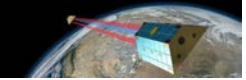


# Precise Orbit Determination

**Adrian Jäggi**

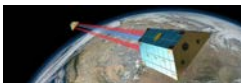
Astronomical Institute  
University of Bern



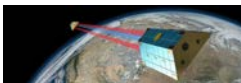


# Lecture Contents

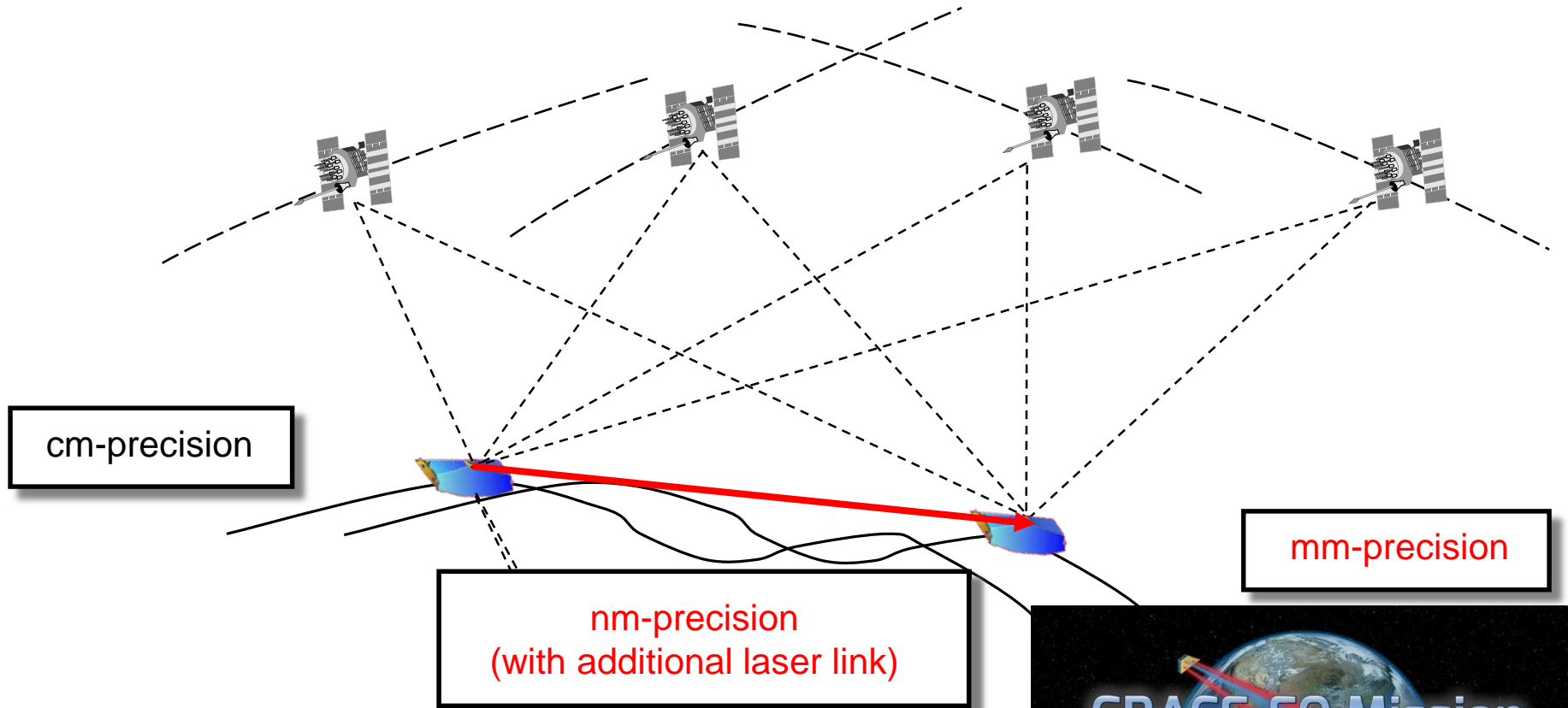
1. Introduction
2. Global Positioning System
3. Different Orbit Representations
4. Principles of Orbit Determination
5. GPS-based LEO POD
6. Orbit Validation



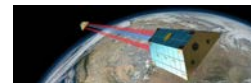
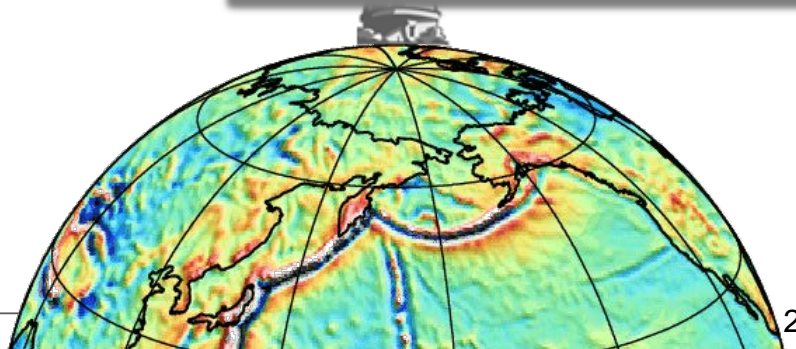
# A multitude of Earth Observation Satellites



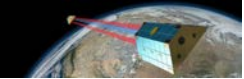
# Precise Tracking Data in Near Earth Space



**GRACE-FO Mission**  
LISA Technology  
Sheds Light on Climate Change

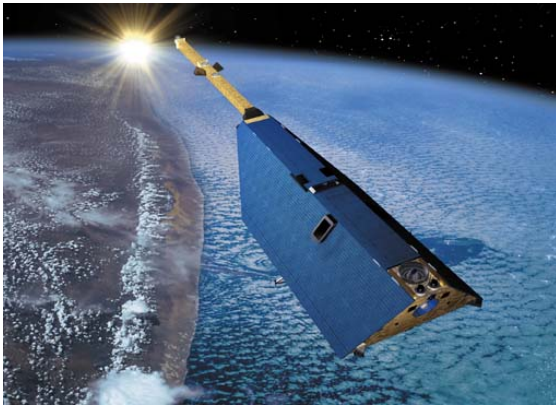






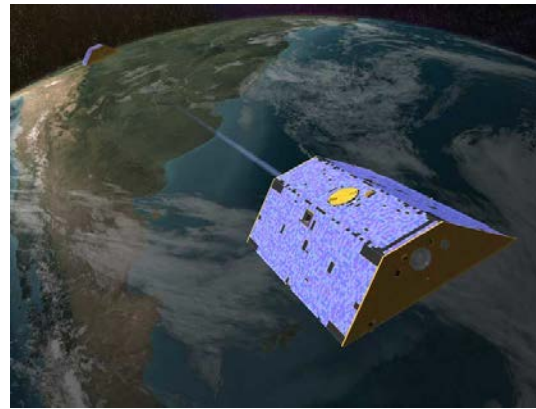
# Low Earth Orbiters (LEOs)

## CHAMP



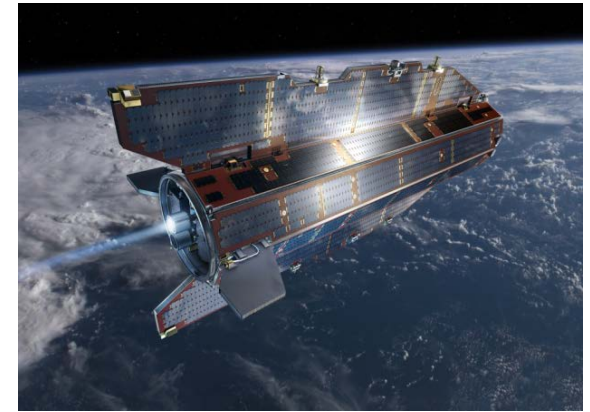
**CH**allenging  
**Mini**satellite **P**ayload

## GRACE



**G**ravity **R**ecovery **A**nd  
**C**limate **E**xperiment

## GOCE



**G**ravity and  
steady-state **O**cean  
**C**irculation **E**xplorer

Of course, there are many more missions equipped with GPS receivers

### Jason



### Jason-2



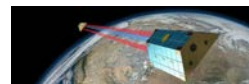
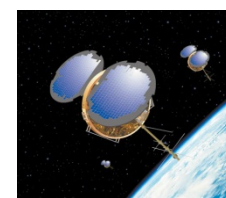
### MetOp-A

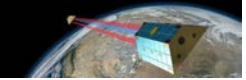


### Icesat



### COSMIC





# LEO Constellations

## TanDEM-X



## Swarm



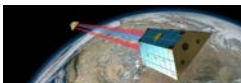
## Sentinel



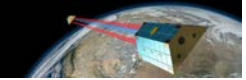
## GRACE-FO



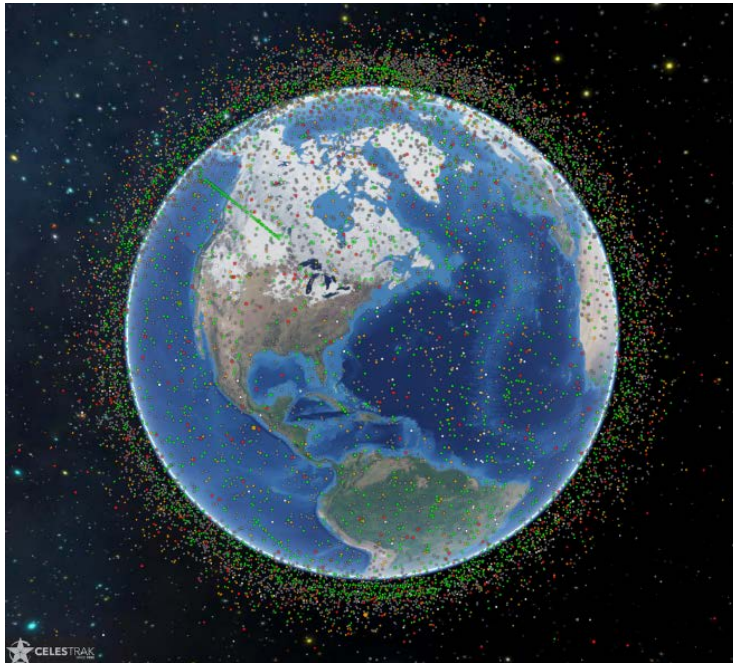
## COSMIC-2







# LEO CubeSats



A multitude of **CubeSats** already exist or are planned for the future

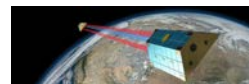
## Spire

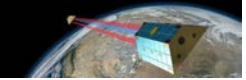
e.g.



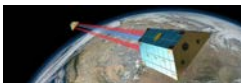
~ 100 Spire satellites are already in different orbits, offering dual-frequency high-quality GPS data

**Many more will follow ...**





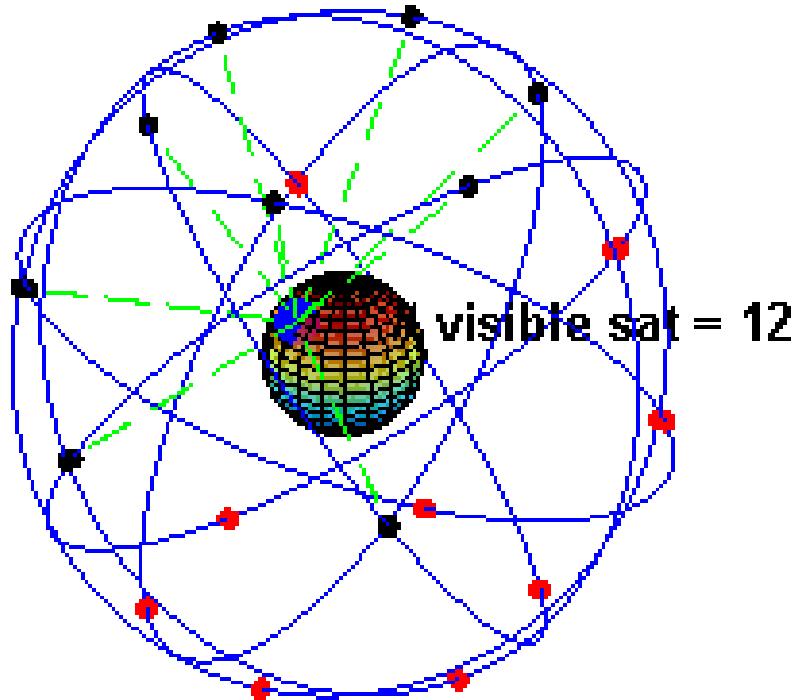
# Global Positioning System



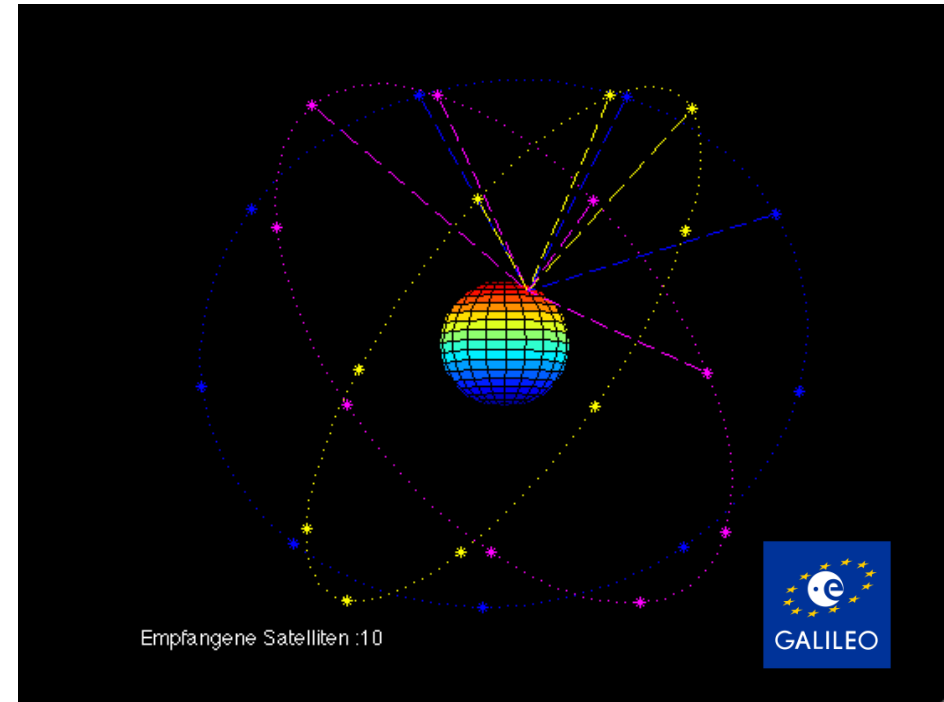


# Introduction to GPS

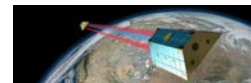
GPS

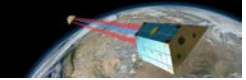


Galileo



Other **Global Navigation Satellite Systems** (GNSS) are also available (GLONASS, Galileo, Beidou), but for a long time no multi-GNSS spaceborne receivers were in orbit. This changed with the launches of Fengyun-3, COSMIC-2, Sentinel-6.



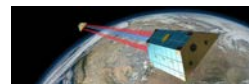


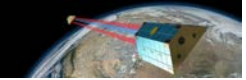
# Introduction to GPS

**GPS: Global Positioning System**

**Characteristics:**

- Satellite system for (real-time) **Positioning** and **Navigation**
- **Global** (everywhere on Earth, up to altitudes of 5000km) and **at any time**
- **Unlimited** number of users
- **Weather-independent** (radio signals are passing through the atmosphere)
- **3-dimensional position, velocity** and **time** information





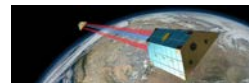
# GPS Segments

The GPS consists of **3 main segments**:

- **Space Segment:** the satellites and the constellation of satellites
- **Control Segment:** the ground stations, infrastructure and software for operation and monitoring of the GPS
- **User Segment:** all GPS receivers worldwide and the corresponding processing software


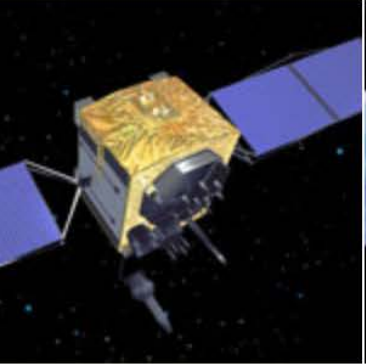
We should add an important **4th segment**:

- **Ground Segment:** all civilian permanent networks of reference sites and the international/regional/local services delivering products for the users

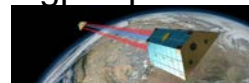


# Space Segment

- The space segment nominally consists of **24 satellites**, presently: 31 active GPS satellites
- Constellation design: at least **4 satellites** in view from **any location** on the Earth at **any time**

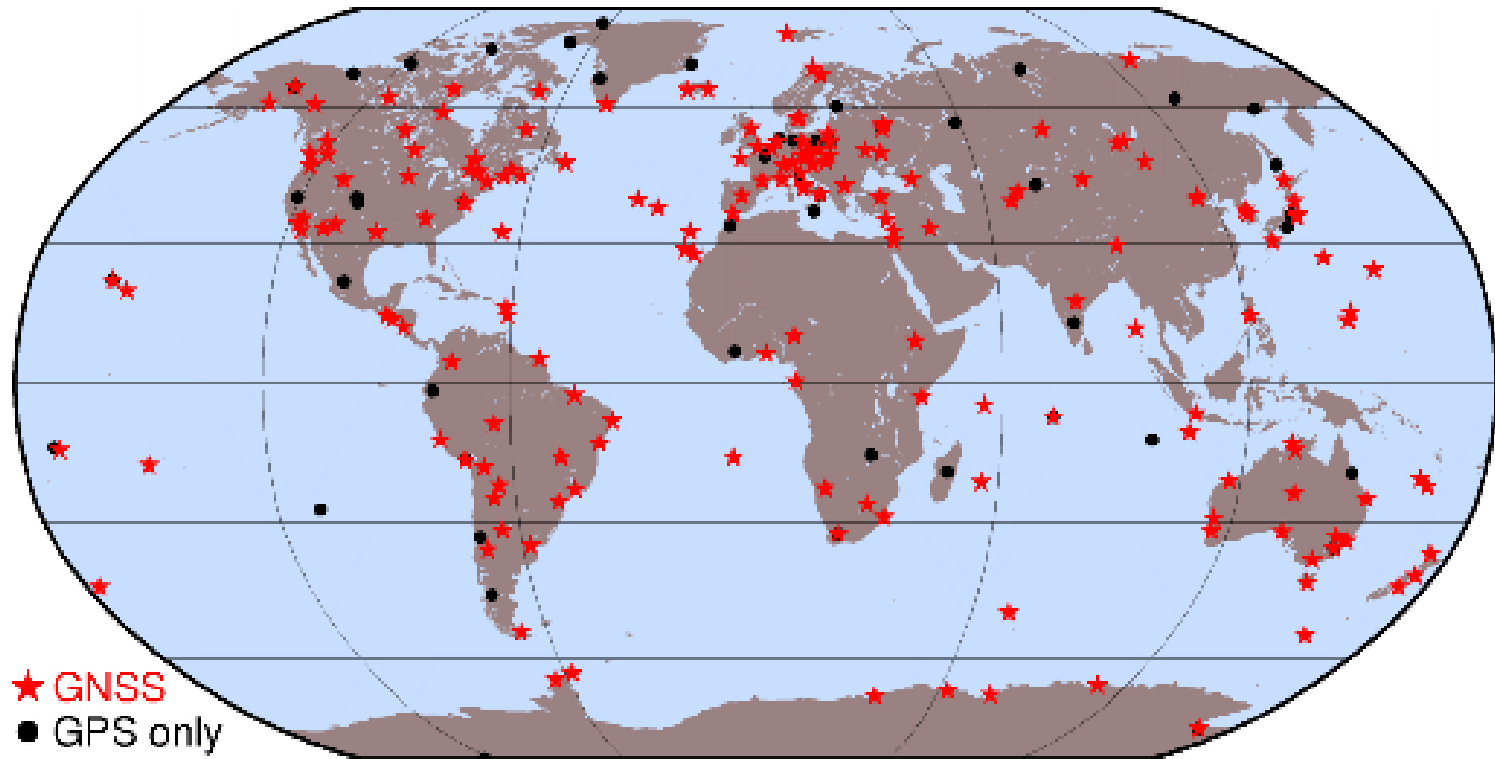
LEGACY SATELLITES		MODERNIZED SATELLITES		
				
BLOCK IIA	BLOCK IIR	BLOCK IIR-M	BLOCK IIF	GPS III/IIIF
0 operational	7 operational	7 operational	12 operational	5 operational

