 60], so whole-cell peak sodium
current may not reflect the full available sodium current. Taken together, it is crucial to keep in mind that ion current densities between normal and detubulated cardiomyocytes should not be compared without taking into account the limitations mentioned above.

Immunohistochemistry data also require careful interpretation. Ideally, any staining for T-tubular proteins should be combined with a fiducial T-tubular marker and a knock-out model to assess the specificity of the antibody. Most publications addressing ion channel expression in T-tubules miss one or both of these quality controls. As a good example, Eichel et al. [70] show that the staining for the ion channel regulator CASK follows a regular striated pattern in cardiomyocytes, which may indicate T-tubular localization. The signal however did not colocalize with RyR2 and Ca\textsubscript{v}1.2, so the authors concluded that CASK was not expressed at the T-tubules [70]. We will therefore interpret published data on ion channel expression in T-tubules carefully.

### 3.4.1 Calcium-handling proteins

The **voltage-gated calcium channel** consists of an \(\alpha\)-subunit with multiple auxiliary \(\beta\)-, \(\alpha_\delta\)- and \(\gamma\)-subunits [71]. The predominant cardiac splice variant of the \(\alpha\)-subunit is Ca\textsubscript{v}1.2 (reviewed in [5]). Electrophysiological studies in detubulated and normal rat cardiomyocytes have determined that 60-75% of calcium current is conducted by calcium channels in the T-tubules [26, 72]. The authors did not normalize the calcium influx [26] or whole-cell calcium peak current [72] to the cell capacitance; the difference in calcium influx or peak current between control and detubulated cardiomyocytes already indicated which fraction of calcium channels was localized in the T-tubules. The previously stated problems with normalizing whole-cell currents by capacitance (**Section 3.4**) do therefore not apply to these studies.

Most calcium channels in the T-tubular membrane are close (~ 15 nm) to RyR2 in the SR, forming dyads [73]. This close proximity allows calcium ions entering the cell through Ca\textsubscript{v}1.2 to quickly and efficiently bind RyR2, followed by calcium-induced calcium release from the SR and subsequent sarcomere shortening. RyR2 and Ca\textsubscript{v}1.2 stay so close together because junctophilin-2 bridges the membranes of the SR and T-tubule to stabilize the dyad [39]. RyR2 and Ca\textsubscript{v}1.2 both bind Bin1, a banana-shaped protein that assists in trafficking and clustering of calcium channels [41].

The **sodium-calcium exchanger NCX** is predominantly expressed in the T-tubules. With detubulation experiments in rat cardiomyocytes, 60% of total NCX current – both outward and inward – was determined to originate from the T-tubules, corresponding to a ratio NCX current carried by the T-tubules/surface membrane of 1.5 [74]. Another study stated on the basis of similar detubulation experiments that the ratio of NCX-carried current in T-tubular to that carried in surface membrane varies from 1.7 to 25 [75]. Both
aforementioned publications determine T-tubular NCX currents by comparing current densities, which methodological problems have been outlined extensively in Section 3.4. It is therefore likely that T-tubular NCX currents make up less than 60% of total. When we assume that T-tubules account for 50% of membrane surface [30], while the T-tubular membrane surface measured in [74] only was 32% of total, the T-tubular NCX current may be overestimated twofold. In other words, T-tubular NCX current may make up 30% rather than 60% of total current. Still, this does not necessarily conflict with the generally accepted idea that NCX plays an important role in calcium handling in T-tubules.

T-tubular NCX expression correlates with the function of T-tubules as the main stage for calcium handling: a relatively high concentration of Ca,1.2 and RyR2 at the T-tubules facilitates a higher calcium conduction at the T-tubules than at the surface sarcolemma. This also requires more NCX activity to remove cytoplasmic calcium again. Indeed, NCX seems to be close to dyads, as well as SERCA, which pumps calcium ions back into the SR (reviewed in [27]), facilitating fast contraction as well as fast relaxation of the sarcomeres. Ca,1.2 and the NCX isoform NCX1 are both reported inside and outside of caveolae [37, 76]. Dyadic Ca,1.2 and NCX1 however most likely occur outside of caveolae, since caveolar Ca,1.2 seems to play a role in the hypertrophy pathway and not to participate in contraction [37].

Ca\(^{2+}\) ATPase also extrudes calcium ions from the cytoplasm. Its activity is reported to be confined to the T-tubules in rat cardiomyocytes based on recordings in cardiomyocytes with the Ca\(^{2+}\) ATPase blocker carboxyeosin [77].

