



Effective detection of weak terahertz pulses in electro-optic sampling at kilohertz repetition rate

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The standard terahertz (THz) detection mechanism known as electro-optic sampling can be improved in sensitivity by biasing the polarization of the sampling field. In this work, we show theoretically and experimentally how weak signals can be amplified without inducing distortions. Our study identifies the influence of THz field strength, the polarization quality, and biasing amplitude on signal amplification and distortion. Here we present a distortion-free amplification of a factor of 28 while at the same time reducing the measurement time significantly. © 2019 Optical Society of America

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1. INTRODUCTION

Measuring terahertz (THz) fields in the picosecond range requires the mixing of a THz pulse with a sampling light pulse of much shorter duration. One method to measure such a pulse is electro-optic sampling (EOS). EOS is based on the optical Kerr effect in the limit of the THz electric field oscillating slowly compared with the gating light pulse duration. Concise and in-depth description of the process has been given in the frequency domain [1] from the optical intensities measured in the process [2], and starting from the polarization components [3]. One typical use of this detection scheme is employed in optical pump THz probe studies, where electronic properties of materials are optically probed, making it an especially powerful tool in optoelectronics research. For many studied materials, the lifetimes of optical excitations reach microseconds; thus, amplified Ti:Sa laser systems with low repetition rates in the kilohertz (kHz) range are often necessary in materials science. Compared with MHz systems, acquiring enough laser shots to reach an adequate signal-to-noise ratio significantly increases measurement time, motivating the improvement of the detection scheme. Several methods were previously employed, all relying on the excess of sampling pulse intensity available when using kHz repetition rate systems, where the damage thresholds of the optics or detectors in use are the limiting factors. Typically, the sampling pulse intensity is attenuated to avoid saturation of the detector. As EOS basically is a translation of field strength into change of polarization, it is possible to take advantage of the excess energy to improve detection sensitivity with careful manipulation of the polarization in the detection process. Previous studies achieved that by biasing the input polarization with an optical retarder before [4] and/or after the optical mixing within the detection crystal [3,5]. In a comparable scheme, a

series of wedges was used to control the polarization [6]. The bias can also be introduced using nearly crossed polarizers [3]. Unbalanced detection, which is employed in THz spectrometers based on asynchronous optical sampling (ASOPS) [7], is also closely related to these. Within these schemes, distortions of the measured THz pulse have been identified as a key limiting factor. Distortion-free amplification techniques have been reported, but the usability is limited since either two measurements are needed to reconstruct the true waveform [5] or they are not applicable to detect weak THz fields [4]. In this paper, we show how the strength of the THz field itself is a key parameter to limit amplification and distortion as hypothesized before [6]. Additionally, we show how distortion-free amplification is possible at a small THz field, allowing us to detect weak signals with increased sensitivity. To this purpose, we used a simple scheme presented before by Johnson *et al.* [5] and Brunner *et al.* [3] and apply the Mueller formalism to propagate the polarization state through the detection. This allows us (i) to account for the imperfect polarization of the sampling pulse (also defined as gate pulse in this work) and (ii) to understand the origin and the link between distortions, THz strength, and amplification in detail. We present the experimental setup followed by its theoretical description and demonstrate its applicability in the detection of weak THz fields.

2. METHODS

A. Setup for Detectivity Enhanced Electro-Optic Sampling of THz Pulses in the kHz Range

Figure 1 shows the typical setup for EOS with bias as published by Johnson *et al.* [5]. We take our 1 kHz, 35 fs light pulse at 800 nm and split it into two beams for the THz generation

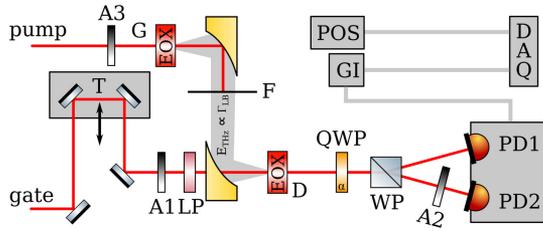


Fig. 1. Experimental setup: generation crystal, G; detection crystal, D; quarter-wave plate, QWP; polarizing beam splitter, WP; two balanced photodiodes, PD1, PD2; variable attenuators, A1, A2, A3; linear polarizer, LP; PTFE filter, F; gated integrator, GI; position encoding, POS; directly from linear translation stage, T; data acquisition, DAQ.