<https://www.gps.gov/systems/gps/space/>

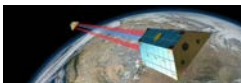




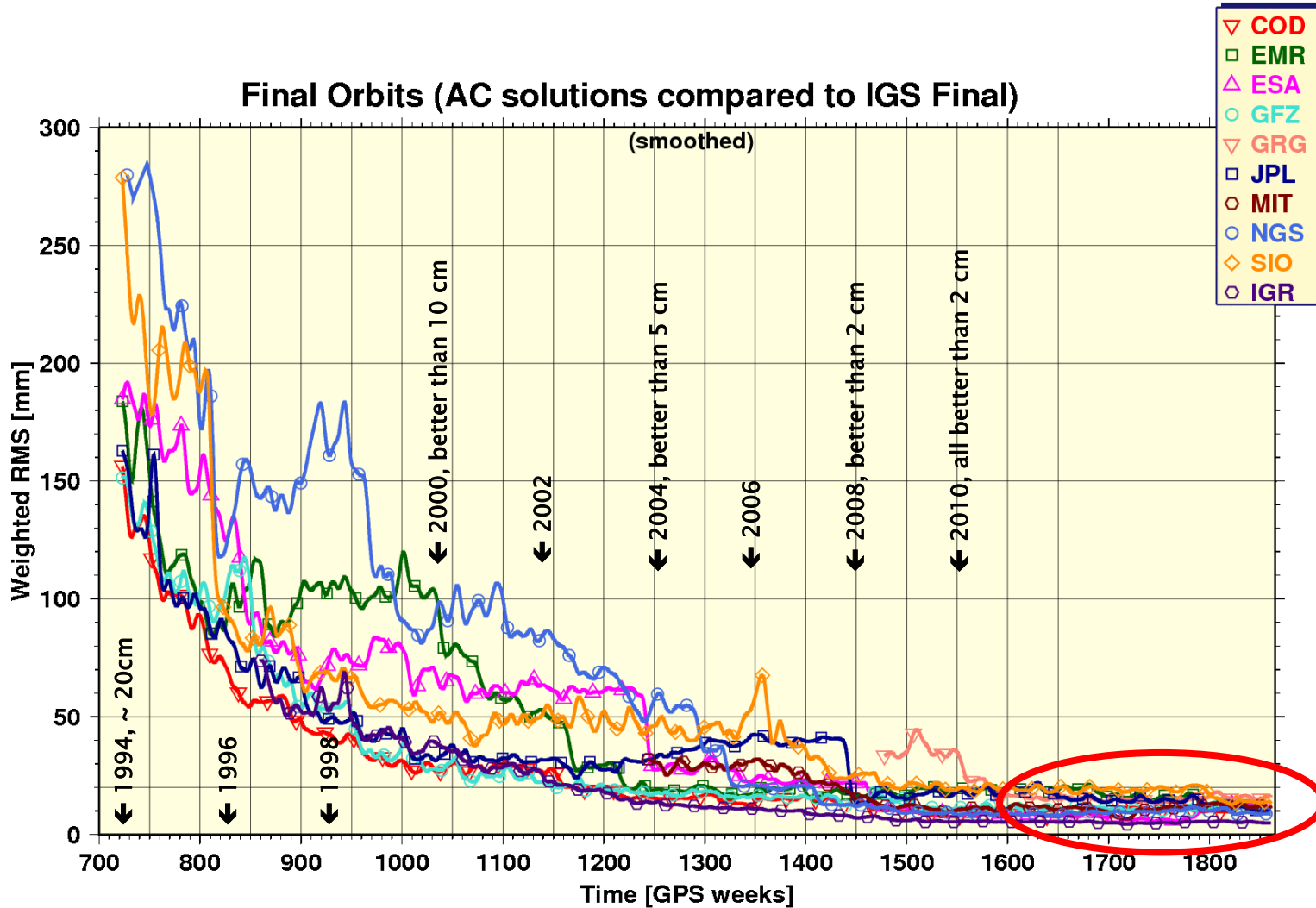
# Global Network of the IGS



IGS stations used for computation of final orbits at CODE (Dach et al., 2009)

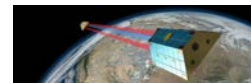


# Performance of IGS Final Orbits

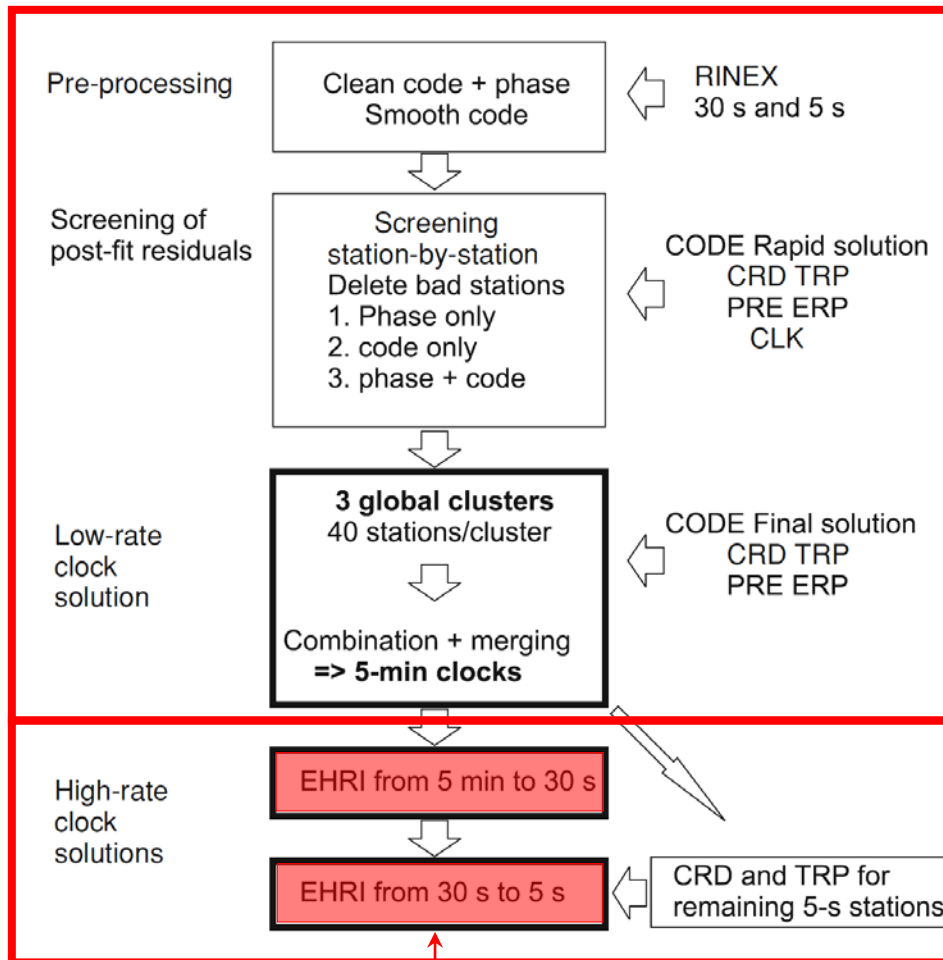


<http://acc.igs.org/>

NOAA NGS, 19.09.2015 19:19 (GMT)



# Computation of High-Rate Clocks

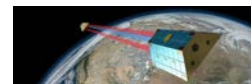


(Bock et al., 2009)

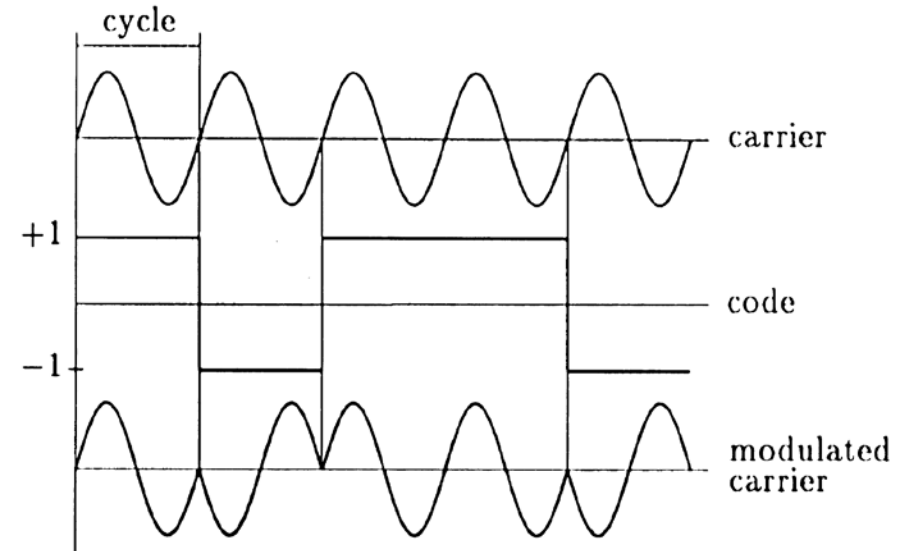
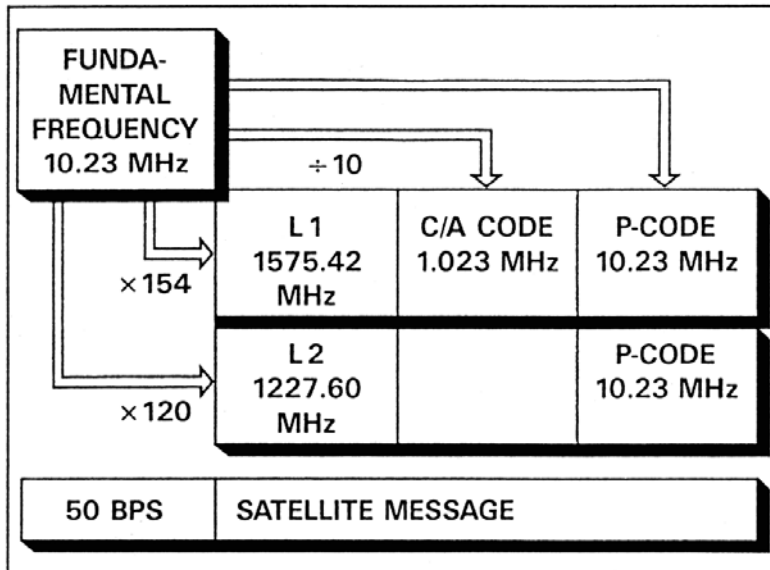
The final clock product with 5 min sampling is based on undifferenced GPS data of typically 120 stations of the IGS network

The IGS 1 Hz network is finally used for clock densification to 5 sec

The 5 sec clocks are interpolated to 1 sec as needed for 1 Hz kinematic LEO POD



# GPS Signals



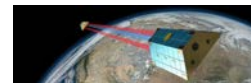
Bits encoded on carrier by phase modulation:

Signals driven by an **atomic clock**

Two **carrier signals** (sine waves):

- **L1**:  $f = 1575.43$  MHz,  $\lambda = 19$  cm
- **L2**:  $f = 1227.60$  MHz,  $\lambda = 24$  cm

- **C/A-code** (Clear Access / Coarse Acquisition)
- **P-code** (Protected / Precise)
- **Broadcast/Navigation Message**

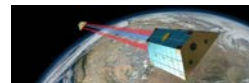
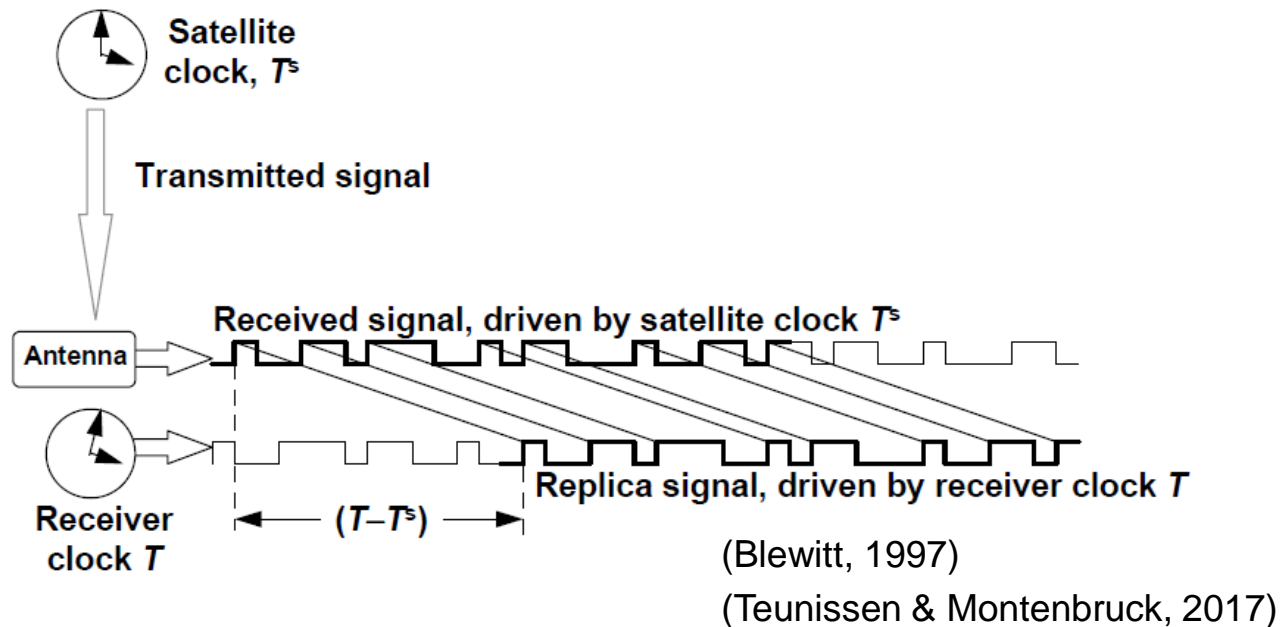


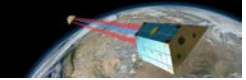


# Pseudorange / Code Measurements

Code Observations  $P_i^k$  are defined as:

$$P_i^k \doteq c (T_i - T^k)$$





# Code Observation Equation

$$P_i^k = \rho_i^k - c \cdot \Delta t^k + c \cdot \Delta t_i$$

$t_i, t^k$  GPS time of reception and emission

$\Delta t^k$  Satellite clock offset  $T^k - t^k$

$\Delta t_i$  Receiver clock offset  $T_i - t_i$

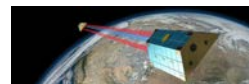
$\rho_i^k$  Distance between receiver and satellite  $c (t_i - t^k)$

**Known** from ACs or IGS:

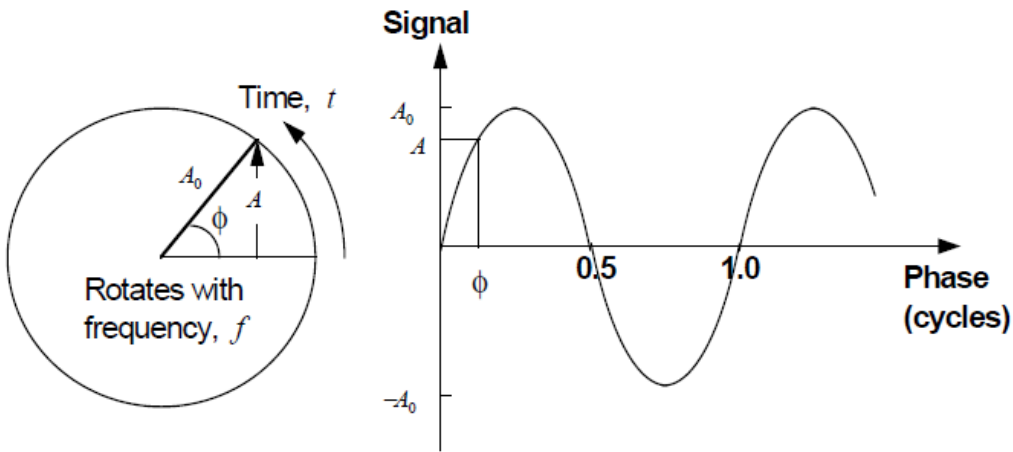
- satellite positions  $(x^{kj}, y^{kj}, z^{kj})$
- satellite clock offsets  $\Delta t^{kj}$

**4 unknown parameters:**

- receiver position  $(x_i, y_i, z_i)$
- receiver clock offset  $\Delta t_i$



# Carrier Phase Measurements (1)



Phase  $\phi$  (in cycles) increases linearly with time  $t$ :

$$\phi = f \cdot t$$

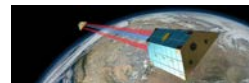
where  $f$  is the frequency

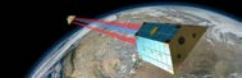
The **satellite** generates with its clock the phase signal  $\phi^k$ . At emission time  $T^k$  (in satellite clock time) we have

$$\phi^k = f \cdot T^k$$

The same phase signal, e.g., a wave crest, propagates from the satellite to the receiver, but the receiver measures only the fractional part of the phase and does not know the **integer number of cycles**  $N_i^k$  (phase ambiguity):

$$\phi_i^k = \phi^k - N_i^k = f \cdot T^k - N_i^k$$





## Carrier Phase Measurements (2)

The **receiver** generates with its clock a **reference phase**. At time of reception  $T_i$  of the satellite phase  $\phi_i^k$  (in receiver clock time) we have:

$$\phi_i = f \cdot T_i$$

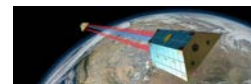
The actual **phase measurement** is the difference between receiver reference phase  $\phi_i$  and satellite phase  $\phi_i^k$ :

$$\psi_i^k = \phi_i - \phi_i^k = f \cdot T_i - (f \cdot T^k - N_i^k) = f \cdot (T_i - T^k) + N_i^k$$

Multiplication with the wavelength  $\lambda = c/f$  leads to the **phase observation equation** in meters:

$$\begin{aligned} L_i^k &= \lambda \cdot \psi_i^k = c \cdot (T_i - T^k) + \lambda \cdot N_i^k \\ &= \rho_i^k - c \cdot \Delta t^k + c \cdot \Delta t_i + \lambda \cdot N_i^k \end{aligned}$$

Difference to the pseudorange observation: **integer ambiguity term**  $N_i^k$



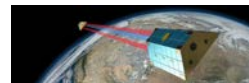
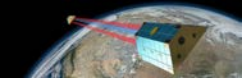


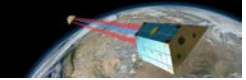
## Detailed Observation Equation

$$L_i^k = \rho_i^k - c \cdot \Delta t^k + c \cdot \Delta t_i + \cancel{I_i^k} + \cancel{I_i^k} + \lambda \cdot N_i^k + \Delta_{rel} - c \cdot b^k + c \cdot b_i + m_i^k + \epsilon_i^k$$

$\rho_i^k$	Distance between satellite and receiver	←	Satellite positions and clocks
$\Delta t^k$	Satellite clock offset wrt GPS time	←	are known from the IGS
$\Delta t_i$	Receiver clock offset wrt GPS time		
<del><math>T_i^k</math></del>	<del>Tropospheric delay</del>	←	Not existent for LEOs
<del><math>I_i^k</math></del>	<del>Ionospheric delay</del>	←	Cancels out (first order only) when forming the ionosphere-free linear combination:
$N_i^k$	Phase ambiguity		
$\Delta_{rel}$	Relativistic corrections		
$b^k$	Delays in satellite (cables, electronics)		
$b_i$	Delays in receiver and antenna		
$m_i^k$	Multipath, scattering, bending effects		
$\epsilon_i^k$	Measurement error		

$$L_c = \frac{f_1^2}{f_1^2 - f_2^2} L_1 - \frac{f_2^2}{f_1^2 - f_2^2} L_2$$





## Geometric Distance

**Geometric distance**  $\rho_{leo}^k$  is given by:

$$\rho_{leo}^k = |\mathbf{r}_{leo}(t_{leo}) - \mathbf{r}^k(t_{leo} - \tau_{leo}^k)|$$

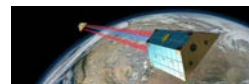
$\mathbf{r}_{leo}$  Inertial position of LEO antenna phase center at reception time

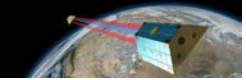
$\mathbf{r}^k$  Inertial position of GPS antenna phase center of satellite  $k$  at emission time

$\tau_{leo}^k$  Signal traveling time between the two phase center positions

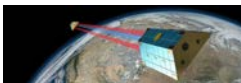
Different ways to represent  $\mathbf{r}_{leo}$ :

- **Kinematic** orbit representation
- **Dynamic** or **reduced-dynamic** orbit representation





# Different Orbit Representations



# Kinematic Orbit Representation (1)

**Satellite position**  $\mathbf{r}_{leo}(t_{leo})$  (in inertial frame) is given by:

$$\mathbf{r}_{leo}(t_{leo}) = \mathbf{R}(t_{leo}) \cdot (\mathbf{r}_{leo,e,0}(t_{leo}) + \delta\mathbf{r}_{leo,e,ant}(t_{leo}))$$

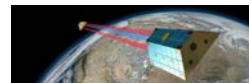
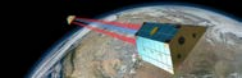
$\mathbf{R}$  Transformation matrix from Earth-fixed to inertial frame

$\mathbf{r}_{leo,e,0}$  LEO center of mass position in Earth-fixed frame

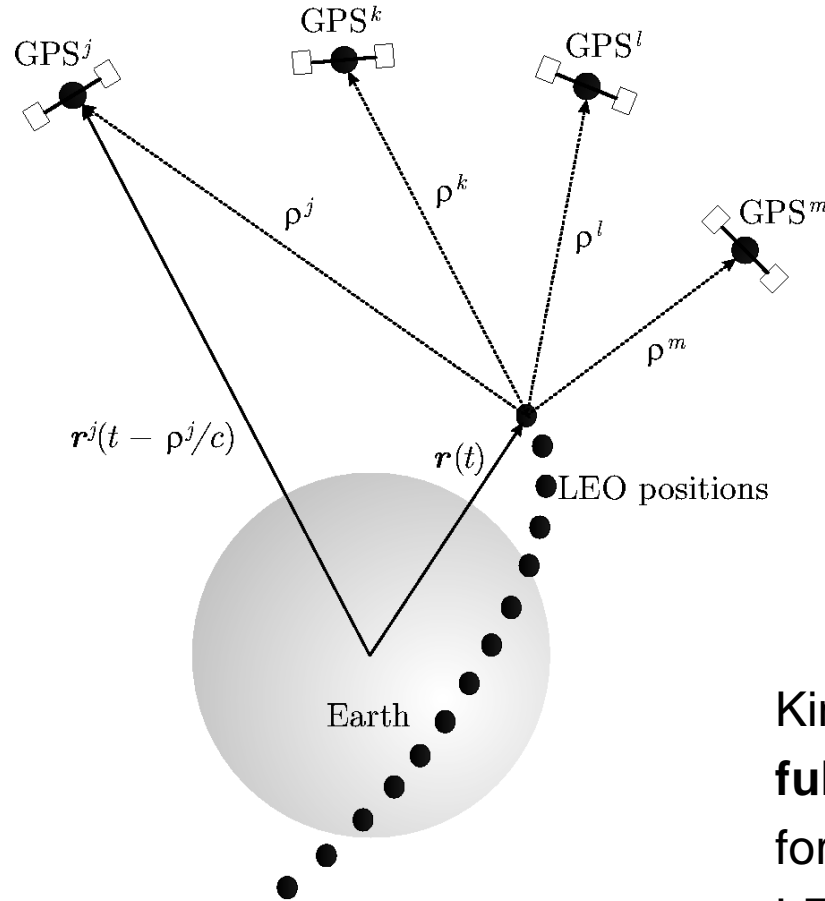
$\delta\mathbf{r}_{leo,e,ant}$  LEO antenna phase center offset in Earth-fixed frame

**Kinematic positions**  $\mathbf{r}_{leo,e,0}$  are estimated for each **measurement epoch**:

- Measurement epochs **need not** to be identical with nominal epochs
- Positions are **independent** of models describing the LEO dynamics. Velocities and accelerations cannot be provided in a "strict" sense.