3.4.2 Sodium-handling proteins
The cardiac voltage-gated sodium channel Na,1.5 is crucial for generating the rapid upstroke of the action potential in cardiac cells. It is the first channel to open when the membrane depolarizes. It is unclear whether voltage-gated sodium channels are required in T-tubules, since the T-tubular membrane will depolarize quickly due to passive conduction alone, yet T-tubular Na, channels would support conduction (see Section 3.3). Data regarding T-tubular sodium channel expression are multi-interpretable or inconclusive. Brette et al. concluded that 29% of sodium current is present in T-tubules because sodium current density in detubulated cardiomyocytes and normal cells were similar while cell capacitance decreased by 29% after detubulation [67]. After correcting for the larger T-tubular membrane area in untreated cells, however, current density would decrease. As a result, the calculated T-tubular sodium current fraction would be lower. We may also conclude that T-tubules do not contain Na, channels because the same study showed that whole-cell sodium current did not differ between normal and detubulated cells [67].
When we consider immunohistochemistry data for Na\textsubscript{v}1.5, we have to conclude that a reliable co-staining of Na\textsubscript{v}1.5 and a T-tubular marker is still missing. Only a striated pattern of Na\textsubscript{v}1.5 has supported the in our view premature conclusion that Na\textsubscript{v}1.5 would be expressed in T-tubules [78-80].

Although Na\textsubscript{v}1.5 is the main cardiac isoform of the voltage-gated sodium channel family, other isoforms have been reported in the heart [81]. It is important to keep in mind however that the expression of neuronal isoforms remains controversial. A recent report reported only Na\textsubscript{v}1.5 and Na\textsubscript{v}1.4 expression using next-generation RNA sequencing data from murine cardiomyocytes [82], yet other reports suggest that the neuronal isoforms Na\textsubscript{v}1.1, Na\textsubscript{v}1.3 and Na\textsubscript{v}1.6 may be enriched in T-tubules [83-85] and even contribute to arrhythmias [86]. These neuronal channels have a lower activation threshold and inactivate more rapidly, which may facilitate action potential propagation into the T-tubules [5] – yet the fast passive propagation of depolarization renders T-tubular sodium channels relatively unessential. The immunohistochemistry data showing Na\textsubscript{v}1.1, -1.3, and -1.6 in T-tubules of murine cardiomyocytes need careful interpretation because no knock-out control is presented for the antibodies [84, 86]. Another publication presents striated patterns of neuronal isoforms in immunohistochemistry data [85], whereas pre-absorbing the antibodies with peptides inhibits any staining, suggesting specificity of the antibodies [85]. Westenbroek et al. quantified the relative expression of Na\textsubscript{v}1.1-6 in murine cardiomyocytes [87]. Na\textsubscript{v}1.5 was only found at the intercalated disc and lateral membrane. Na\textsubscript{v}1.1 and -1.3 showed a striated pattern, but compared to Na\textsubscript{v}1.5, the signal intensity was relatively low, as was expected from protein expression data [88]. A T-tubular marker was however missing [87].

On a functional level, current conducted by neuronal channels was recorded in murine ventricular cardiomyocytes treated with β-scorpion toxin, a specific activator of neuronal sodium channels [85]. Na\textsubscript{v}1.6-deficient mice moreover showed prolonged calcium transients in cardiomyocytes, suggesting a functional link between Na\textsubscript{v}1.6 and excitation-contraction coupling [89]. In wild type mice, Na\textsubscript{v}1.6 staining showed a striated pattern that weakly colocalized with α-actinin, but we have to consider that α-actinin is not an ideal T-tubular marker. Na\textsubscript{v}1.6-deficient mice did not show any striated pattern, which suggests the anti-Na\textsubscript{v}1.6 antibody binds specifically. In a model for CPVT (catecholaminergic polymorphic ventricular tachycardia), an arrhythmic persistent sodium current seemed to be conducted by neuronal Na\textsubscript{v} isoforms [90]. The immunofluorescence data from this paper however show too high background signals to allow any conclusions on Na\textsubscript{v} isoform expression.

Taken together, despite many reports on the expression of sodium channels in T-tubules, T-tubular Na\textsubscript{v} expression remains debatable. A de facto inward sodium current by NCX should be expected in any case because of the high electrochemical gradient.
Lastly, NKA is expressed in T-tubules [91]. Especially the α₂-subunit is enriched in T-tubules, whereas the α₁-subunit is more uniformly expressed over the sarcolemma [91].