(pump) and detection (gate). The pump generates the THz pulse in the generation crystal (G) via optical rectification in a 1 mm thick ZnTe crystal (EOX). The THz pulse is imaged with two off-axis parabolic mirrors into our detection crystal (D). A PTFE filter (F) transmits the THz pulse and blocks the residual 800 nm light. The linearly polarized gating beam (through LP: colorPol VIS-IR polarizer, CODIXX) is spatially and temporally overlapped with the THz pulse. The collinear spatial and directional overlap is achieved by drilling a hole into our focusing parabola and directing the gate through it. Temporal overlap is achieved by modifying the length of the gate's optical pathway with a translation stage (T). Scanning the gate arrival time over the THz waveform reconstructs the electric field by measuring its effect on the birefringence on the detection crystal [1,2]. This THz-induced birefringence (Γ_{LB}) is then analyzed by a series of quarter-wave plates (QWPs), Wollaston prisms (WPs), and balanced photodiodes (PD1, PD2).

Classically, the birefringent ZnTe detection crystal is oriented along the polarization of the incoming detection beam and the wave plate at 45° , such that ideally perfectly circular polarized light is split 50:50 from the WP [2]. Without THz, the diodes are balanced so the signal $S = I_1 - I_2 = 0$, with I_1, I_2 being the intensities on PD1 and PD2, respectively. With THz, that difference leads to a finite value of S .

In our implementation, (T) is a magnetic delay stage (V-508 PIMag, PI Instruments), with which fast scanning is employed to evenly distribute slow pulse-to-pulse fluctuations of the laser system over the whole waveform [8]. Here, rapid scanning is achieved by simultaneously reading the voltage from our gated integrator (GI) connected to the balanced PDs and the position of the stage, so that the position is recorded for every laser shot. We derive the position by externally sensing several contacts in the cable connecting stage and controller PI PIMAG C981, while pulling the voltage level to transistor–transistor logic (TTL) levels to read the position with the counter on a digital input.

To take special care about the polarization state within the detection, we sequentially add optics into the system. First, we ensure that LP exactly aligns with one output port of the WP such that all intensity is on PD2. Then, we introduce the

detection crystal EOX-D along one of its crystal axes to ensure minimal residual birefringence. Lastly, the QWP is added in a motorized rotational mount (K10CR1, Thorlabs) to reproducibly set angles with high precision.

With A2 set to no attenuation, the overall used intensity is set at A1 to ensure linear behavior on the diodes PD1 and PD2. To amplify the measured signal, we bias the gate by setting the wave plate at an angle of $\alpha \neq 45^\circ$ and attenuate (A2) after the WP to reduce I_2 and thus rebalance the intensities on the detector, as described in Refs. [3,5]. For a given damage threshold of EOX-D, one can increase the input intensity of the gating beam and still be at balanced detection. In an extreme case, one removes the QWP such that $I_2 = I_0$ and $I_1 = 0$ without THz. This resembles the cross-polarizer measurement scheme [9]. With strong THz pulses and a sufficiently clean polarization state, this leads to the detection of purely positive signals (homodyne) corresponding to strong distortions reported in Ref. [5]. For a set of wave plate angles, we measure the THz waveform keeping the intensity constant on the diodes.