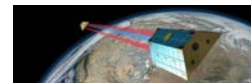


## Kinematic Orbit Representation (2)



A kinematic orbit is an ephemeris at **discrete** measurement epochs

Kinematic positions are **fully independent** on the force models used for LEO orbit determination (Svehla and Rothacher, 2004)





# Kinematic Orbit Representation (3)

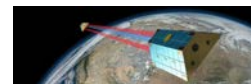
Measurement epochs  
(in GPS time)

Positions (km)  
(Earth-fixed)

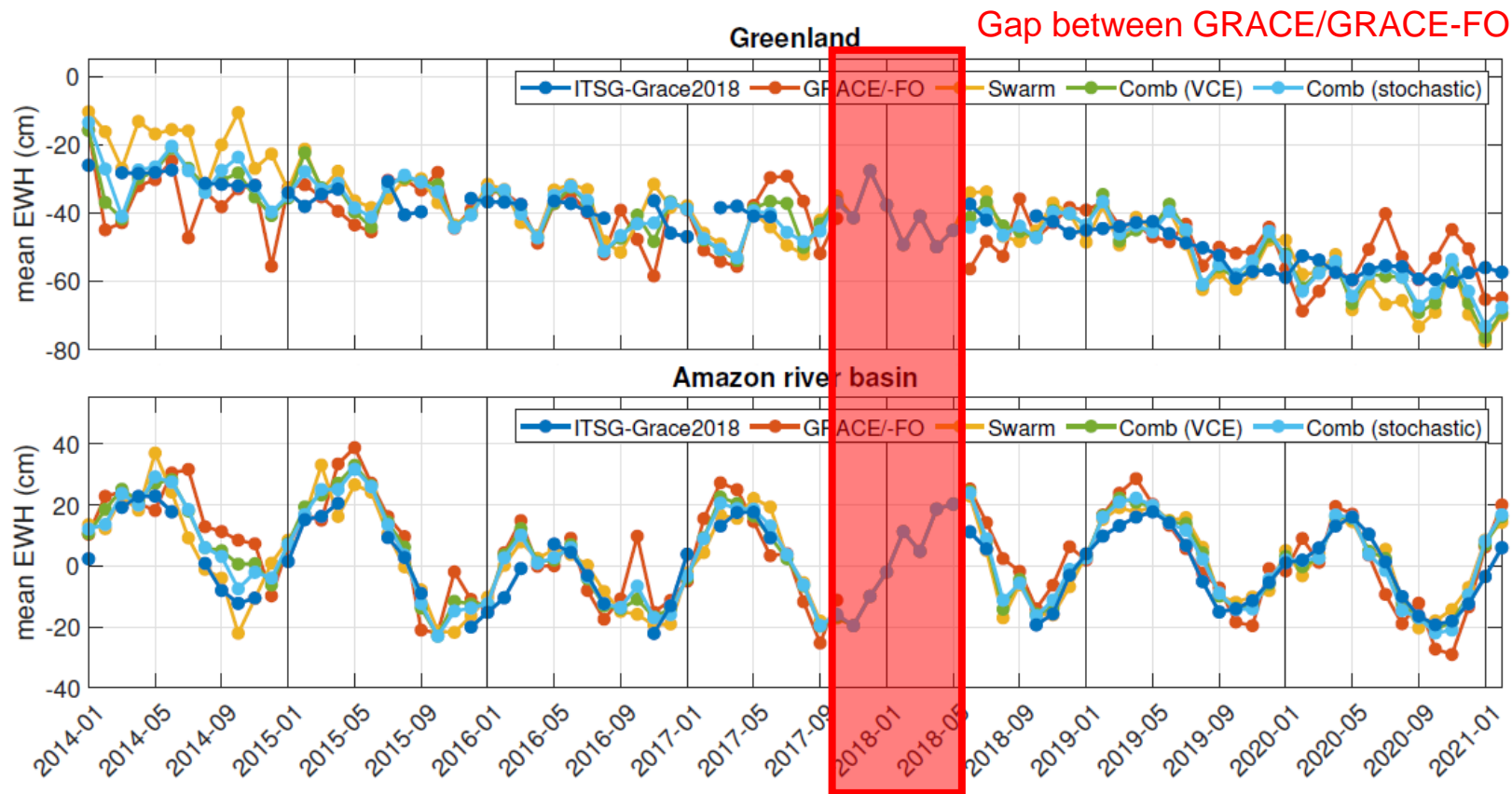
* 2009 11 2	0	0	0.80678020		
PL15	-390.612059	6623.987679	73.104149	193219.797196	
* 2009 11 2	0	0	1.80678020		
PL15	-389.240315	6624.166512	65.402457	193219.799413	
* 2009 11 2	0	0	2.80678020		
PL15	-387.868014	6624.336133	57.700679	193219.801634	
* 2009 11 2	0	0	3.80678020		
PL15	-386.495163	6624.496541	49.998817	193219.803855	
* 2009 11 2	0	0	4.80678019		
PL15	-385.121760	6624.647724	42.296889	193219.806059	
* 2009 11 2	0	0	5.80678019		
PL15	-383.747819	6624.789703	34.594896	193219.808280	
* 2009 11 2	0	0	6.80678019		
PL15	-382.373332	6624.922464	26.892861	193219.810495	
* 2009 11 2	0	0	7.80678019		
PL15	-380.998306	6625.046003	19.190792	193219.812692	
* 2009 11 2	0	0	8.80678019		
PL15	-379.622745	6625.160329	11.488692	193219.814899	
* 2009 11 2	0	0	9.80678018		
PL15	-378.246651	6625.265448	3.786580	193219.817123	

Clock correction to nominal epoch ( $\mu$ s), e.g., to epoch 00:00:03

Excerpt of kinematic GOCE positions at begin of 2 Nov, 2009

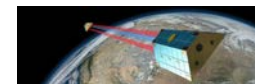


# Recovery of Large-Scale Time-VARIABLE Gravity Signals

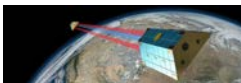
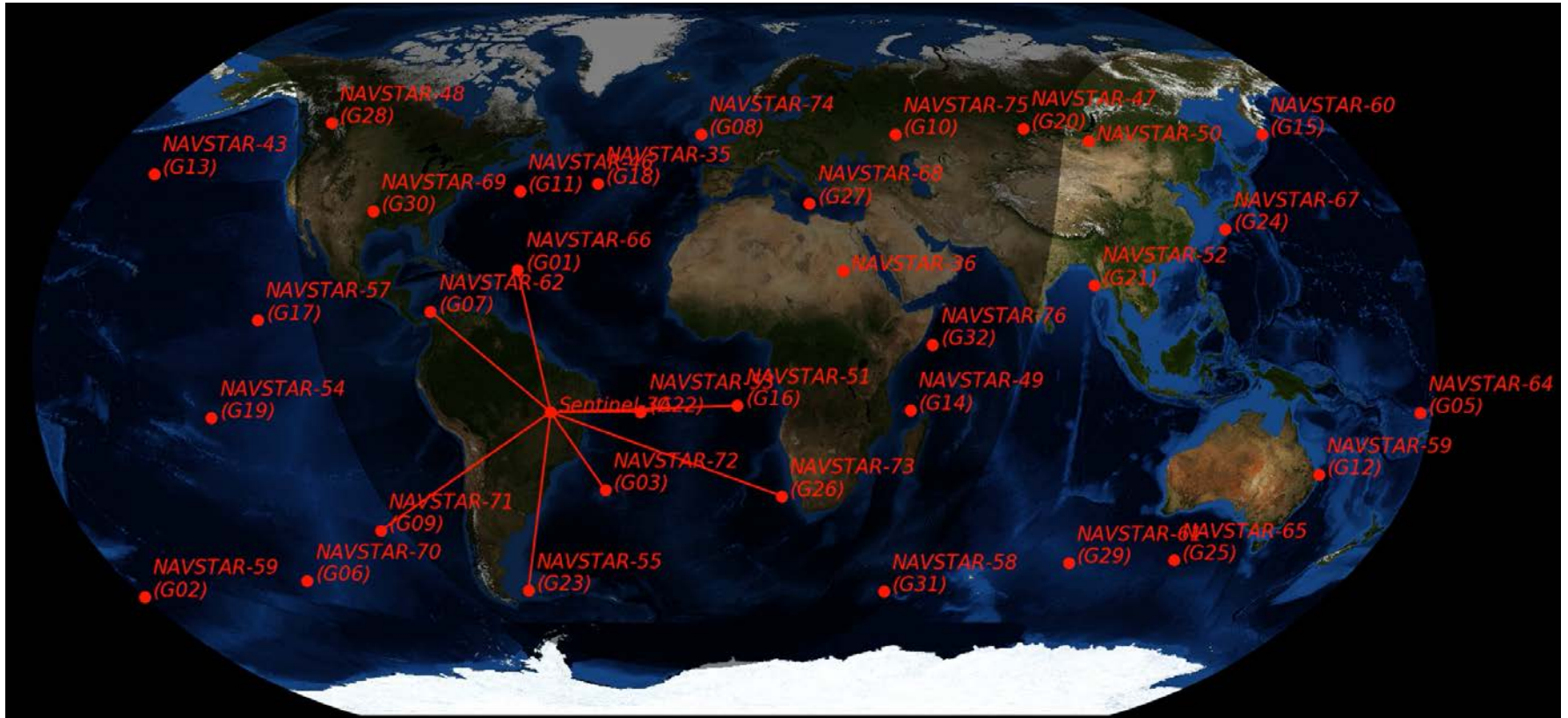


Kinematic positions allow it to recover the **long wavelength** part of the Earth's time-variable gravity field. The scatter of monthly gravity field solutions is larger than from dedicated GRACE/GRACE-FO data, but trends and annual signals may be derived remarkably well.

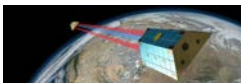
(Grombein et al., 2022)



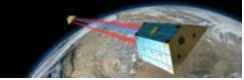
# Example: Sentinel-3A GPS Tracking



# Example: Sentinel-6A multi-GNSS Tracking







# Simulated Data Set (1)

The simplified **Code Observation Equation** of the simulation reads as

$$\bar{\rho}_{k,i} = \sqrt{(x_{k,i} - x_{leo,i})^2 + (y_{k,i} - y_{leo,i})^2 + (z_{k,i} - z_{leo,i})^2} + c_{leo,i} \quad , \quad k = 1, \dots, n_{sat}$$

with

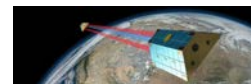
- $(x_{k,i}, y_{k,i}, z_{k,i})$       Known inertial position of GPS satellite  $k$  at epoch  $i$
- $(x_{leo,i}, y_{leo,i}, z_{leo,i})$       Unknown inertial LEO position at epoch  $i$
- $c_{leo,i}$       Unknown clock correction of LEO receiver at epoch  $i$

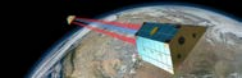
The simplified **Phase Observation Equation** of the simulation reads as

$$\bar{\lambda}_{k,i} = \sqrt{(x_{k,i} - x_{leo,i})^2 + (y_{k,i} - y_{leo,i})^2 + (z_{k,i} - z_{leo,i})^2} + c_{leo,i} + b_k \quad , \quad k = 1, \dots, n_{sat}$$

with the additional parameter

- $b_k$       Unknown constant phase bias to the satellite  $k$





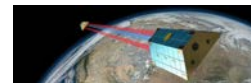
# Simulated Data Set (2)

Tabelle 1: Code observations  $\rho'_{k,i}$  contained in the file OBS\_CODE.txt. They are stored in template.m in the array obs\_code(kepo, isat). In analogy the phase observations are contained in the file OBS\_PHASE.txt and are stored in the array obs\_phase(kepo, isat). The observation times are also computed in the source code by  $t = 10*(kepo-1)$ ,  $kepo = 1, \dots, nepo$ .

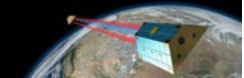
Time $t_i$ (sec)	Sat. 1 (m)	Sat. 2 (m)	Sat. 3 (m)	Sat. 4 (m)	Sat. 5 (m)	etc.
0.0000	19447557.3266	0.0000	0.0000	0.0000	21654965.4010	...
10.0000	19446601.2222	0.0000	0.0000	0.0000	21678374.5718	...
20.0000	19446475.1291	0.0000	0.0000	0.0000	21702249.7103	...
...	...	...	...	...	....	...

Tabelle 2: Positions  $(x_{k,i}, y_{k,i}, z_{k,i})$  of the GPS satellites contained in the file GPS.txt. They are stored in the array r\_gps(coord, iepo, ksatsat).

$x$ GPS 1 (m)	$y$ GPS 1 (m)	$z$ GPS 1 (m)	$x$ GPS 2	$y$ GPS 2	$z$ GPS 2	etc.
26235000.0000	0.0000	0.0000	-789889.9012	-15196033.7951	-21702185.3748	...
26234971.0435	22468.8157	32088.7944	-751120.9506	-15196458.0361	-21702791.2539	...
26234884.1742	44937.5818	64177.5179	-712350.3922	-15196849.7458	-21703350.6732	...
...	...	...	...	...	....	...







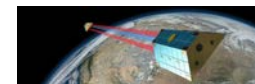
# Simulated Data Set (3)

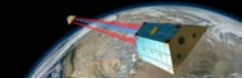
Tabelle 3: Positions  $(x_{leo,i}, y_{leo,i}, z_{leo,i})$  and velocities of the LEO satellite contained in the file `LEO.txt`. The positions are stored in the array `r_leo(coord,iepo)`. The velocities are not needed for this project unless the orbit differences shall be plotted in a radial, along-track, cross-track frame instead of the inertial  $x,y,z$  frame.

$x$ LEO (m)	$y$ LEO (m)	$z$ LEO (m)	$v_x$ LEO (m/s)	$v_y$ LEO (m/s)	$v_z$ LEO (m/s)
6824717.7284	1203381.8712	0.0000	-0.0000	0.0000	7621.8949
6824309.0437	1203309.8091	76217.4278	-81.7361	-14.4123	7621.4385
6823083.0403	1203093.6316	152425.7274	-163.4621	-28.8228	7620.0693
...	...	...	...	...	...

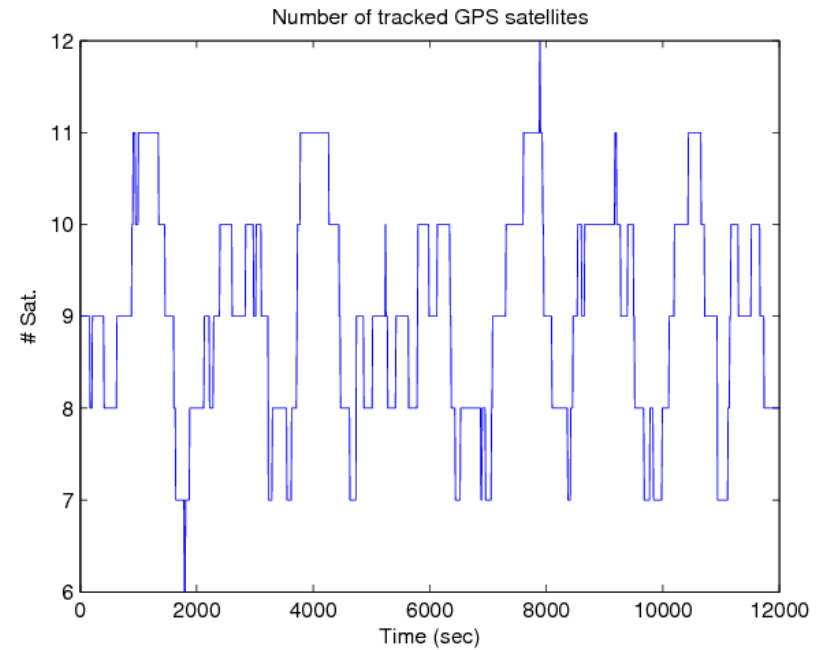
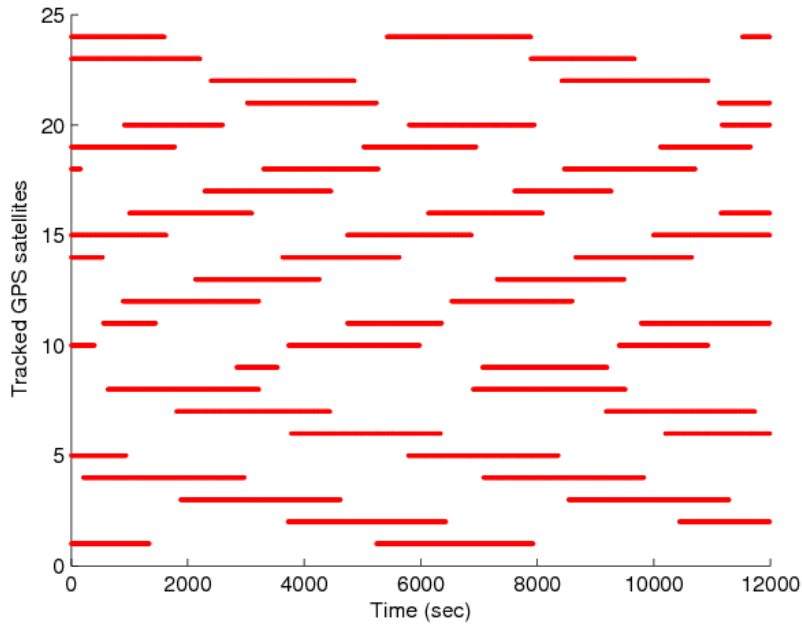
Tabelle 4: Biases  $b_k$  of the phase observations to satellite  $k$  contained in the file `BIASES.txt`. They are stored in the array `true_bias(ksat)`.

$b_1$ (m)	$b_2$ (m)	$b_3$ (m)	$b_4$ (m)	etc.
2319	-1149	2918	-10144	...

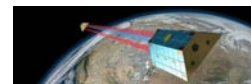




# Simulated Data Set (4)



Tracking scenario of the simulated data set (left). Up to 12 GPS satellites are at maximum simultaneously visible from the LEO satellite (right). The viewing geometry is continuously changing due to the orbital motion of all satellites.



# Dynamic Orbit Representation (1)

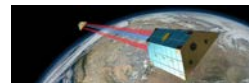
**Satellite position**  $\mathbf{r}_{leo}(t_{leo})$  (in inertial frame) is given by:

$$\mathbf{r}_{leo}(t_{leo}) = \mathbf{r}_{leo,0}(t_{leo}; a, e, i, \Omega, \omega, u_0; Q_1, \dots, Q_d) + \delta\mathbf{r}_{leo,ant}(t_{leo})$$

$\mathbf{r}_{leo,0}$	LEO center of mass position
$\delta\mathbf{r}_{leo,ant}$	LEO antenna phase center offset
$a, e, i, \Omega, \omega, u_0$	LEO initial osculating orbital elements
$Q_1, \dots, Q_d$	LEO dynamical parameters

**Satellite trajectory**  $\mathbf{r}_{leo,0}$  is a particular solution of an **equation of motion**

- One set of **initial conditions** (orbital elements) is estimated per arc.  
Dynamical parameters of the force model may be estimated on request.



## Dynamic Orbit Representation (2)

**Equation of motion** (in inertial frame) is given by:

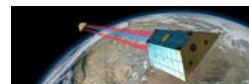
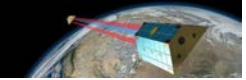
$$\ddot{\mathbf{r}} = -GM \frac{\mathbf{r}}{r^3} + \mathbf{f}_1(t, \mathbf{r}, \dot{\mathbf{r}}, Q_1, \dots, Q_d)$$

with initial conditions

$$\mathbf{r}(t_0) = \mathbf{r}(a, e, i, \Omega, \omega, u_0; t_0)$$

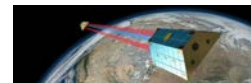
$$\dot{\mathbf{r}}(t_0) = \dot{\mathbf{r}}(a, e, i, \Omega, \omega, u_0; t_0)$$

The **acceleration**  $\mathbf{f}_1$  consists of **gravitational** and **non-gravitational** perturbations taken into account to model the satellite trajectory. Unknown parameters  $Q_1, \dots, Q_d$  of force models may appear in the equation of motion together with deterministic (known) accelerations given by analytical models.

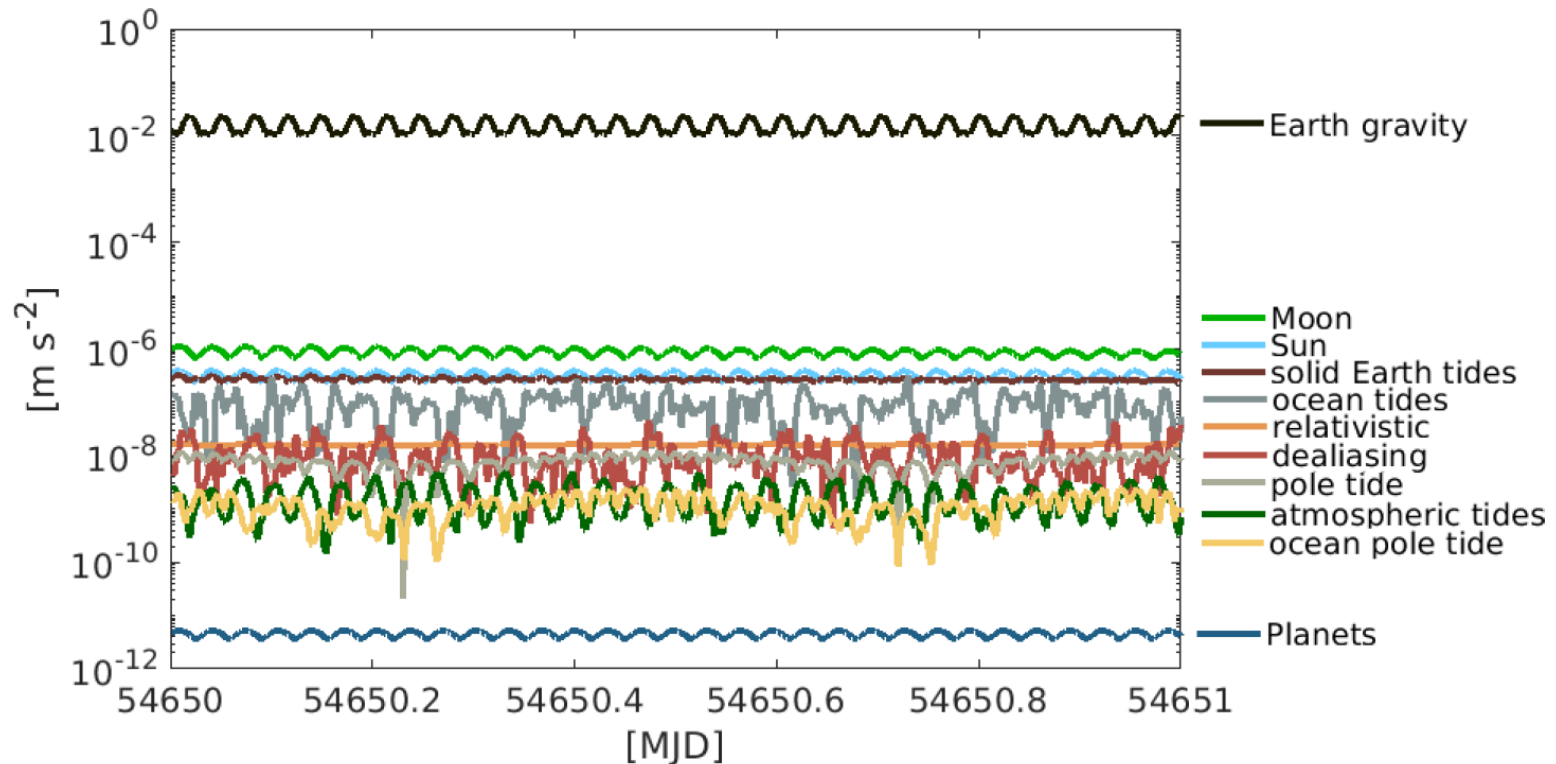


# Perturbing Accelerations of a LEO Satellite

Force	Acceleration (m/s <sup>2</sup> )
Central term of Earth's gravity field	8.42
Oblateness of Earth's gravity field	0.015
Atmospheric drag	0.00000079
Higher order terms of Earth's gravity field	0.00025
Attraction from the Moon	0.0000054
Attraction from the Sun	0.0000005
Direct solar radiation pressure	0.000000097
...	...