### 3.4.3 Potassium-handling proteins

Although the subcellular localization of potassium channels is less studied, several potassium currents seem to be conducted in the T-tubular systems of mice and rats. These include the inward rectifier current $I_{K1}$ [58] (conducted by $K_{ir}2.1$, 2.2, and/or 2.3), the steady-state current $I_{ss}$ ($K_{ir}1.2$ and/or -2.1), the transient outward current $I_{to}$ ($K_{ir}4.2$ and/or -4.3), the slow delayed rectifier current $I_{Ks}$ ($K_{ir}7.1$), and the rapid delayed rectifier current $I_{Kr}$ ($K_{ir}11.1/hERG$) [30]. Based on detubulation experiments, $I_{K1}$, $I_{to}$, $I_{Kr}$, and $I_{Ks}$ seem to be divided equally over T-tubular and surface sarcolemma [92]. The inward-rectifying current $I_{K1}$ is especially high in T-tubules because potassium accumulates in the small extracellular space during the repolarization phase of the action potential, increasing its driving force [58]. Other currents were found to be enriched in rat T-tubules, including $I_{ss}$ [92]. Interestingly, only $K_{ir}2.1$-encoding mRNAs were detected in isolated murine cardiomyocytes [82], suggesting $K_{ir}1.2$ may not contribute to $I_{ss}$ in these cells.

The T-tubular fraction of $I_{KATP}$ has, to our knowledge, not yet been determined. It is however likely that at least one of the $I_{KATP}$-conducting channels $K_{ir}6.1$ and -6.2 is expressed in T-tubules, since $I_{KATP}$, as well as $I_{K1}$, decreases along with the loss of T-tubules in cardiomyocytes that have been taken in culture (reviewed in [5]). $K_{ir}6.2$ has been associated with the T-tubular protein ankyrin-B, although these immunohistochemistry data lack a fiducial T-tubular marker [93].

Considering immunohistochemistry data, T-tubular expression has been shown relatively convincingly for $K_{ir}11.1$ by co-staining with myosin-binding protein C [94]. A thorough live-cell imaging study for $K_{ir}2.1$ and $K_{ir}1.4$ in rat cardiomyocytes shows a T-tubular-like pattern, including axial elements, but unfortunately lacks a T-tubular marker [95]. $K_{ir}2.1$, -2.3, -4.2 and TASK-1 show clear co-localization with the membrane marker wheat germ agglutinin in murine (K₂2.1), canine (K₂2.1 and -2.3) and rat (K₂4.2 and TASK-1) ventricular cardiomyocytes [96-99]. For K₂4.3 only a striated pattern was shown without T-tubular marker in canine cells [100]. Note that potassium channel composition probably differs between species, since K₂2.3 mRNA was not detected in isolated murine cardiomyocytes [82] while K₂2.3 protein was found in canine cells [97]. Interestingly, TASK-1 has been recently defined as an atrial-specific current [101], which seems to contradict these immunohistochemistry and mRNA expression data – although it must be noted that we did not find any recordings of ventricular TASK-1 currents in the literature.
3.4.4 Other ion-handling proteins

Chloride channels are not well researched in T-tubules. Scanning ion conductance microscopy (SICM) experiments revealed that these channels are close to the openings of T-tubules at the lateral membrane [102]. mRNA data show that the voltage-sensitive chloride channels CLCN4, -7, -3, -6 and -1 are expressed in murine ventricular myocytes (in order of expression level) [82], but their roles in cardiac physiology remain to be investigated.

Another ion channel family which roles in cardiac physiology are just emerging is the transient receptor potential (TRP) superfamily of non-selective cation channels. The vanilloid subfamily member TRPV4 is expressed in T-tubules of aged mice and human cardiomyocytes [103]. This is based on immunohistochemistry data in which TRPV4 and the membrane marker caveolin-3 are co-stained. TRPV4 seems to play a role in calcium cycling, but in the aged heart it may contribute to stress-induced cell damage. Of note, TRPV4-encoding mRNA is not found in murine cardiomyocytes [82], which illustrates that even among mouse models, ion channel composition may differ.

3.5 Ion dynamics in T-tubules

The small luminal space of T-tubules can affect the ion diffusion and therefore the driving force of ion channels (discussed in Section 3.2). Diffusion rates of ions are slower in T-tubules than in other extracellular spaces, with a markedly higher delay for divalent than monovalent ions [18, 19]. Functionally, an extracellular change in calcium concentration reaches the T-tubules 2.3 seconds later in guinea pig cardiomyocytes, indicating a “fuzzy space” [104]. This diffusion delay is related to the membrane microfolds that are shaped by Bin1 [42]. These microfolds are proposed to protect against arrhythmias, since excitability of the cell would increase when ions can diffuse faster (reviewed in [5]). The finding that T-tubules of animals with higher heart rates are narrower than in low-heart-rate animals [5] might indicate that the relatively high calcium depletion and potassium accumulation rates in the narrow tubules of high-heart-rate animals are of evolutionary benefit.