B. Theoretical Description of Detectivity Enhancement in Electro-Optic Sampling

The effect of the THz field in the detection crystal (EOX-D) is to create a linear birefringence Γ_{LB} at 45° [10], proportional to the THz electric field strength. This way, it modulates the initially horizontally polarized gating beam. With QWP set at α and the intensity on the diodes kept constant with A1 and A2, propagation of the polarization state through the optical elements as a Stokes vector in Mueller formalism yields the signal S . With perfect polarizers ($r = 0$, r being the extinction ratio of the polarizing optics) and in the limit of small Γ_{LB} ($\Gamma_{LB} \ll 45^\circ$), we obtain

$$S(\alpha, \Gamma_{LB}) = \frac{\Gamma_{LB}}{\sin(2\alpha)} + \frac{1}{2} \cot(2\alpha)^2 \cdot \Gamma_{LB}^2.$$

The first term, linear in Γ_{LB} , is referred to as the heterodyne (or self-heterodyne) signal. It is linear in Γ_{LB} (THz field strength) and thus induces no distortion. The second term is the so-called homodyne signal. The homodyne term is quadratic in Γ_{LB} and is responsible for the distortions. At small angles α , the signal diverges because this equation ignores finite gating light intensity at the input. Subtracting two measurements at $\pm\alpha$ doubles $S(\alpha, \Gamma_{LB})$ while eliminating the homodyne component leading to an undistorted signal [3,5]. On the contrary, adding them yields twice the purely positive homodyne signal,

$$S(\alpha, \Gamma_{LB}) - S(-\alpha, \Gamma_{LB}) = \frac{2\Gamma_{LB}}{\sin(2\alpha)}, \quad \text{and}$$

$$S(\alpha, \Gamma_{LB}) + S(-\alpha, \Gamma_{LB}) = \cot(2\alpha)^2 \cdot \Gamma_{LB}^2.$$

In the Mueller formalism (see Appendix A), we account for imperfections in the polarization ($r > 0$). In theory, the signal is as follows:

$$S_r(\alpha, \Gamma_{LB}) = \frac{\cos(2\alpha)^2(-2 + 4r^2 - 2r^4) + (2 - 4r^2 + 2r^4)^2 \cos(\Gamma_{LB}) + (-2 + 4r^2 - 2r^4) \sin(2\alpha) \cdot \sin \Gamma_{LB}}{-(1+r)^4 + (1-r)^4 \cos(2\alpha)^4}. \quad (1)$$

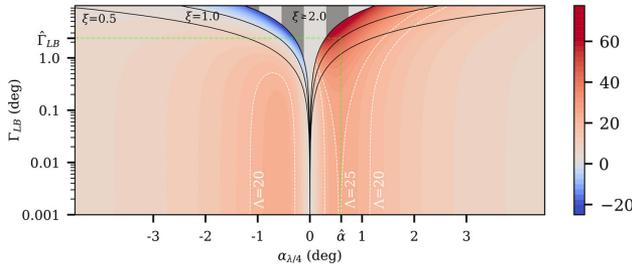


Fig. 2. 2D representation of amplification Λ_r [Eq. (2)] depending on Γ_{LB} and α at $r = 10^{-4}$. Black lines depict distortion $\xi_r(\alpha, \Gamma_{LB}) = 0.5, 1.0$ and $\xi_r(\alpha, \Gamma_{LB}) \geq 2.0$, respectively [Eq. (3)]. In the shaded region with $\xi_r \geq 2.0$, distortions dominate the signal. $\hat{\Gamma}_{LB} = 2.37^\circ$ corresponds to the peak field strength of the pulses used in the measurements presented in this work. $\hat{\alpha}$ depicts the optimized angle in this case. The green dashed line indicates both in the graph. White dashed lines depict boundaries of $\Lambda_r(\alpha, \Gamma_{LB}) = 20$ and $\Lambda_r(\alpha, \Gamma_{LB}) = 25$, respectively.

We define the amplification Λ at any $\pm\alpha$ relative to the signal measured at $\alpha = \pm 45^\circ$, respectively,

$$\Lambda_r(\alpha, \Gamma_{LB}) = \frac{S_r(\Gamma_{LB}, \pm\alpha)}{S_r(\Gamma_{LB}, \alpha = \pm 45^\circ)}. \quad (2)$$

The distortion ξ is defined as the ratio of homodyne (distorted) to the self-heterodyne (undistorted) signal,

$$\xi_r(\alpha, \Gamma_{LB}) = \left| \frac{S_r(\alpha, \Gamma_{LB}) + S_r(-\alpha, \Gamma_{LB})}{S_r(\alpha, \Gamma_{LB}) - S_r(-\alpha, \Gamma_{LB})} \right|. \quad (3)$$

In this way, the case $\xi = 0$ contains no homodyne signal resulting in no distortions, while for $\xi > 0$, the influence of the homodyne signal distorts the measured waveform. ξ increases with Γ_{LB} ; however, at small Γ_{LB} , we already find significant amplification with only negligible distortions, which is the focus of this study. To validate this theory experimentally, measurements are taken up to a maximally applied field (peak of the THz waveform) of $\hat{\Gamma}_{LB} = 2.37^\circ$ (in the distortion regime).