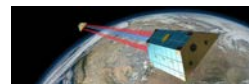


# Perturbing Accelerations of a LEO Satellite



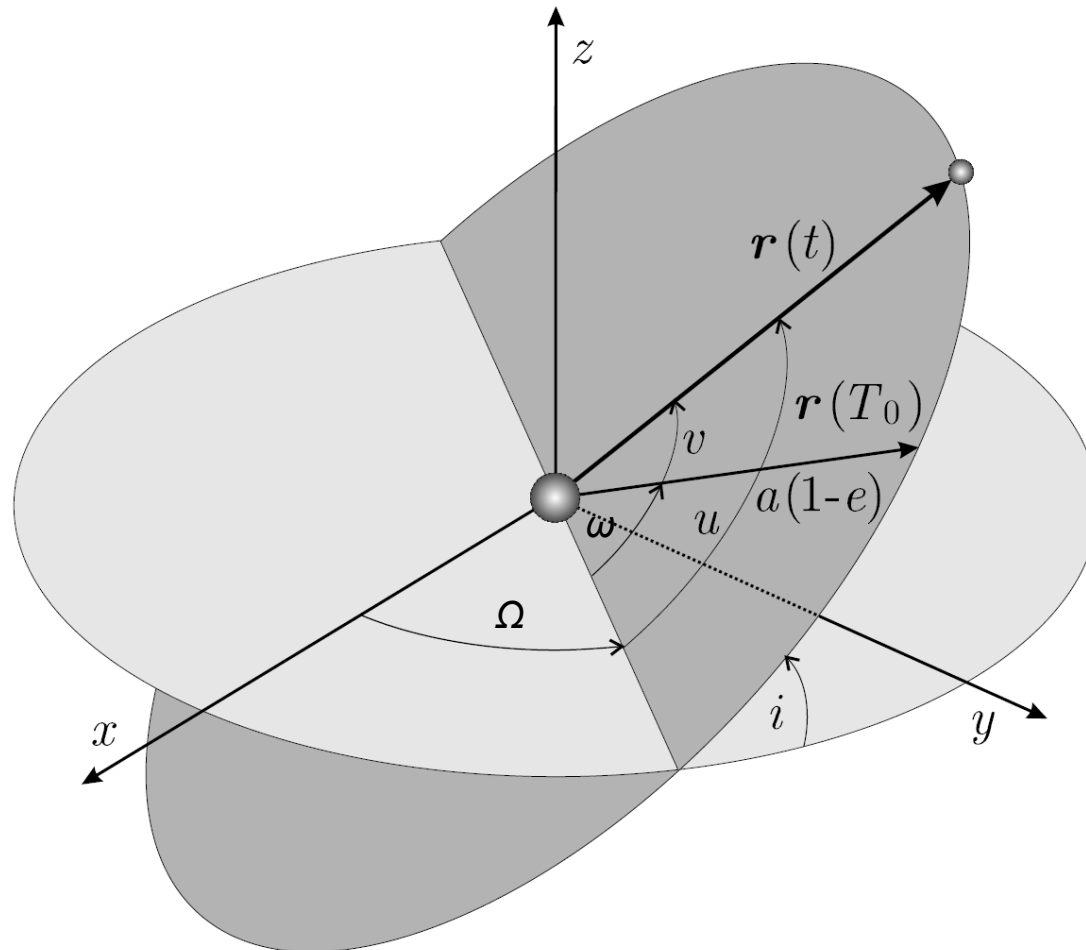
Norm of the **COST-G benchmark accelerations** along a GRACE satellite orbit. The benchmark data set may be used as a reference data set and provides the opportunity to test the implementation of corresponding background models.

(Mayer-Gürr and Kvas, 2019; Lasser et al., 2020)

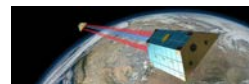




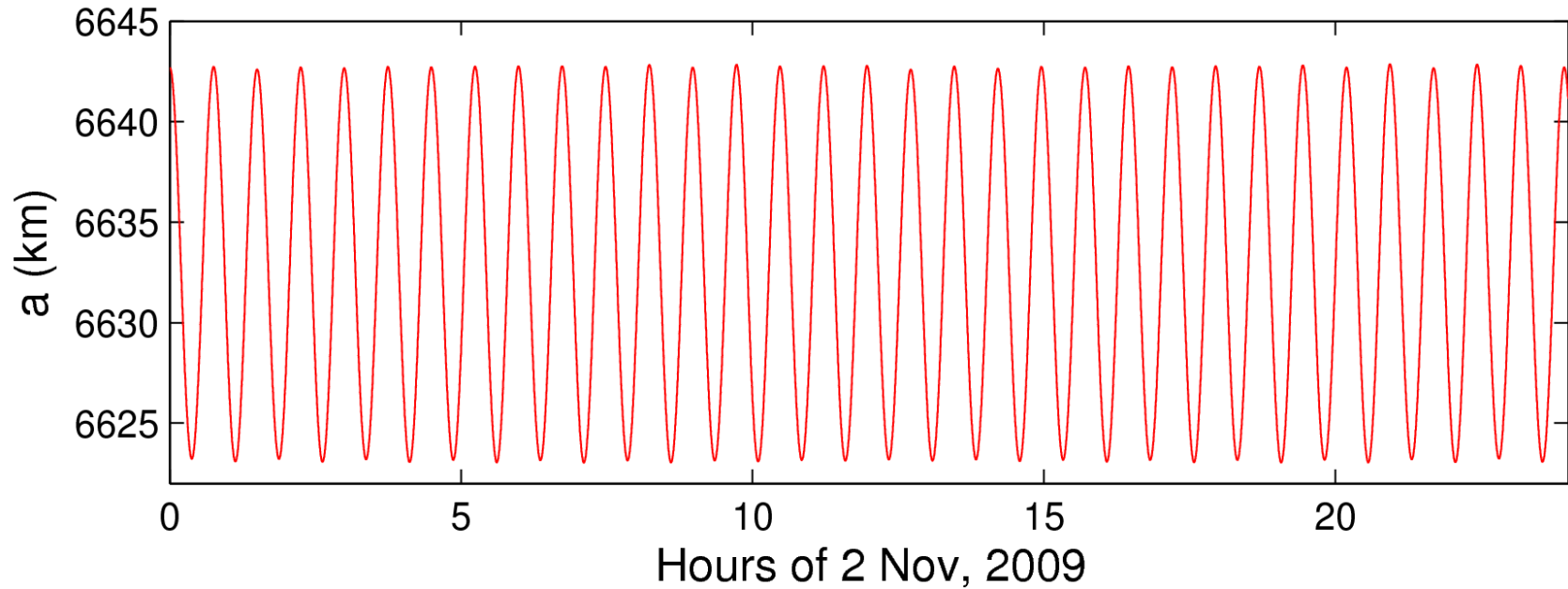
# Osculating Orbital Elements



(Beutler, 2005)

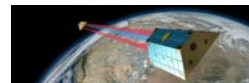


# Osculating Orbital Elements of GOCE (1)

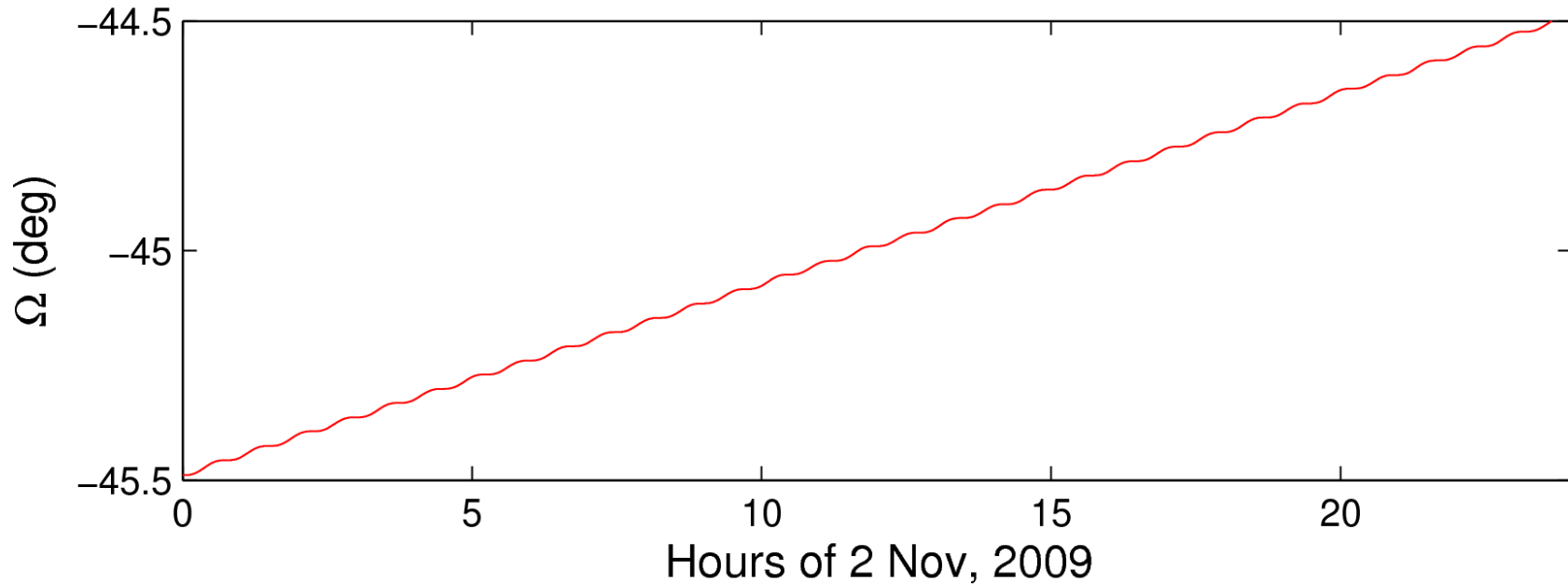


## Semi-major axis:

Twice-per-revolution variations of about  $\pm 10$  km around the mean semi-major axis of 6632.9 km, which corresponds to a mean altitude of 254.9 km

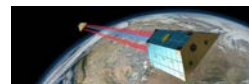


## Osculating Orbital Elements of GOCE (2)

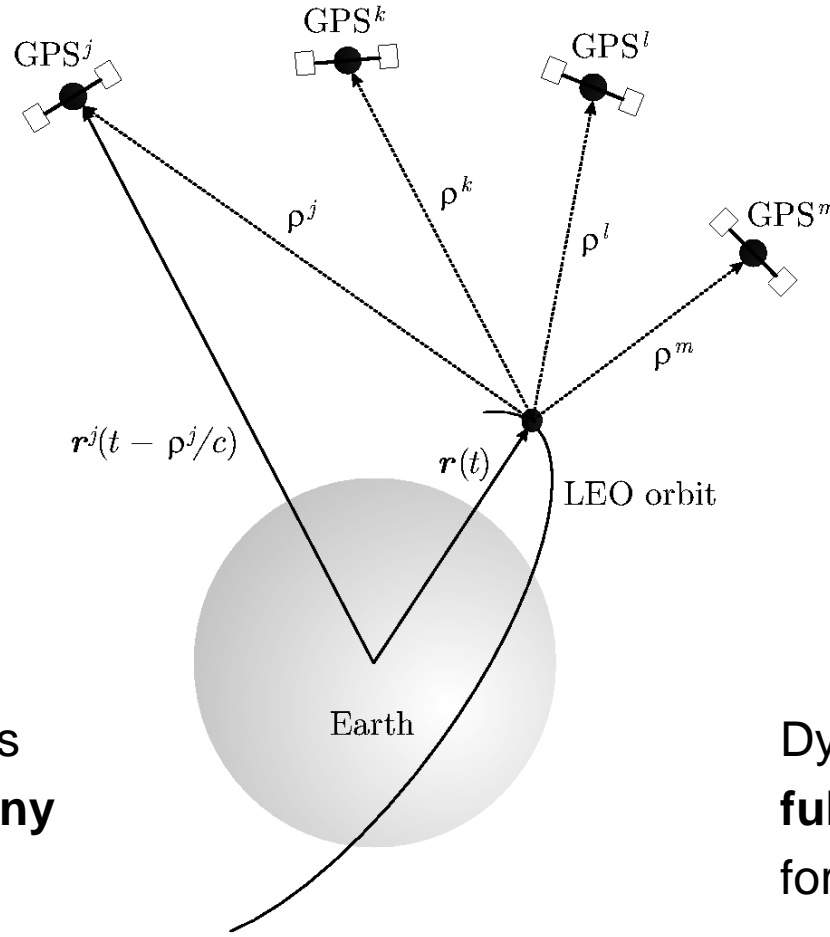


### Right ascension of ascending node:

Twice-per-revolution variations and linear drift of about  $+1^\circ/\text{day}$  ( $360^\circ/365\text{days}$ ) due to the sun-synchronous orbit

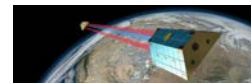


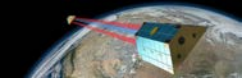
## Dynamic Orbit Representation (3)



Dynamic orbit positions may be computed at **any epoch** within the arc

Dynamic positions are **fully dependent** on the force models used, e.g., on the gravity field model





# Reduced-Dynamic Orbit Representation (1)

**Equation of motion** (in inertial frame) is given by:

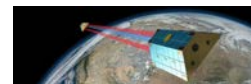
$$\ddot{\mathbf{r}} = -GM \frac{\mathbf{r}}{r^3} + \mathbf{f}_1(t, \mathbf{r}, \dot{\mathbf{r}}, Q_1, \dots, Q_d, P_1, \dots, P_s)$$

$P_1, \dots, P_s$

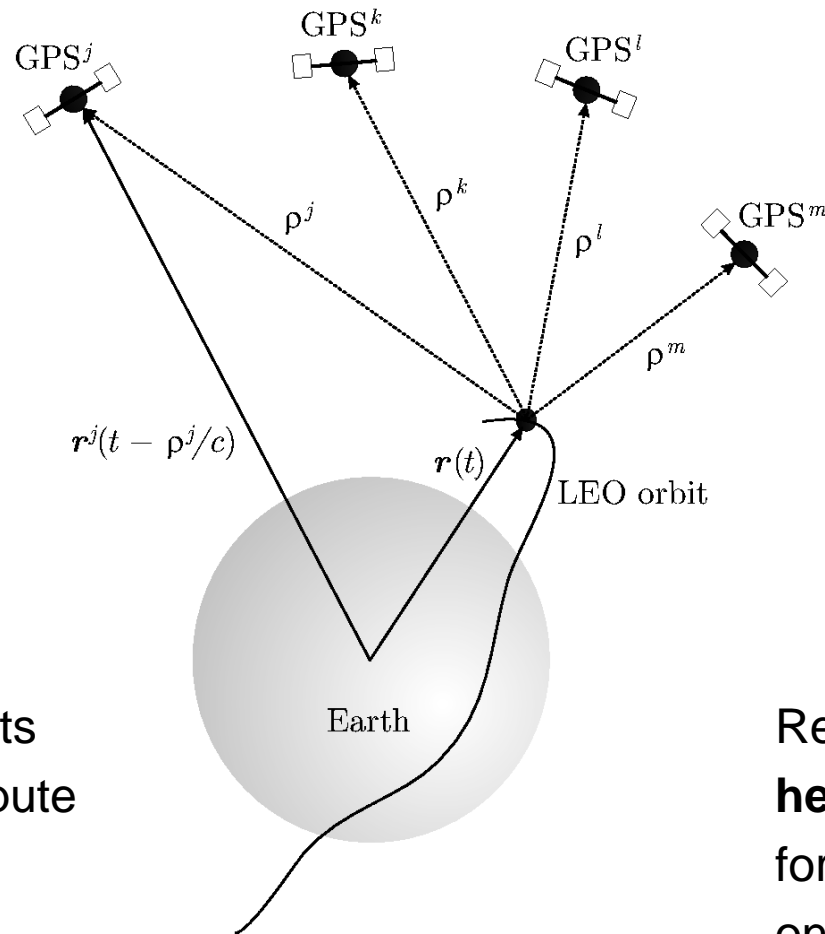
Pseudo-stochastic parameters

**Pseudo-stochastic** parameters are:

- additional empirical parameters characterized by a priori known **statistical properties**, e.g., by expectation values and a priori variances
- useful to **compensate** for deficiencies in dynamic models, e.g., deficiencies in models describing non-gravitational accelerations
- often set up as **piecewise constant accelerations** to ensure that satellite trajectories are continuous and differentiable at any epoch



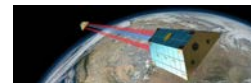
## Reduced-Dynamic Orbit Representation (2)



Reduced-dynamic orbits are well suited to compute LEO orbits of **highest quality**

(Jäggi et al., 2006; Jäggi, 2007)

Reduced-dynamic orbits **heavily depend** on the force models used, e.g., on the gravity field model





# Reduced-dynamic Orbit Representation (3)

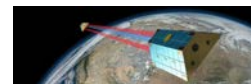
Positions (km) &  
Velocities (dm/s)  
(Earth-fixed)

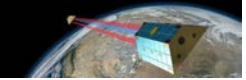
Position epochs  
(in GPS time)

* 2009 11 2	0	0	0.00000000		
PL15	-391.718353	6623.836682	79.317661	999999.999999	
VL15	13710.157683	1908.731015	-77015.601314	999999.999999	
* 2009 11 2	0	0	10.00000000		
PL15	-377.980705	6625.284690	2.298385	999999.999999	
VL15	13764.602016	987.250587	-77021.193676	999999.999999	
* 2009 11 2	0	0	20.00000000		
PL15	-364.190222	6625.811136	-74.721213	999999.999999	
VL15	13815.825127	65.631014	-77016.232293	999999.999999	
* 2009 11 2	0	0	30.00000000		
PL15	-350.350131	6625.415949	-151.730567	999999.999999	
VL15	13863.820409	-855.995477	-77000.719734	999999.999999	
* 2009 11 2	0	0	40.00000000		
PL15	-336.463660	6624.099187	-228.719134	999999.999999	
VL15	13908.581905	-1777.497047	-76974.660058	999999.999999	
* 2009 11 2	0	0	50.00000000		
PL15	-322.534047	6621.861041	-305.676371	999999.999999	
VL15	13950.104280	-2698.741871	-76938.058807	999999.999999	
* 2009 11 2	0	1	0.00000000		
PL15	-308.564533	6618.701833	-382.591743	999999.999999	
VL15	13988.382807	-3619.598277	-76890.923043	999999.999999	

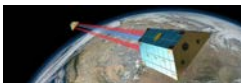
Clock corrections  
are not provided

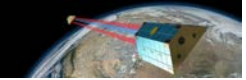
Excerpt of reduced-dynamic GOCE positions at begin of 2 Nov, 2009





# Principles of Orbit Determination





# Principle of Orbit Determination

The **actual orbit**  $\mathbf{r}(t)$  is expressed as a truncated Taylor series:

$$\mathbf{r}(t) = \mathbf{r}_0(t) + \sum_{i=1}^n \frac{\partial \mathbf{r}_0}{\partial P_i}(t) \cdot (P_i - P_{0,i})$$

$\mathbf{r}_0(t)$

A priori orbit

$\frac{\partial \mathbf{r}_0}{\partial P_i}(t)$

Partial derivative of the a priori orbit  $\mathbf{r}_0(t)$  w.r.t. parameter  $P_i$

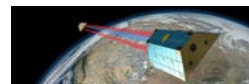
$P_{0,i}$

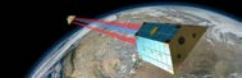
A priori parameter values of the a priori orbit  $\mathbf{r}_0(t)$

$P_i$

Parameter values of the improved orbit  $\mathbf{r}(t)$

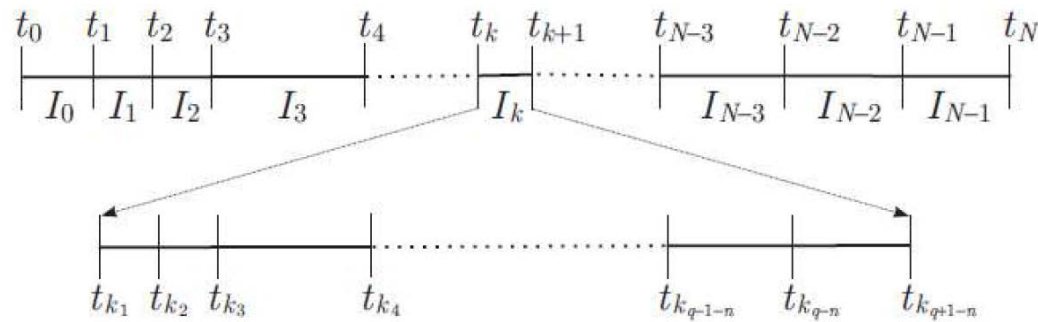
A **least-squares** adjustment of spacecraft tracking data yields **corrections** to the a priori parameter values  $P_{0,i}$ . Using the above equation, the improved (linearized) orbit  $\mathbf{r}(t)$  may be eventually computed.





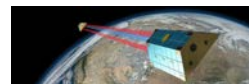
# Numerical Integration (1)

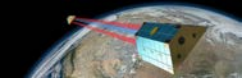
**Collocation algorithms** (one particular class of numerical integration techniques) are subsequently used to briefly illustrate the principles of numerical integration:



The original interval is divided into  $N$  integration intervals. For each interval  $I_k$  a further subdivision is performed according to the order  $q$  of the adopted method. At these points  $t_{k_j}$  the numerical solution is requested to solve the differential equation system of order  $n$ .