Theoretically, calcium influx during an action potential may deplete the T-tubular lumen of calcium and reduce the driving force of the calcium channel. This may cause self-attenuation of the calcium current and thus shorten the action potential [45], in a manner similar to what was proposed in the context of intercalated discs for the late sodium current during the plateau phase (see Section 3.2) [62]: the sooner the calcium channels close, the sooner the potassium current repolarizes the membrane. A strong depletion of calcium and accumulation of potassium in the T-tubular lumen is however unlikely due to the physiological need for significant electrochemical gradients of calcium and potassium, which will be secured by NCX and
An important general remark is that we cannot speak of a shorter or longer action potential in T-tubules compared to the rest of the sarcolemma as one might think from the finding that detubulated rat cells show shorter action potentials due to the lost calcium current [26]. The membrane domains of the cell are electrically too tightly coupled [105, 106]. Detubulation may however affect action potential duration of other species differently as rat cardiomyocytes show a relatively unique calcium handling profile compared to other mammals, which is partly explained by a reduced NCX/SERCA2a ratio [67].

3.6 T-tubular remodeling in disease

Heart failure has been associated with T-tubular remodeling, characterized by T-tubular widening and an increase in tortuosity [107], the loss of the microfold regulator Bin1 [42], collagen deposition in T-tubules [11, 51], dyad uncoupling [108], physical obstruction of T-tubules [47], and sheet-like T-tubular remodeling in end-stage heart failure [36].

The loss of Bin1 observed in human heart failure has been suggested to be arrhythmogenic as Bin1 knockout mice show a loss of microfolds and are more vulnerable to arrhythmias due to changes in ion dynamics as discussed in Section 3.5 [42]. Sheet-like remodeling is similarly expected to be arrhythmogenic [36]. On the other hand, collagen deposition in T-tubules, which is also seen in human heart failure, and physical obstruction of T-tubules would again slow ion diffusion [10, 11, 51]. Understanding the interplay of these effects will need further investigation.

T-tubular remodeling likely disturbs optimal excitation-contraction coupling as calcium dynamics greatly depend on T-tubular microstructure [109] and on the association and membrane expression of voltage-gated calcium channels and the sodium-calcium exchanger [110]. Indeed, a hallmark of T-tubular dysfunction is the increase in orphaned RyR2 at the Z-lines [108]. This causes a dysregulation of calcium release as diffusion distances for calcium increase, and the calcium transient slows and broadens. In turn, contraction of the cell desynchronizes and slows. Expectedly, detubulated cardiomyocytes show comparable contraction kinetics as failing cells (reviewed in [10]). Interestingly, in end-stage heart failure patients, relieving the mechanical load with a left-ventricular assist device (LVAD) improves cardiac function but only if the T-tubular system is intact, suggesting that calcium handling deficits may be reversible whereas sheet-like T-tubular remodeling is not [36].

It may be valuable to compare depolarization delays in T-tubules from healthy and diseased cells to assess the passive electrical effects of T-tubular remodeling, although other processes associated with T-tubular remodeling may play greater roles in pathophysiology. Since heart failure is also associated with a reduction of Bin1 expression, and Bin1 is important for voltage-gated calcium channel trafficking to the T-tubules.
[33], the resulting dysregulation of this trafficking may further impair calcium dynamics. In summary, the T-tubular remodeling associated with heart failure probably affects cardiomyocyte function via different pathways, yet the effects of T-tubular remodeling on T-tubular depolarization remains to be determined.

4. T-tubules in atrial cardiomyocytes

Although T-tubules support a high heart rate in mammalian ventricular cardiomyocytes, heart rates of atrial cells are just as high with a far less dense T-tubular system [20]. Atrial cells in many species can be subdivided in three sub-populations based on their T-tubular network: dense, semi-dense and disorganized, or absent [8]. In pig and rat atria, cardiomyocytes with a well-organized T-tubular network populated the epicardial layer, whereas the endocardial layer mostly contained empty cardiomyocytes [111]. The degree of tubulation furthermore correlated well with synchronicity of calcium release. This may suggest that T-tubules at the epicardial layer are involved in synchronizing the contraction across the atrial wall [111], and that β-adrenergic stimulation has a faster and greater effect on cells close to the epicardium. Interestingly, one recent study [112] reports that human and murine atrial cardiomyocytes contain relatively many axial tubules, which run over the long axis of the cell. The sparse transverse tubules conduct the excitation of the cell from the surface sarcolemma to the axial tubules, while axial tubules contain much more dyads than transverse tubules, thus allowing quick excitation-contraction coupling around these axial tubules [112].