In Fig. 2, we display a 2D map of Λ_r in function of α and positive Γ_{LB} , with r set to $r = 10^{-4}$, which is close to our experimental conditions. We show the area with small α to highlight the regime where significant amplification Λ_r is possible. Λ_r is color coded from blue to red going from negative to positive amplification. Negative amplification is present when the homodyne contribution to the signal takes over. The black line at $\xi_r = 1.0$ for $\alpha < 0$ marks this case where $\Lambda_r = 0$, and, thus, the sign flips.

Distortions increase with higher Γ_{LB} and smaller α . The solid black lines depict signals with distortions of $\xi_r = 0.5$, $\xi_r = 1.0$, and $\xi_r \geq 2.0$, respectively. The area for $\xi_r \geq 2.0$, where the homodyne signal is twice the heterodyne contribution, is grayed out for illustrative purposes because the amplification diverges at large Γ_{LB} . Contour lines of $\Lambda_r = 20$ and $\Lambda_r = 25$ are shown as white dashed lines. The area below $\xi_r = 0.5$ is where distortions are small, such that areas of strong amplification without significant distortions can readily be identified.

We observe that when Γ_{LB} is small, we can use small α without distorting the recovered THz field significantly.

Furthermore, there is an optimal angle at $\hat{\alpha} = \sqrt{r}$ for which the strongest amplification is achieved while distortions can be neglected (if Γ_{LB} is sufficiently small). Then, the amplification at this angle converges to $\hat{\Lambda} = 1/(4\sqrt{r})$ when $\Gamma_{LB} \rightarrow 0$. In the particular case of this figure with $r = 10^{-4}$, the angle is $\hat{\alpha} = 0.01(0.6^\circ)$ and leads to an amplification of $\hat{\Lambda} = 25$. Because the amplification only approaches $\hat{\Lambda} = 25$, we do not find exactly the same for both $\pm\alpha$. This is due to the small homodyne contribution where the sign of the effect depends on the relative sign between α and Γ_{LB} . Once the introduced optical bias in α does not compete with the field strength Γ_{LB} , we are in the limit of distortion-free amplification of weak fields.

3. RESULTS AND DISCUSSION

In Fig. 3(a), we show the THz electric field transients at a representative selection of angles α of the QWP. The subtracted signals of the respective $\pm\alpha$ counterparts are shown in Fig. 3(b), showing nondistorted signals following [5]. The inset in Fig. 3(b) shows these same waveforms normalized to demonstrate that fact. We find a distortion-free amplification factor of ~ 24 , comparing the subtracted signals at 0.7° with a single measurement at 45° , which closely agrees with the predictions made in the previous section. We present two ways of avoiding the second measurement to reconstruct the waveform. First, if the applied field is weak compared with the optical bias, and the single measurement with $+\alpha$ will be distortion-free, meaning the measurement at $-\alpha$ will have the opposite sign such that the subtracted signals will just be 2 times the measurement at $+\alpha$. Second, the applied field can be reconstructed by fitting Eq. (1) to the distorted measurement data as shown in Fig. 4(b). For strongly distorted signals where the sign flips ($\Gamma_{LB} \geq \alpha$), the mapping with $S_r(\alpha, \Gamma_{LB})$ is not bijective anymore; thus, fitting requires knowledge of the original signals' shape. From