(Beutler, 2005)





## Numerical Integration (2)

**Initial value problem** in the interval  $I_k$  is given by:

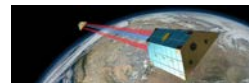
$$\ddot{\mathbf{r}}_k = \mathbf{f}(t, \mathbf{r}_k, \dot{\mathbf{r}}_k)$$

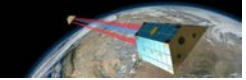
with initial conditions

$$\mathbf{r}_k(t_k) \doteq \mathbf{r}_{k0} \quad \text{and} \quad \dot{\mathbf{r}}_k(t_k) \doteq \dot{\mathbf{r}}_{k0}$$

where the initial values are defined as

$$\mathbf{r}_{k0}^{(i)} = \begin{cases} \mathbf{r}_0^{(i)} & ; k = 0 \\ \mathbf{r}_{k-1}^{(i)}(t_k) & ; k > 0 \end{cases}$$





## Numerical Integration (3)

The **collocation method** approximates the solution in the interval  $I_k$  by:

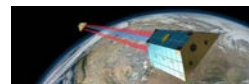
$$\mathbf{r}_k(t) \doteq \sum_{l=0}^q \frac{1}{l!} (t - t_k)^l \mathbf{r}_{k0}^{(l)}$$

The coefficients  $\mathbf{r}_{k0}^{(l)}$ ,  $l = 0, \dots, q$  are obtained by requesting that the numerical solution assumes the initial values and solves the differential equation system at  $q - 1$  different epochs  $t_{kj}$ ,  $j = 1, \dots, q - 1$ . This leads to the conditions

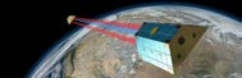
$$\sum_{l=2}^q \frac{(t_{kj} - t_k)^{l-2}}{(l-2)!} \mathbf{r}_{k0}^{(l)} = \mathbf{f}(t_{kj}, \mathbf{r}_k(t_{kj}), \dot{\mathbf{r}}_k(t_{kj})), \quad j = 1, \dots, q - 1.$$

They are non-linear but can be solved efficiently by an iterative procedure.

(Beutler, 2005)







## Partial Derivatives

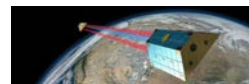
The partial of the  $r$ -th observation w.r.t. orbit parameter  $P_i$  may be expressed as

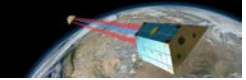
$$\frac{\partial F_r(\mathbf{X})}{\partial P_i} = (\nabla (F_r(\mathbf{X})))^T \cdot \frac{\partial \mathbf{r}_0}{\partial P_i}(t)$$

with the gradient given by

$$(\nabla (F_r(\mathbf{X})))^T = \left( \frac{\partial F_r(\mathbf{X})}{\partial r_{0,1}} \quad \frac{\partial F_r(\mathbf{X})}{\partial r_{0,2}} \quad \frac{\partial F_r(\mathbf{X})}{\partial r_{0,3}} \right)$$

if the observations only depend on the geocentric position vector and are referring to only one epoch. The gradient only depends on the type of observations used, whereas the second term is independent of the observation type and related to the **variational equations**. This separates the observation-specific (**geometric**) part from the **dynamic** part.





## Variational Equations

For each orbit parameter  $P_i$  the corresponding **variational equation** reads as

$$\ddot{\mathbf{z}}_{P_i} = \mathbf{A}_0 \cdot \mathbf{z}_{P_i} + \mathbf{A}_1 \cdot \dot{\mathbf{z}}_{P_i} + \frac{\partial \mathbf{f}_1}{\partial P_i}$$

with  $\mathbf{z}_{P_i}(t) \doteq \frac{\partial \mathbf{r}_0}{\partial P_i}(t)$  and the 3 x 3 matrices defined by

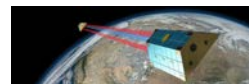
$$A_{0[i;k]} \doteq \frac{\partial f_i}{\partial r_{0,k}} \quad \text{and} \quad A_{1[i;k]} \doteq \frac{\partial f_i}{\partial \dot{r}_{0,k}}$$

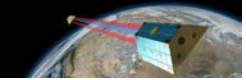
$f_i$   $i$ -th component of the total acceleration  $\mathbf{f}$

$r_{0,k}$   $k$ -th component of the geocentric position  $\mathbf{r}_0$

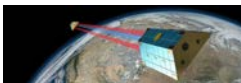
For each orbit parameter  $P_i$  the **variational equation** is a linear differential equation system of second order in time. Their solutions are all needed for orbit determination.

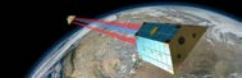
(Jäggi, 2007)





# GPS-based LEO POD

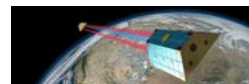




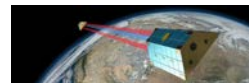
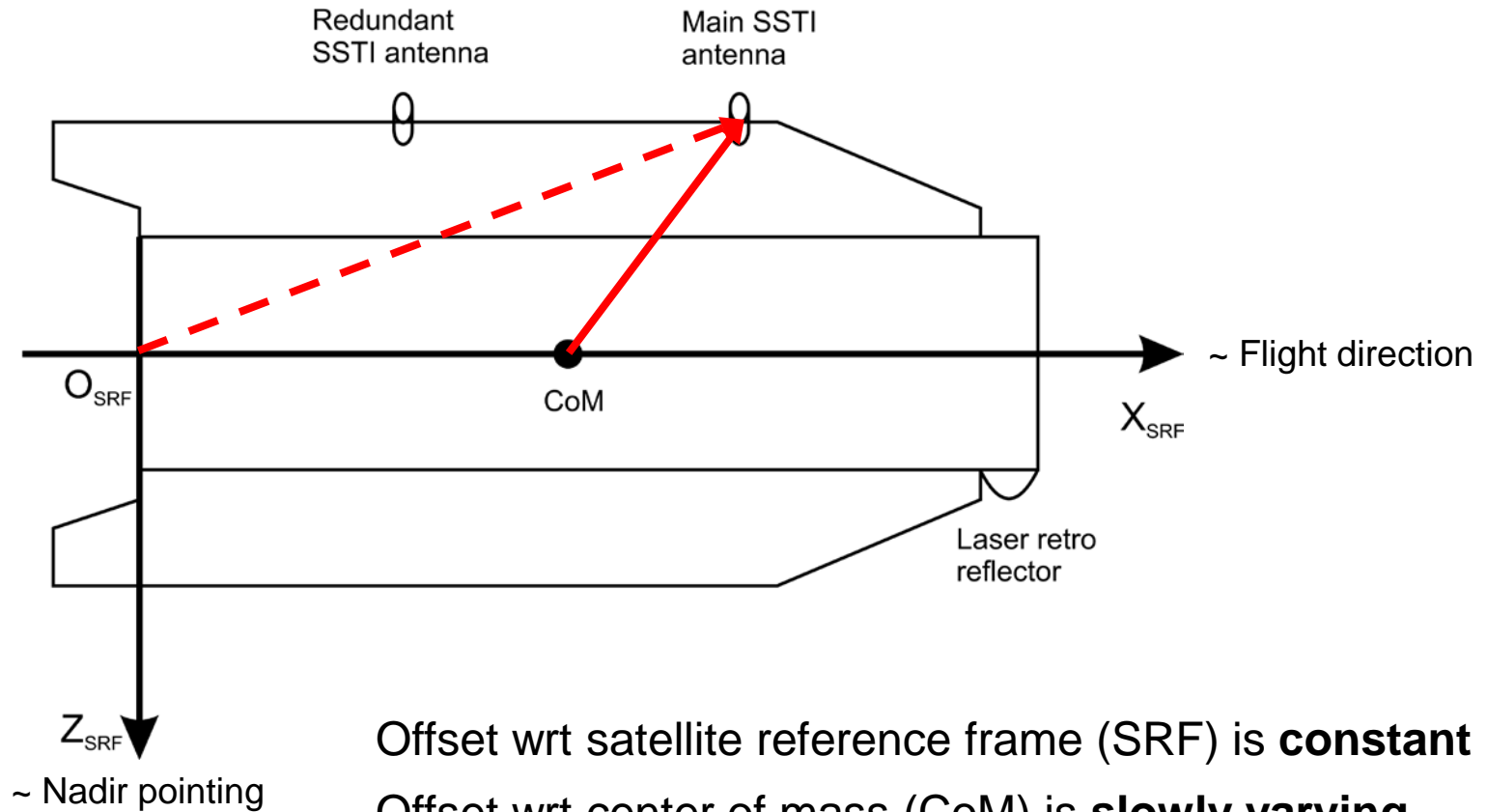
# LEO Sensor Offsets

## Phase center offsets $\delta r_{leo,ant}$ :

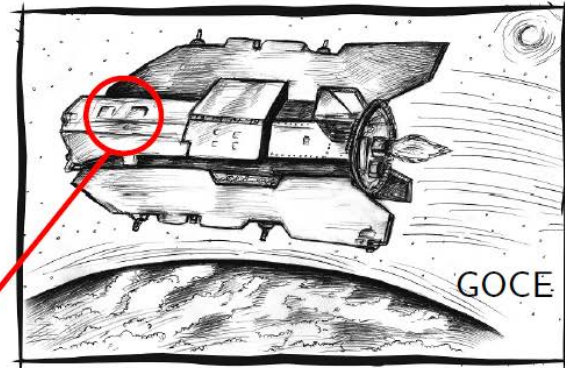
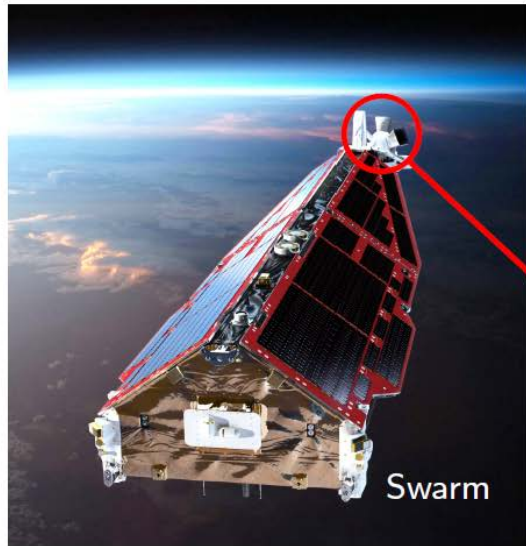
- are needed in the inertial or Earth-fixed frame and have to be transformed from the satellite frame using **attitude data** from the star-trackers
- consist of a frequency-independent **instrument offset**, e.g., defined by the center of the instrument's mounting plane (CMP) in the satellite frame
- consist of frequency-dependent **phase center offsets** (PCOs), e.g., defined wrt the center of the instrument's mounting plane in the antenna frame (ARF)
- consist of frequency-dependent **phase center variations** (PCVs) varying with the direction of the incoming signal, e.g., defined wrt the PCOs in the antenna frame



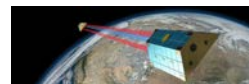
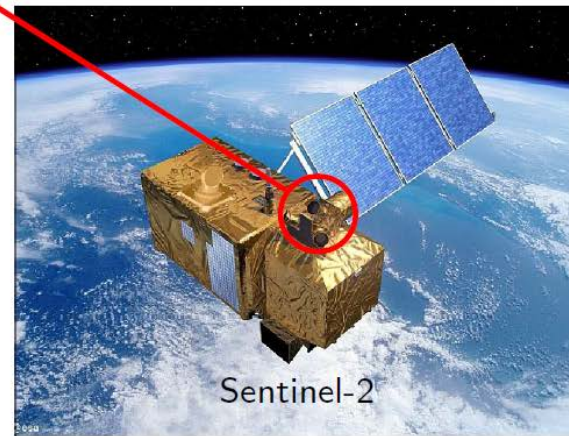
# Example: GOCE Sensor Offsets (1)



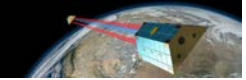
# Example: GOCE Sensor Offsets (2)



Star tracker

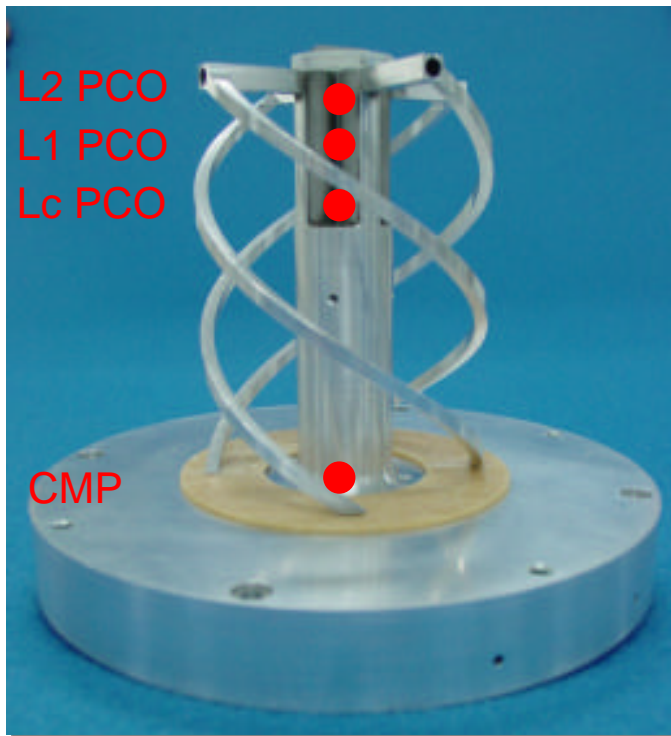






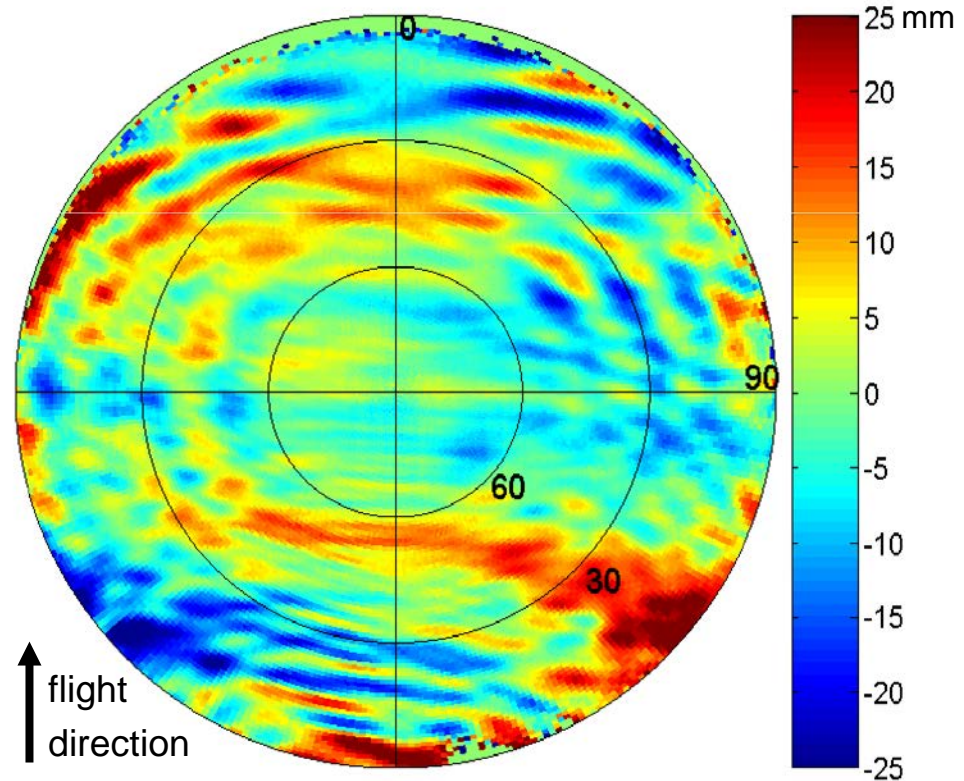
# Spaceborne GPS Antennas: GOCE

## L1, L2, Lc phase center offsets:

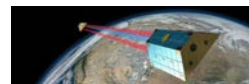


Measured from ground calibration in anechoic chamber

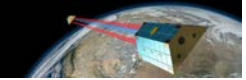
## Lc phase center variations



Empirically derived during orbit determination according to Jäggi et al. (2009)

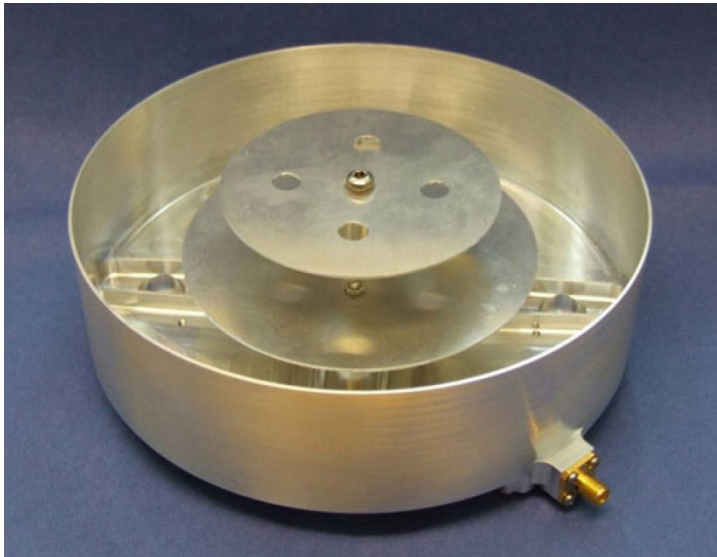






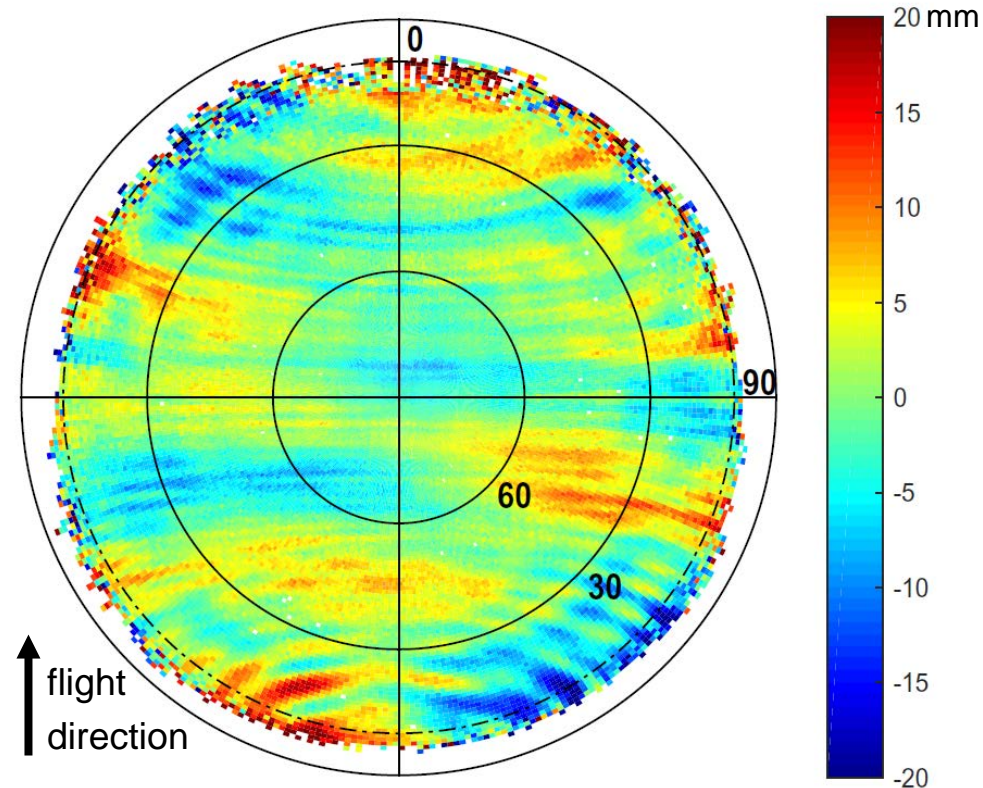
# Spaceborne GPS Antennas: Swarm

## Swarm GPS antenna

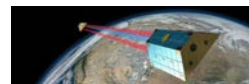


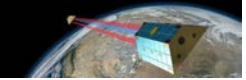
Multipath shall be minimized by chokering

## Lc phase center variations

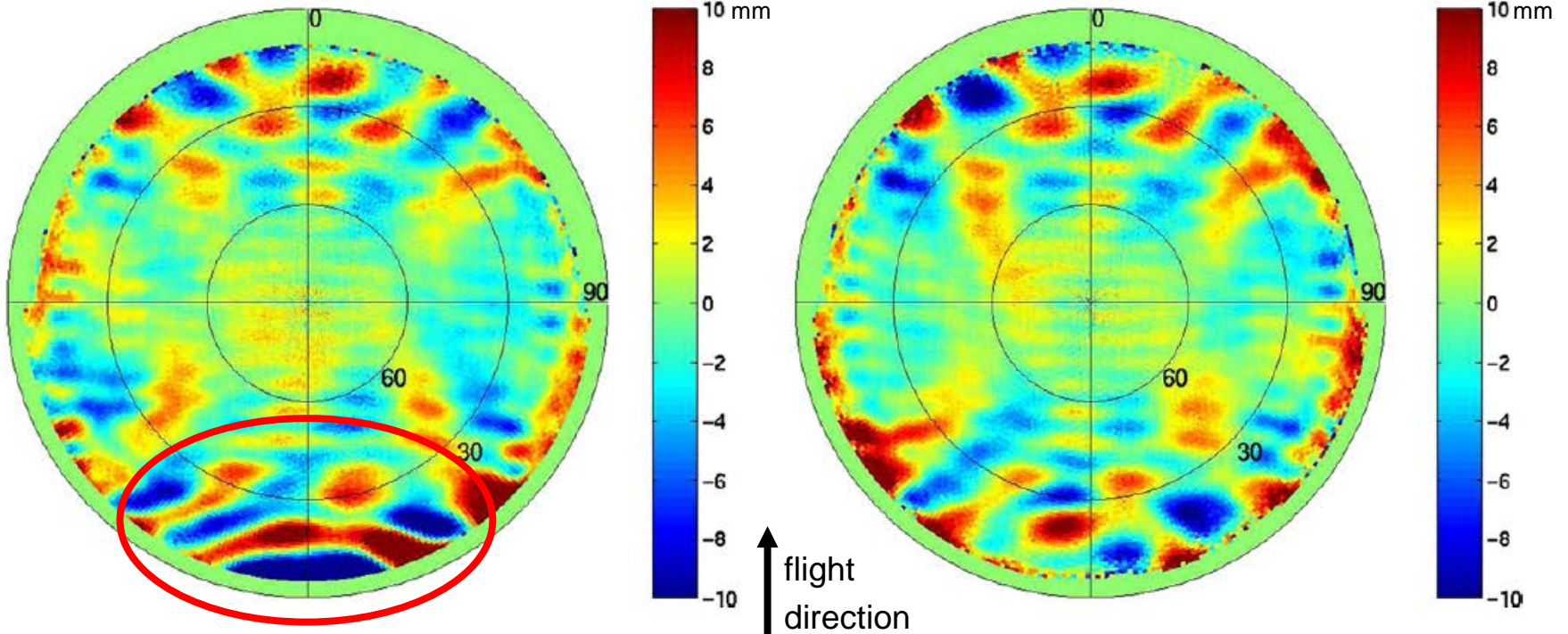


Empirically derived during orbit determination according to Jäggi et al. (2016)