5. Discussion

This work assessed passive and active electrical properties of T-tubules from which we distilled the physiological function of T-tubules. Generally, the roles of T-tubules are twofold: firstly, T-tubules provide a large calcium-handling stage to allow the cardiomyocyte to generate a relatively high force, and secondly, the relatively large membrane surface facilitates the cell to respond faster to signaling molecules than non-tubulated cells. Importantly, T-tubules are not crucial to sustain a high heart rate – atrial cardiomyocytes contract just as fast [20], and non-tubulated cardiomyocytes from certain avian species even faster [113, 114]. Within mammalian ventricular cardiomyocytes however there is a relationship between heart rate and T-tubular density.

The publications discussed in this review strengthen the hypothesis that T-tubules facilitate rapid and efficient excitation-contraction coupling throughout the entire volume of the cell [26]. The length constant of ~ 68 μm [51] and time constant of ~ 200 μs [58] indicate that electrotonic responses allow depolarization of the entire T-tubular membrane within physiological timescales.
When assessing conduction velocity in cardiac tissue, T-tubules most likely slow conduction from a theoretical point of view. Since a tubulated cardiomyocyte has a greater membrane surface, this additional capacitive load may slow down and delay depolarization. Ion channels in T-tubules, especially Na$_v$ channels, will mitigate this delay. Given that the concentration of Na$_v$ channels in the T-tubule is most likely lower than that in the surface sarcolemma, and assuming T-tubular membrane accounts for 50% of the membrane (which is the case at least in mice and rats [30]), T-tubules will slow down conduction velocity by 29-50%, while loss of T-tubules may accelerate conduction.

Next, we evaluated evidence regarding the expression of voltage-gated sodium channels in the T-tubules. Some neuronal isoforms seem to be expressed in low levels; especially Na$_v$1.6 seems to be expressed in the T-tubules of a mouse model [89]. It remains unclear if these low levels could carry enough current to initiate an action potential. The expression of neuronal channels moreover varies greatly between species and even mouse strains. Concerning the cardiac isoform Na$_v$1.5, T-tubular expression remains debatable, despite several reports pointing in that direction. Generally, it is important to keep in mind that sodium channels are not required to depolarize T-tubular membranes. The often-used term “action potential propagation” in T-tubules therefore should be interpreted as “either passive or active”.

In terms of repolarizing currents, the expression of voltage-gated potassium channels differs greatly between species and mouse models. It is probable that the expression levels suffice to repolarize the membrane to balance the high expression level of voltage-gated calcium channels. Considering ion channel expression data in general, the authors would like to stress that critical evaluation of research data is pivotal, especially when functional detubulation experiments and immunohistochemistry data are concerned.

In cardiac disease, T-tubules often remodel [9]. As a result, T-tubules widen, the tortuous membrane microfolds disappear, dyads uncouple, T-tubules close off to the extracellular space, calcium transient weakens, and T-tubules may take on a sheet-like shape while collagen content in the lumen is increased.

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AUTHOR CONTRIBUTIONS

Sarah Vermij: conceptualization, methodology, visualization, writing- original draft preparation, reviewing and editing.
Hugues Abriel: supervision, writing - reviewing and editing.

Jan Kucera: visualization, writing - reviewing and editing.

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REFERENCES


[34] K.R. Levin, E. Page, Quantitative studies on plasmalemmal folds and caveolae of rabbit ventricular myocardiocells, Circ. Res., 46 (1980) 244-255.


Ryanodine receptors after acute junctophilin knockdown in mice, Circulation, 123 (2011) 979-988. DOI: 10.1161/circulationaha.110.006437.


A. Gopinathan, Neuronal space constants in various geometries and physiological conditions, arXiv, (2013).


A. Greer-Short, S.A. George, S. Poelzing, S.H. Weinberg, Revealing the Concealed Nature of Long-QT Type 3 Syndrome, Circ Arrhythm Electrophysiol, 10 (2017) e004400. DOI: 10.1161/CIRCEP.116.004400.