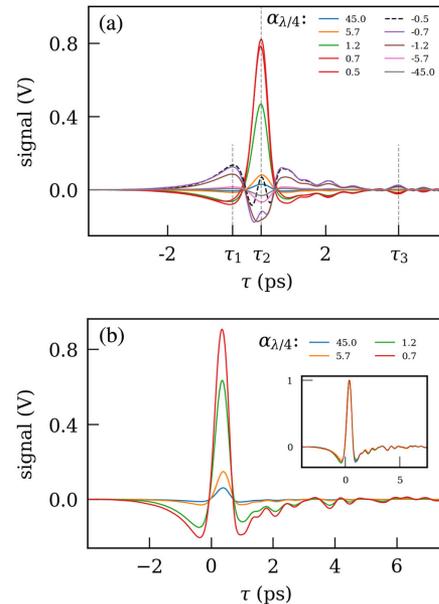


Fig. 3. (a) Collection of the measured signals at constant power on the diodes with different angles α of the quarter-wave plate. (b) Subtracted signals of corresponding quarter-wave plate angles ($\pm\alpha$) with their respective normalized signals in the inset.

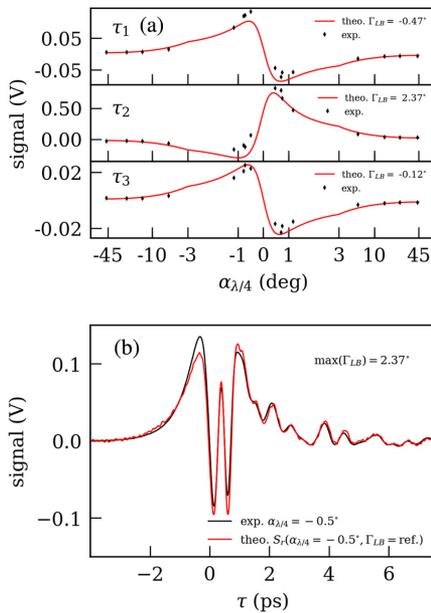


Fig. 4. (a) $S_r(\Gamma_{LB}, \alpha)$ and measured signal for three Γ_{LB} selected from τ in Fig. 3(a). (b) Measurement and calculated signal for the whole waveform at a small angle $\alpha = -0.5^\circ$.

Fig. 3(a), it becomes evident how small angles lead to higher measured signals, including an asymmetry at small angles that appears strongest at the peak of the THz field. We determine Γ_{LB} experimentally by compensating the THz signal rotating the QWP for each τ_i depicted in Fig. 3(a). The wave plate angle is directly related to Γ_{LB} . The τ_i corresponds to a minimum field at τ_1 , the field maximum at τ_2 , and a small nonvanishing field at τ_3 .

Figure 4(a) shows the signal at the 3 times τ_i as a function of α . We also calculate the signal $S_r(\Gamma_{LB}, \alpha)$ from the applied angle α at the QWP and the acting THz field Γ_{LB} from the measurement at 45° and plot it in red on a symbolic log scale [linear x -scale between -3° and 3° continued as $-\log(\alpha)$ and $\log(\alpha)$, respectively] to resolve small angles properly. The kink at 3° is where the scale changes. Note that at the maximum ($\tau_2 \hat{=} \hat{\Gamma}_{LB} = 2.37^\circ$), we are well within the distortion region depicted in Fig. 2 indicated by the green dashed line. To picture this, we show the single measurement at a QWP angle of $\alpha = -0.5^\circ$ in Fig. 4(b) in black. This waveform is highlighted in Fig. 3(a) as a black dashed line. The red line in Fig. 4(b) shows the signal we calculate from a measurement at 45° using Eq. (1). This demonstrates the case where the applied THz field is strong enough to flip the sign of the measured signal against the bias of the small angle.

We find excellent agreement between our experimental data and theory. The largest deviations are found at high Γ_{LB} , where our setup could not be optimized further to follow the large voltage changes with acceptable stability at high measurement speeds and allow for signals differing by orders of magnitude without any change in the electronic amplification. For the minimal field, which resembles the case of an intermediate field strength and the very weak just not vanishing field, we find that the method holds.