# Spaceborne GPS Antennas: GRACE

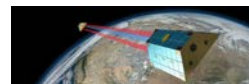


GRACE-A

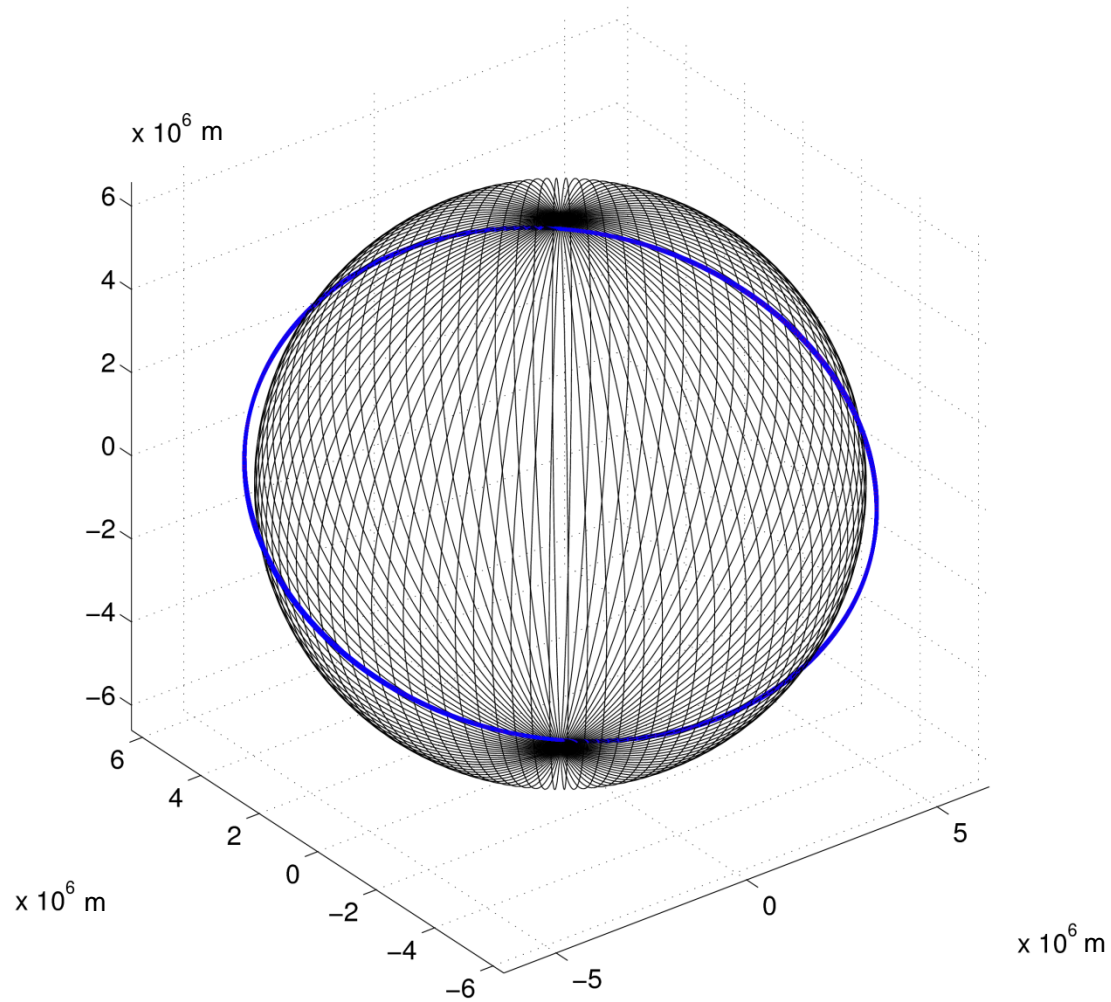
(occultation antenna switched on)

GRACE-B

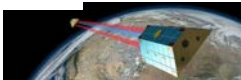
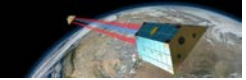
(Jäggi et al., 2009)



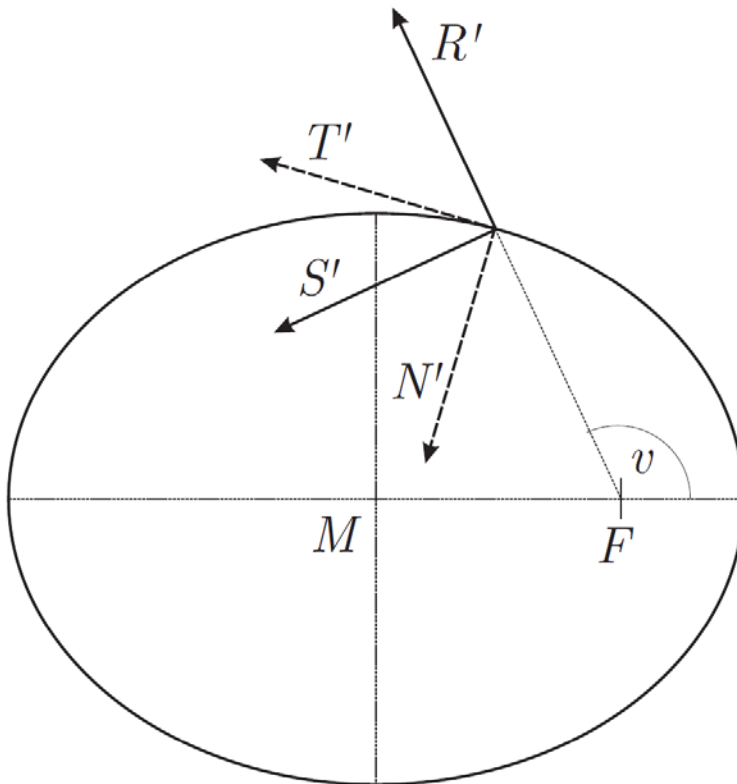
# Visualization of Orbit Solutions



It is more instructive to look at differences between orbits in well suited coordinate systems ...



# Co-Rotating Orbital Frames



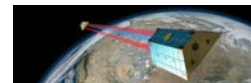
**R, S, C** unit vectors are pointing:

- into the radial direction
- normal to **R** in the orbital plane
- normal to the orbital plane (cross-track)

**T, N, C** unit vectors are pointing:

- into the tangential (along-track) direction
- normal to **T** in the orbital plane
- normal to the orbital plane (cross-track)

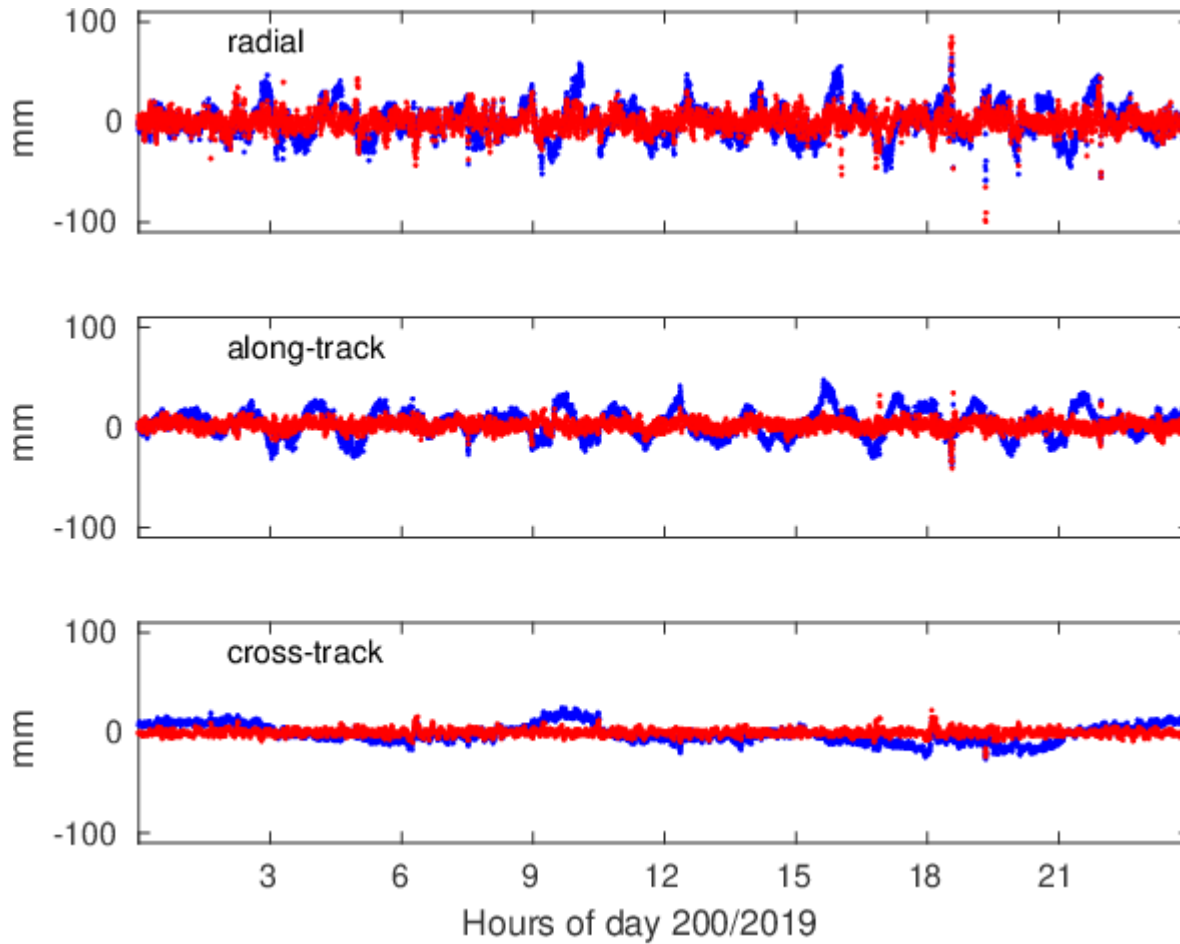
Small eccentricities: **S~T** (velocity direction)



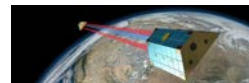


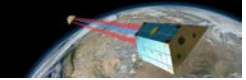
# Orbit Differences KIN-RD (Sentinel-3A)

Differences at epochs of kin. positions



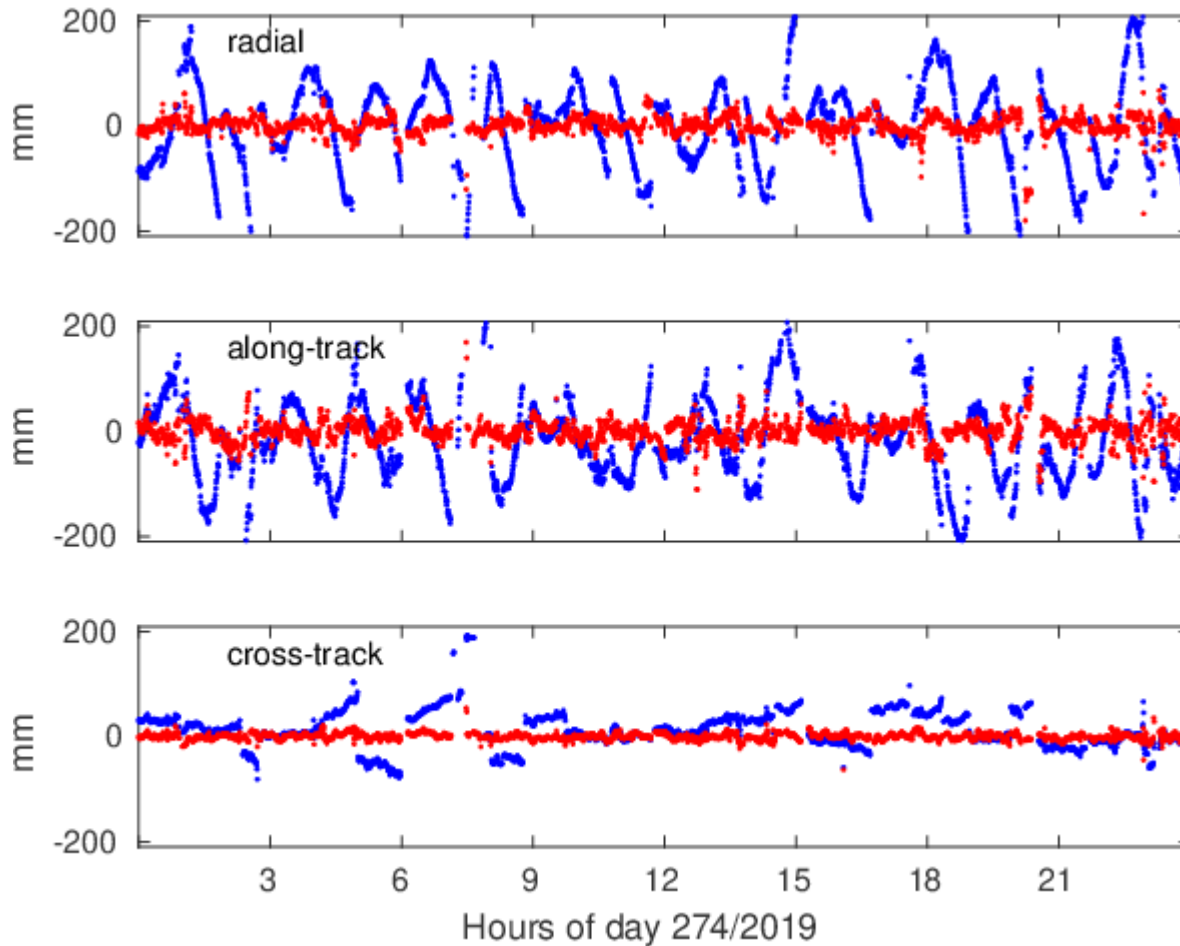
Comparison of **ambiguity-float** solutions and **ambiguity-fixed** solutions.



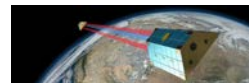


## Orbit Differences KIN-RD (COSMIC-2)

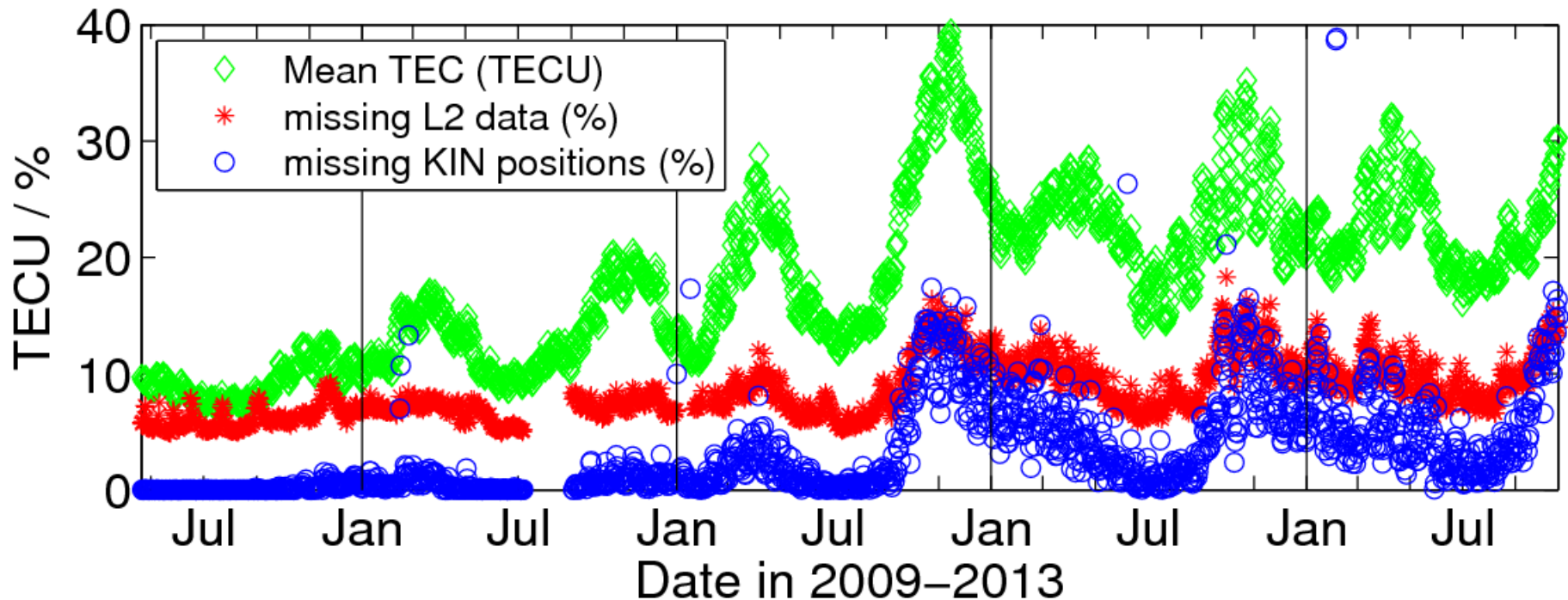
Differences at epochs of kin. positions, FM-1, POD-1



Comparison of **ambiguity-float** solutions and **ambiguity-fixed** solutions. (Jäggi et al. 2021)



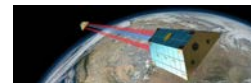
## Orbit Differences KIN-RD (GOCE)



The result illustrates the **consistency** between both orbit-types. The level of the differences is usually given by the quality of the kinematic positions.

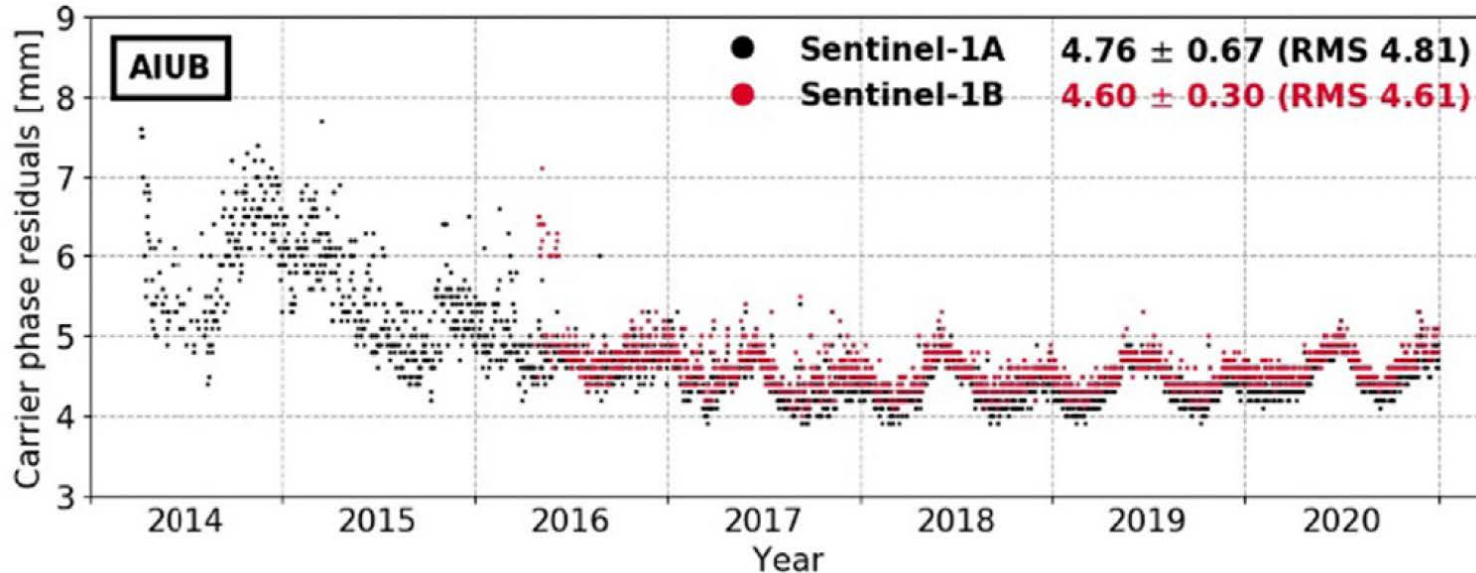
The differences are highly correlated with the **ionosphere activity** and with **data losses on L2**.

(Bock et al., 2014)





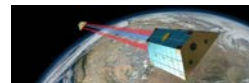
# RMS of Carrier Phase Residuals (Sentinel-1)



The plot illustrates the **quality** of the orbit fit. The level is given by the adopted parametrization, depending on how dynamic the orbit parametrization is.

The variations reveal again the impact of the **ionosphere activity** and also further **modeling deficiencies**.