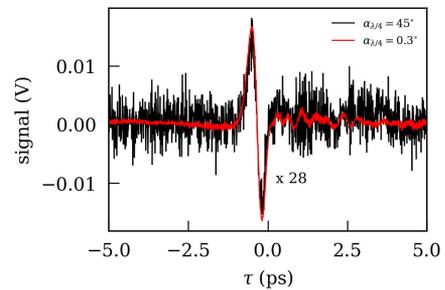


Fig. 5. Collection of measurements of a weak field in the standard method (black) compared with the amplified scheme (red). The black curve is multiplied by our amplification factor of 28 and took 6.5 times as many shots from our laser system.

For the purpose of this study and to rule out the influence of the detection system, we use the same electronic amplification for all angles inducing slight saturation features at measurements with high Γ_{LB} and small angles of α . Optimizing electronics for certain measurement conditions improves the sensitivity beyond what is presented here.

The THz pulses investigated in this study were generated from 800 nm pulses of 37.5 μJ . To significantly reduce Γ_{LB} , we reduced the pump intensity P_G incident on the generation crystal to 0.165 μJ , which corresponds to a reduction of 200 of the maximum of the THz field. The corresponding field strength then is $\Gamma_{LB} = 0.045^\circ$. We measured the THz field at the standard configuration of $\alpha = 45^\circ$ (black) and at $\alpha = 0.3^\circ$. From Fig. 2, we read an amplification factor of 28. In Fig. 5, we show the measurement at 45° times the amplification of 28 in black and the measurement at $\alpha = 0.3^\circ$ in red. While the black line is the collection of about 15 min integration time ($8.7 \cdot 10^5$ shots), the measurement at the amplified scheme took only 133 s ($1.3 \cdot 10^5$ shots). This amplification factor of 28 and the reduced measurement time of a factor of 6.5 demonstrates the improvement at measuring comparably weak THz pulses. The THz fields at the crystal are of the order of $< 5 \text{ kV/cm}^{-1}$, following Ref. [11]. The measurements of field strengths of a few degrees presented here correspond roughly to 1 kV/cm^{-1} per degree of induced birefringence.

The extinction ratio of the detection setup becomes the main concern when aiming at detecting weak THz fields. Linear polarizers can have extinction ratios greater than 10^9 . The WP (as seen in Fig. 1) that separates the two polarizations has an extinction on the order of 10^5 . We observed that the QWP nearly does not affect extinction, but that the somehow reduced extinction ratio obtained in our setup is mainly limited by the detection crystal, which limits the extinction ratio to about 10^4 . A better ZnTe crystal with better optical quality is highly desirable for the detection of even weaker THz fields. As seen by the formula for the maximum amplification, $\Lambda = \frac{1}{4\sqrt{r}}$, a r of 10^{-5} would give amplifications of 79, while realizing $r = 10^{-6}$ would give amplifications of 250.

We applied this amplification to the technique of THz emission spectroscopy, but it could also find interest in the field of THz absorption spectroscopy (steady state or with optical excitation), where relatively thick emitting crystals are used in the generation. Thick THz crystals increase the THz maximum

field emitted while reducing its bandwidth, due to the group velocity mismatch between the THz and the input visible or near-infrared pulse. Combining our technique with the use of thinner emitting crystal could increase the bandwidth detection in THz absorption setups.

Second, in THz spectroscopy of materials in the gas phase, the rotational free induction decay (FID) usually lasts tens of picoseconds (corresponding to absorption lines of tens of GHz) and is usually orders of magnitude weaker than the input field. EOS detection of gas phase FID can be strongly amplified without distortion while the distorted input field can be ignored for analysis [12]. This could increase significantly the detectivity of such setups.

THz spectroscopy is also very popular when used with a high repetition rate system, at MHz frequencies. When EOS is used and when the gating light is in large excess to saturate the detector, our technique can be used. The THz fields at MHz repetition rate are typically weak and far from the limit where detection-induced distortions appear.

4. CONCLUSION AND OUTLOOK

We present a concise understanding of the possible amplification in the EOP of THz pulses by biasing the gating field polarization state. We deduce boundaries within which it is possible to use this amplification technique with a single measurement with negligible distortions. We demonstrate this at a comparably weak field, which is nearly impossible to detect in the standard method. With little additions to the typical detection scheme, sensitivities can strongly be enhanced in a variety of applications not limited to THz emission experiments [13,14], but also in optical pump THz probe experiments significantly reducing measurement time [15,16]. Making these changes to the typical EOP allows for the measurement of weak fields such as those generated in chemical and photophysical processes. It also helps in fiber-coupled detection schemes [17]. The requirements for this application are efficient optics in terms of polarization and the availability of excess gating photons while staying below the damage thresholds of the detection crystal and the optical detector, respectively.

APPENDIX A: MUELLER FORMALISM

By using Stokes vector and Mueller matrices, we derive Eq. (1). The two full matrices for the optical system for each of the balanced detectors consist of the following:

- a linear polarizer at 0° with extinction ratio r : $LP(0, r)$,
- the ZnTe crystal, which acts as a variable wave plate at 45° : $EO(\Gamma_{LB})$,
- a QWP at angle α : $QWP(\alpha)$,
- a linear polarizer for the two PDs of the balanced detector PD1: $LP(0, r)$ and for PD2: $LP(\frac{\pi}{2}, r)$ at 0° and 90° , both with extinction ratio r .

We use two scalars, $A1$ and $A2$, to account for the two attenuators in Fig. 1 and ensure balancing in the very end.

The input Stokes S_{in} vector corresponds to unpolarized light with an intensity depending on the first attenuator A_1 . Without

A_1 , we ensure an intensity of 1 on each PD with $\alpha = 45^\circ$, in absence of a THz field ($\Gamma_{LB} = 0$) by setting $S_{in} = \begin{pmatrix} 4 \\ 0 \\ 0 \\ 0 \end{pmatrix}$.

The intensity on PD1 is the first element ($A1$) of the Stokes vector corresponding to the intensity at PD1 without attenuator,

$$I_{PD1}(\Gamma_{LB}, \alpha, r) = \left(LP\left(\frac{\pi}{2}, r\right) \cdot QWP(\alpha) \cdot EO(\Gamma_{LB}) \cdot LP(0, r) \cdot S_{in} \right), \quad (A1)$$

whereas the intensity without attenuator on PD2 is

$$I_{PD2}(\Gamma_{LB}, \alpha, r) = \left(LP(0, r) \cdot QWP(\alpha) \cdot EO(\Gamma_{LB}) \cdot LP(0, r) \cdot S_{in} \right). \quad (A2)$$

The difference between these two intensities is the last polarizer, which is applied in the formalism and corresponds to the two outputs of the WP. Finally, considering the attenuation, the detected signal calculated with Eq. (1) is the difference between these intensities,

$$S_r(\alpha, \Gamma_{LB}) = A1 \times (I_{PD1} - A2 \times I_{PD2}),$$

with $A1 = \frac{1}{I_{PD1}(0, \alpha, r)}$ and $A2 = \frac{I_{PD1}(0, \alpha, r)}{I_{PD2}(0, \alpha, r)}$ to ensure balancing of the PDs at ($\Gamma_{LB} = 0$).

The Mueller matrices used above are

$$LP(0, r) = \frac{1}{2} \begin{pmatrix} (1+r) & (1-r) & 0 & 0 \\ (1-r) & (1+r) & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix},$$

$$LP\left(\frac{\pi}{2}, r\right) = \frac{1}{2} \begin{pmatrix} (1+r) & (r-1) & 0 & 0 \\ (r-1) & (1+r) & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix},$$

$$QWP(\alpha) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos(2\alpha)^2 & \frac{1}{2} \sin(4\alpha) & -\sin(2\alpha) \\ 0 & \frac{1}{2} \sin(4\alpha) & \sin(2\alpha)^2 & \cos(2\alpha) \\ 0 & \sin(2\alpha) & -\cos(2\alpha) & 0 \end{pmatrix},$$

and

$$EO(\Gamma_{LB}) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos(\Gamma_{LB}) & 0 & \sin(\Gamma_{LB}) \\ 0 & 0 & 1 & 0 \\ 0 & -\sin(\Gamma_{LB}) & 0 & \cos(\Gamma_{LB}) \end{pmatrix}.$$

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