(Fernández et al., 2022)

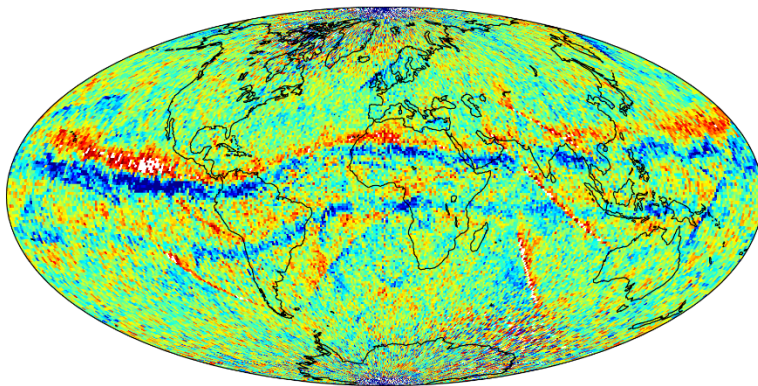


# Consequences of Ionospheric Effects in Orbits

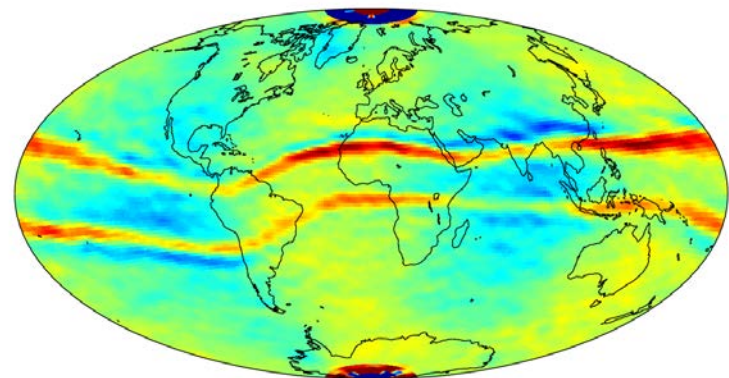
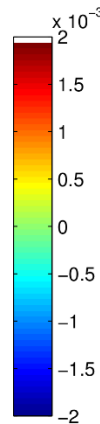
For GOCE systematic effects around the geomagnetic equator were observed in the ionosphere-free GPS phase residuals => **affects kinematic positions**

Degradation of kinematic positions around the geomagnetic equator propagates into gravity field solutions.

mean residuals at ionosphere-crossing: 2011, doys 245–365

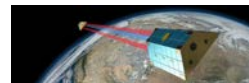


Phase observation residuals  
(- 2 mm ... +2 mm) mapped  
to the ionosphere piercing  
point



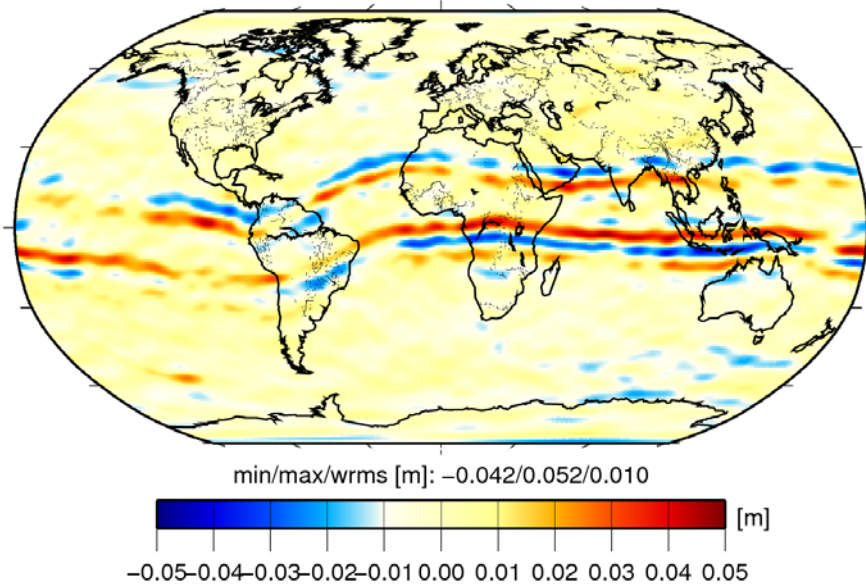
Geoid height differences  
(-5 cm ... 5 cm);  
**R4 period**

(Jäggi et al., 2015)

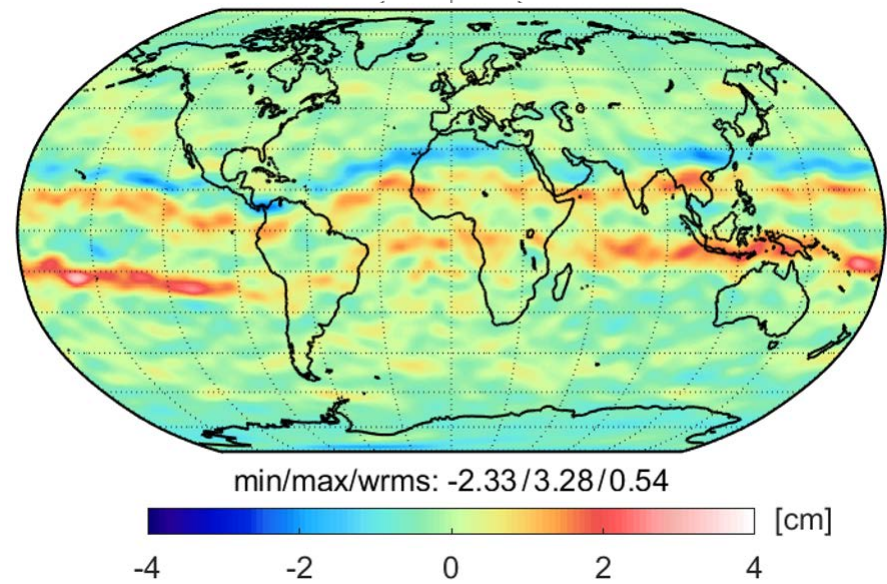


# Systematic Errors in GPS Data (1)

Original GPS Data  
(Swarm)



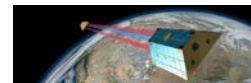
Original GPS Data  
(GRACE)



(Differences wrt GOCO05S, 400 km Gauss smoothing adopted)

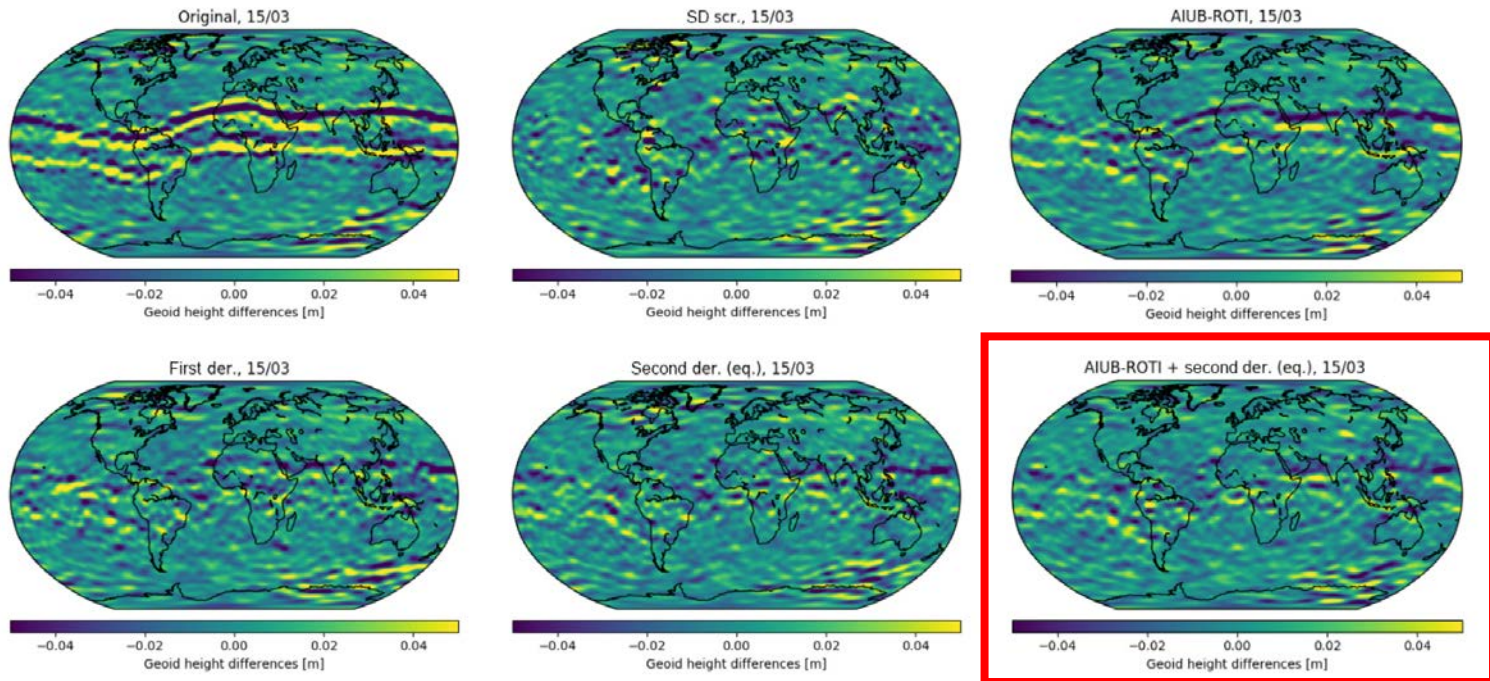
Systematic signatures along the geomagnetic equator are "**not**" visible when using original L1B RINEX GPS data files from the GRACE mission.

(Jäggi et al., 2016)





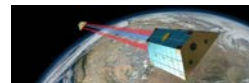
## Systematic Errors in GPS Data (2)



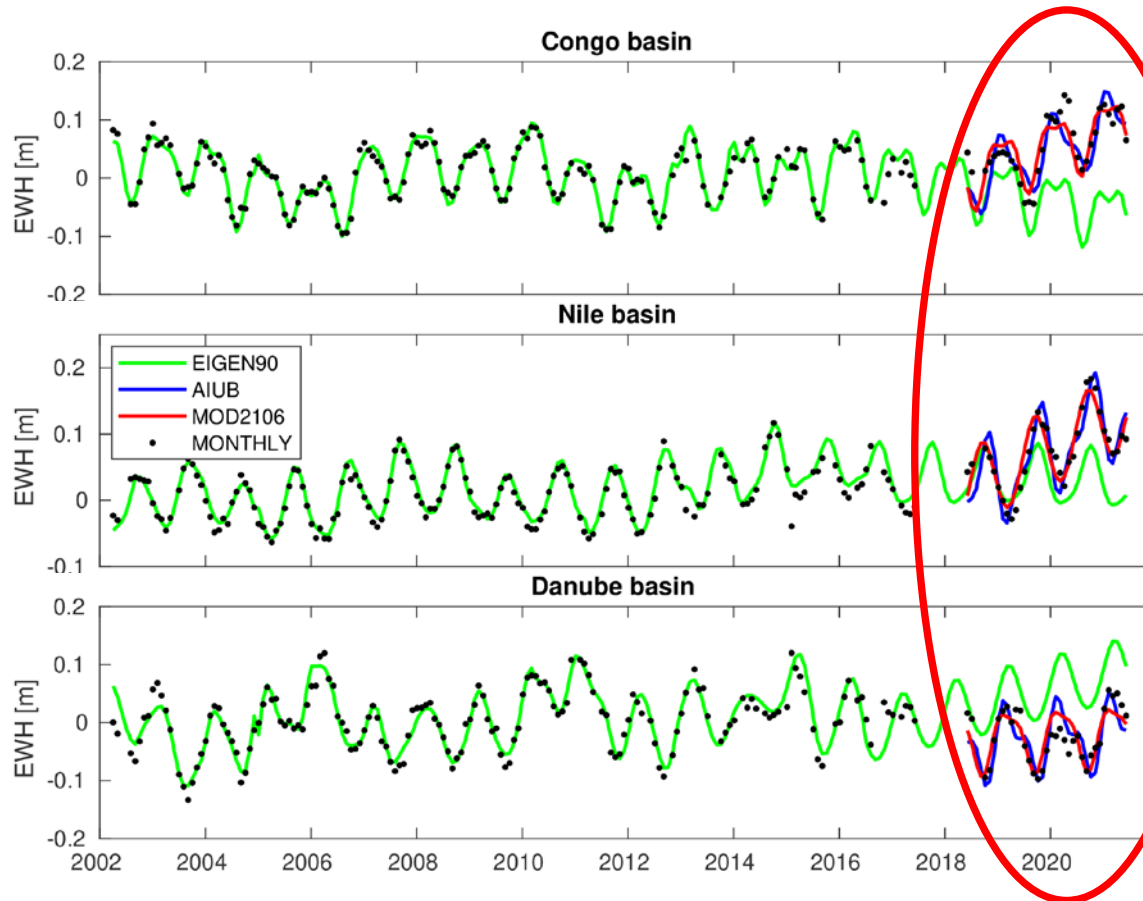
(Differences wrt JPL-GRACE-RL06, 400 km Gauss smoothing adopted)

Systematic signatures along the geomagnetic equator may be efficiently reduced when down-weighting the GPS data using derivatives of the geometry-free linear combination. ROTI-based down-weighting additionally reduces scintillation noise.

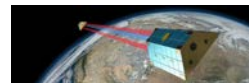
(Schreiter et al., 2019)



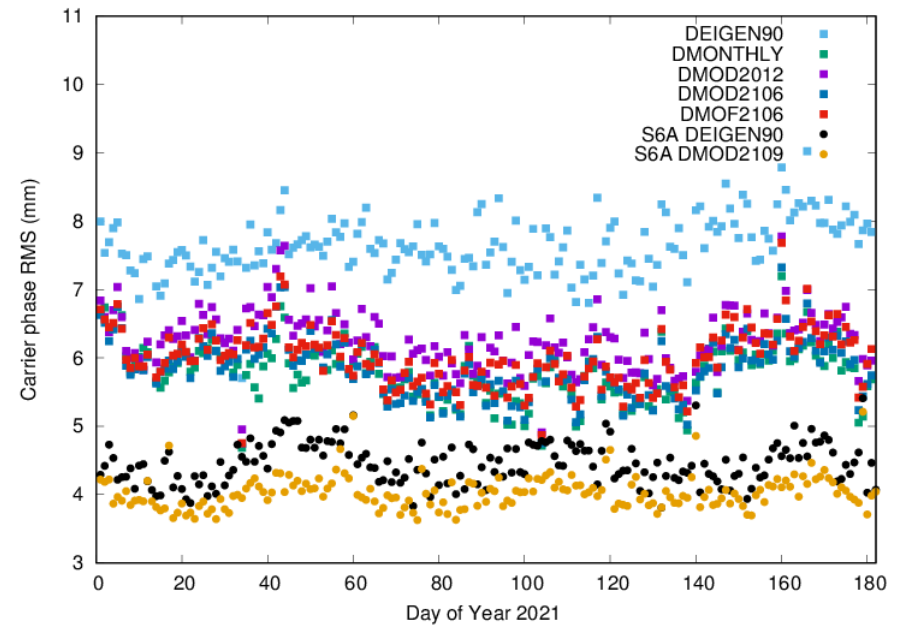
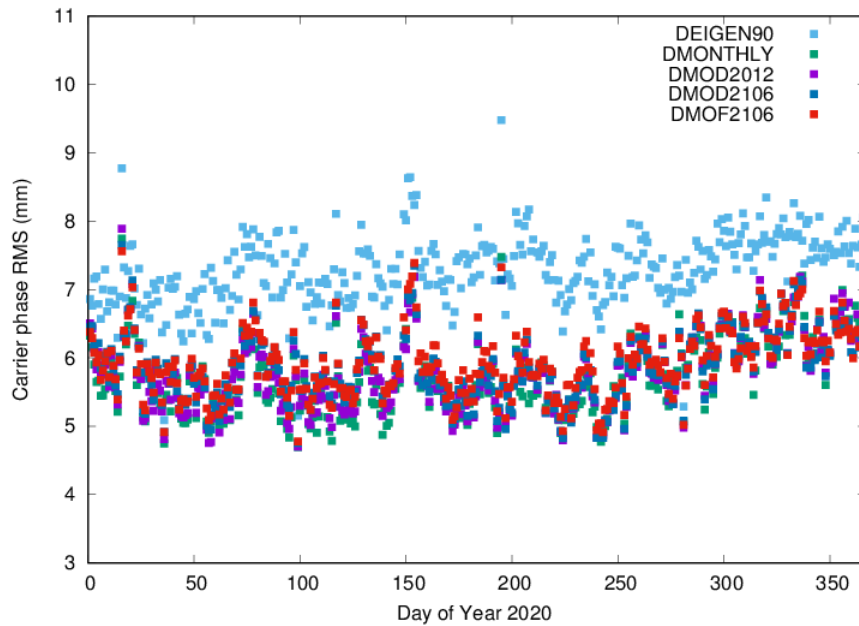
# Sensitivity to the Time-Varyable Gravity Field Model (1)



The predictions of the EIGEN-GRGS.RL04 model (containing no data after 2017) are rather poor as shown here for three river basins. The new fitted signal model (FSM) to COST-G monthly solutions has clear advantages. (Peter et al., 2022)



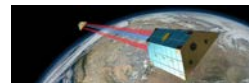
## Sensitivity to the Time-VARIABLE Gravity Field Model (2)

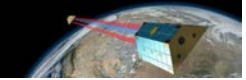


Monthly COST-G GRACE-FO gravity fields (■) outperform GRACE-based static fields (□) with co-estimated time-variations for Sentinel-3B (left and right plot).

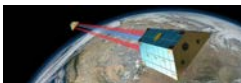
COST-G fitted signal models (FSM) perform comparably good (■ □ ■, left plot), which also holds for Sentinel-6A (●, right plot). In particular the COST-G FSM show also good prediction capabilities (■, right plot).

(Peter et al., 2022)



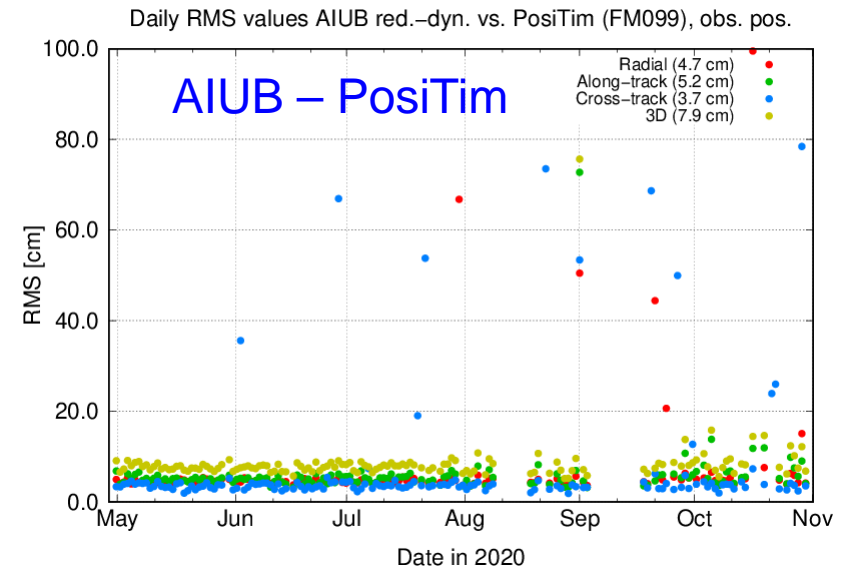
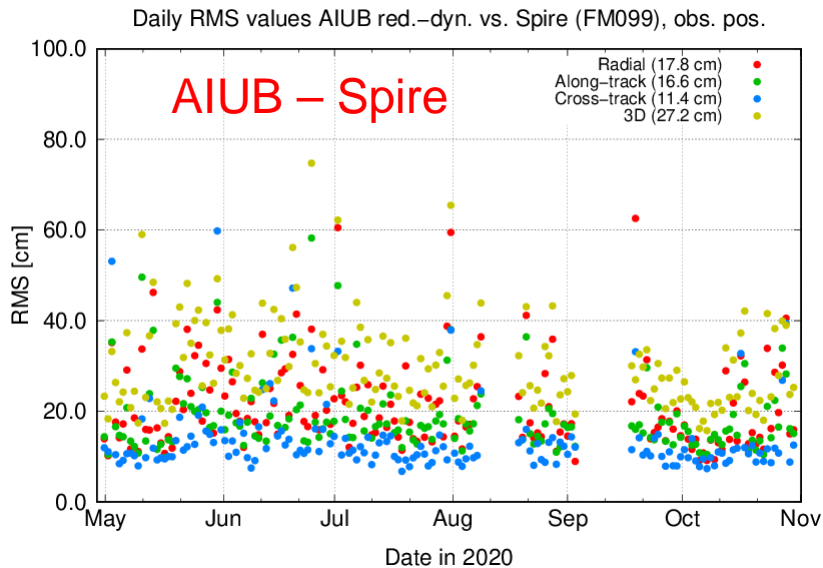


# Orbit Validation





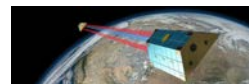
# Orbit Comparisons (Spire FM099)



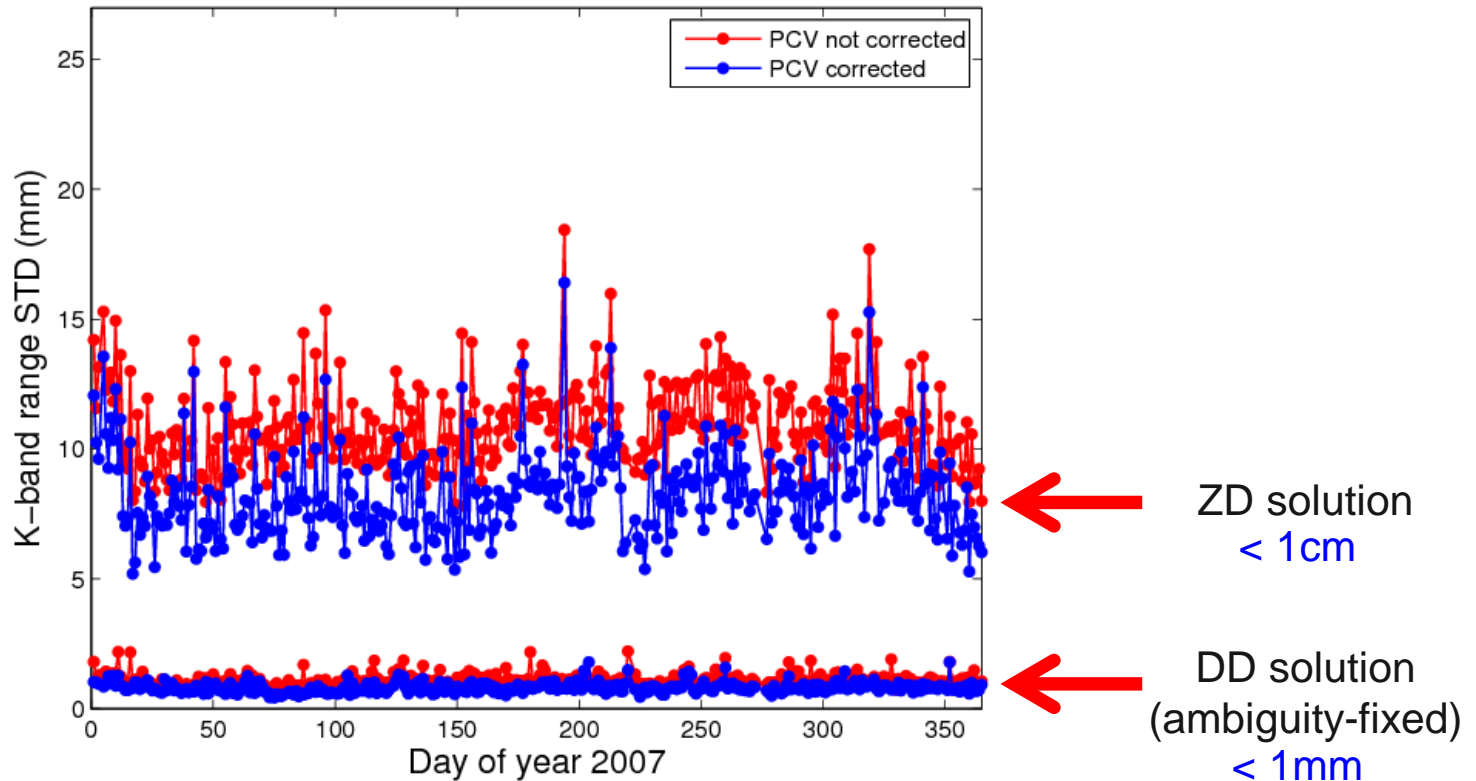
Orbit comparisons between solutions computed **with different software packages** are helpful, especially if no external orbit validations are possible.

The plots show for the example of one Spire satellite that orbit differences **AIUB – PosiTim** are significantly smaller than the orbit differences **AIUB – Spire**.

(Jäggi et al., 2022)

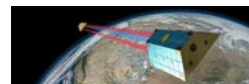


# GRACE Orbit Validation with K-Band

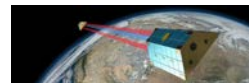
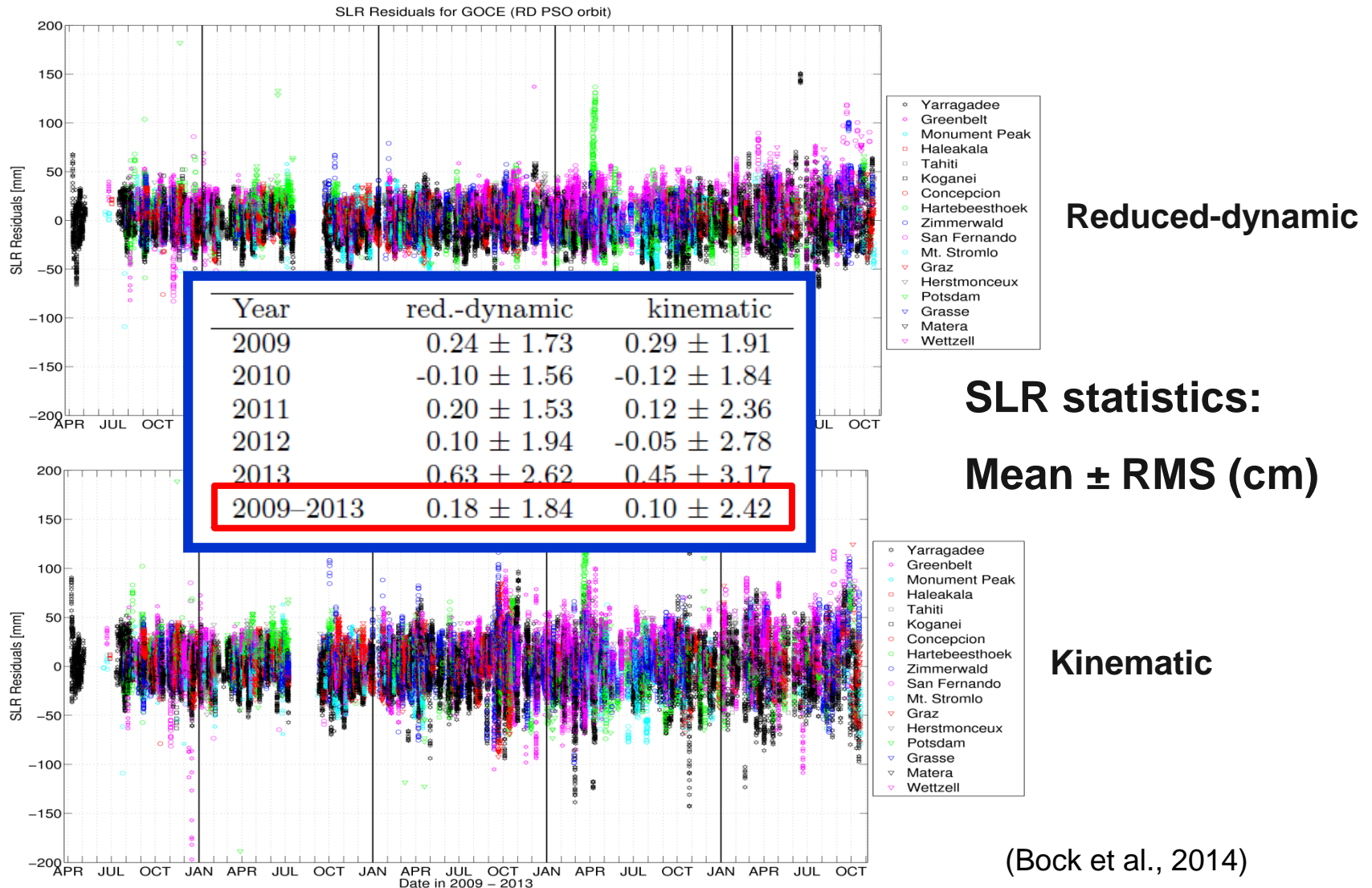


The ultra-precise and continuously available K-Band data allow it to validate the **inter-satellite distances** between the GRACE satellites. Thanks to this validation, e.g., PCV maps were recognized to be crucial for high-quality POD.

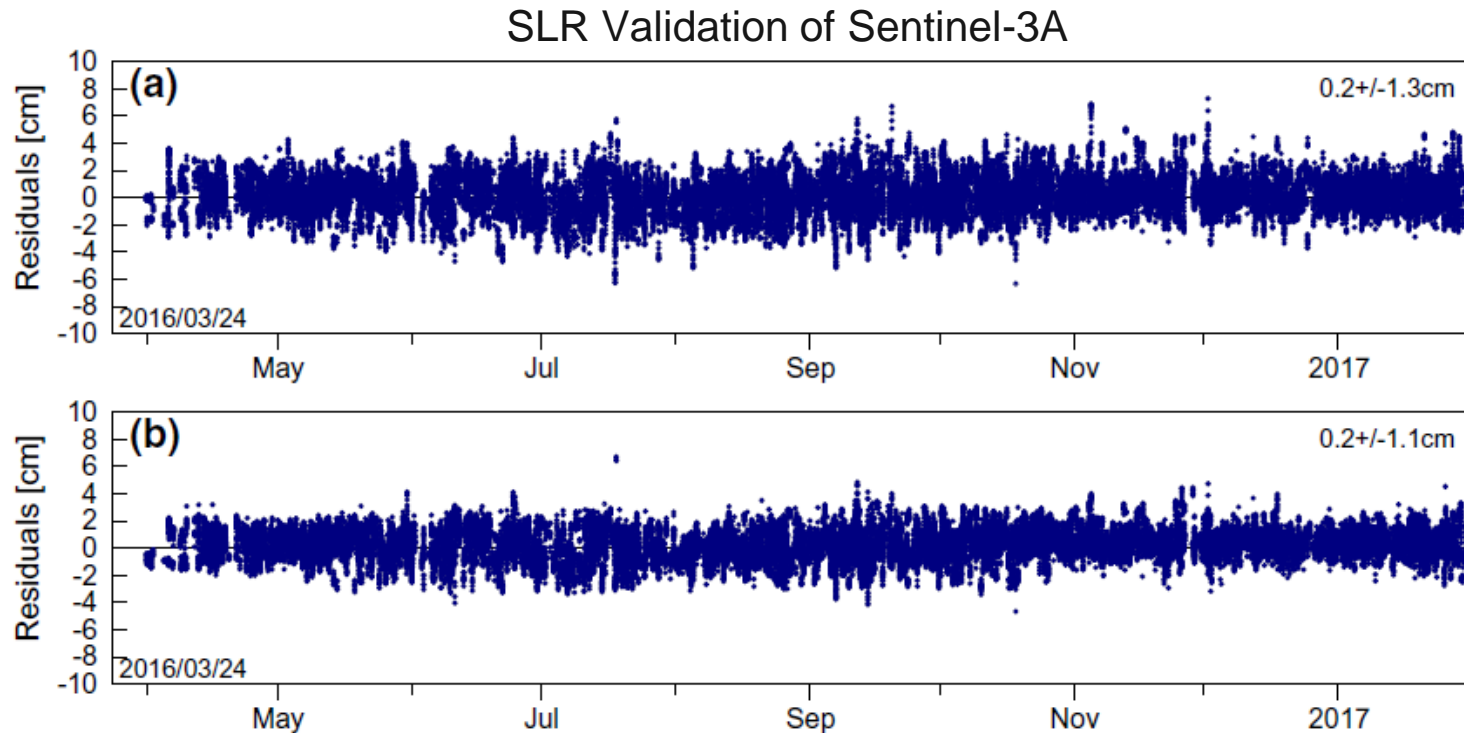
(Jäggi et al., 2009)



# Orbit Validation with SLR (GOCE)

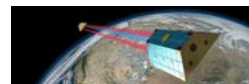


# Impact of Undifferenced Ambiguity Resolution (1)

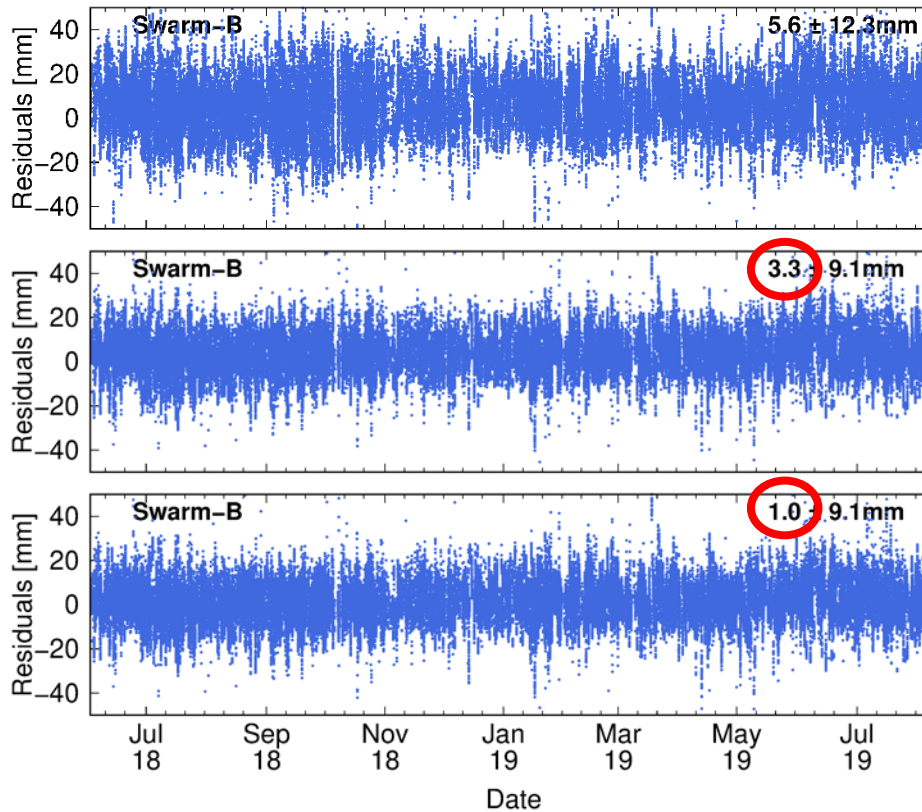
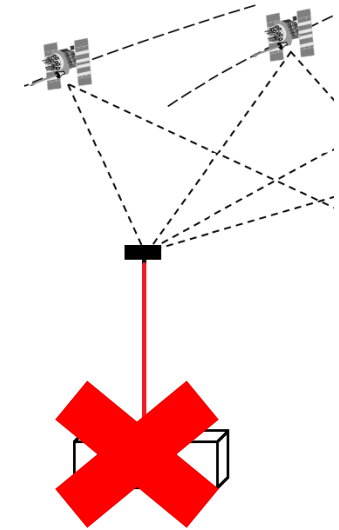
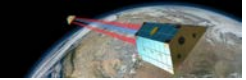


**Single-receiver ambiguity fixing** may be enabled by using phase bias products and corresponding clock products provided by the IGS analysis centers without the need to form any baselines. It allows to identify lateral offsets in the GPS antenna or center-of-mass location and to significantly stabilize the LEO trajectories.

(Montenbruck et al., 2018)



# Impact of Undifferenced Ambiguity Resolution (2)



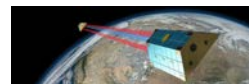
Ambiguity-float, no non-grav. modeling

Ambiguity-fixed, no non-grav. modeling

Ambiguity-fixed, with non-grav. modeling

LEO POD significantly profits from single-receiver ambiguity fixing techniques and high-quality signal-specific phase bias products, e.g., by Schaer et al. (2021).

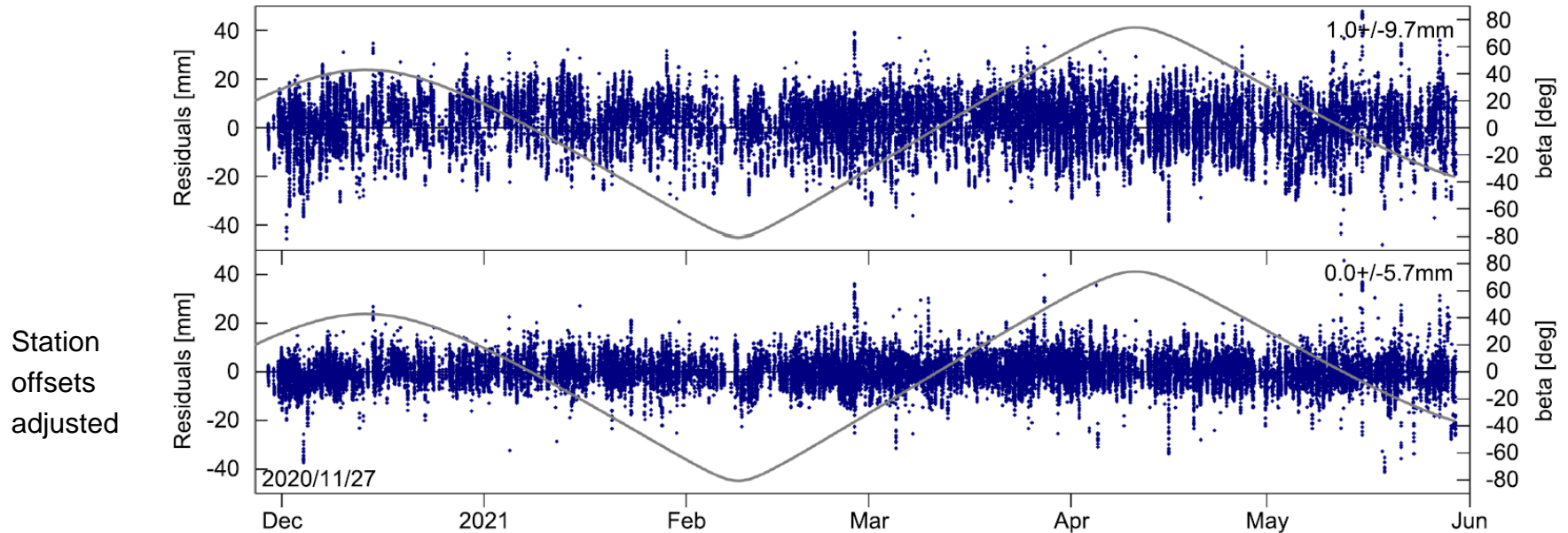
(Arnold et al., 2019; Mao et al., 2021)





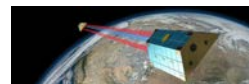
# Multi-GNSS LEO POD

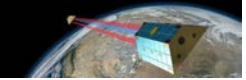
## SLR Validation of Sentinel-6A



While **Galileo** measurements exhibit **30–50% smaller RMS errors** than those of GPS, the POD benefits most from the availability of an increased number of satellites. For Sentinel-6A a 1-cm consistency of ambiguity-fixed GPS-only and Galileo-only solutions with the dual-constellation orbits can be demonstrated.

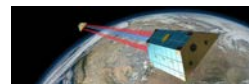
(Montenbruck et al., 2021)



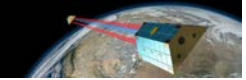


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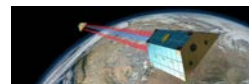


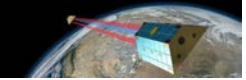




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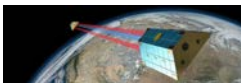
## Literature (3)

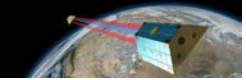
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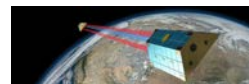
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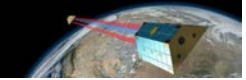




## Literature (4)

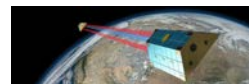
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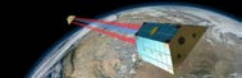




## Literature (5)

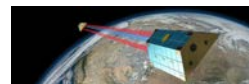
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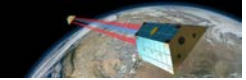




## Literature (6)

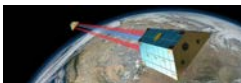
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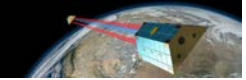


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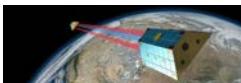
# Kinematic LEO Orbit Products (1)

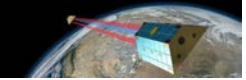
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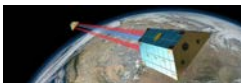
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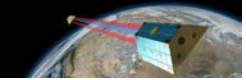
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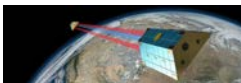
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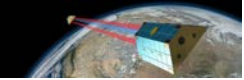




# Kinematic LEO Orbit Products (3)

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# Pocket Guide of Least-Squares Adjustment (1)

The system of **Observation Equations** is given by:

$$L' + \epsilon = F(X)$$

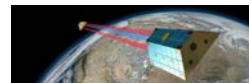
or, if  $F$  is a non-linear function of the parameters, in its **linearized** form:

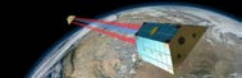
$$L' + \epsilon = F(X_0) + A x$$

$L'$	Tracking observations	$X_0$	A priori parameter values
$\epsilon$	Observation corrections	$x$	Parameter corrections
$F$	Functional model	$X$	Improved parameter values, i.e., $X = X_0 + x$

$$A \doteq \left. \frac{\partial F(X)}{\partial X} \right|_{X=X_0}$$

First design matrix





# Pocket Guide of Least-Squares Adjustment (2)

The system of **Normal Equations** is obtained by minimizing  $\epsilon^T P \epsilon$ :

$$\left( A^T P A \right) x - A^T P l = N x - b = 0$$

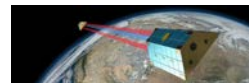
$N \doteq A^T P A$       Normal equation matrix

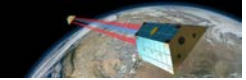
$b \doteq A^T P l$       Right-hand side with "O-C" term  $l \doteq L' - F(X_0)$

$P = \sigma_0^2 C_{ll}^{-1}$       Weight matrix, from covariance matrix  $C_{ll}$  of observations

For a **regular** normal equation matrix the parameter corrections follow as:

$$x = \left( A^T P A \right)^{-1} A^T P l = N^{-1} b$$





# Pocket Guide of Least-Squares Adjustment (3)

The **a posteriori standard deviation of unit weight** is computed as:

$$m_0 = \sqrt{\frac{\epsilon^T P \epsilon}{f}}$$

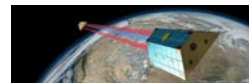
$f$  Degree of freedom (number of observations minus number of parameters)

The **covariance matrix** of the adjusted parameters is given by

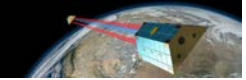
$$C_{xx} = m_0^2 Q_{xx} = m_0^2 N^{-1}$$

and their a posteriori standard deviations follow from the diagonal elements:

$$m_x = \sqrt{C_{xx}} = m_0 \sqrt{Q_{xx}}$$







# Pocket Guide of Least-Squares Adjustment (4)

**Parameter pre-elimination** is useful to handle a large number of parameters efficiently. Let us sub-divide the system of normal equations into two parts:

$$\begin{pmatrix} N_{11} & N_{12} \\ N_{21} & N_{22} \end{pmatrix} \cdot \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}$$

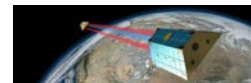
We we may reduce the normal equation system by pre-eliminating epoch-specific parameters  $x_2$  , which yields the modified system of normal equations as

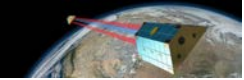
$$N_{11}^* x_1 = b_1^*$$

where

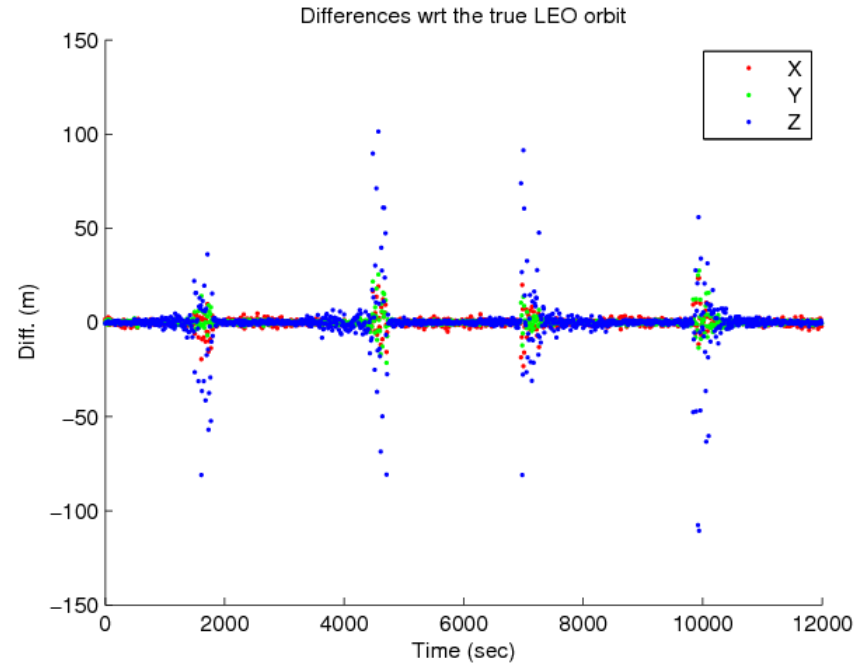
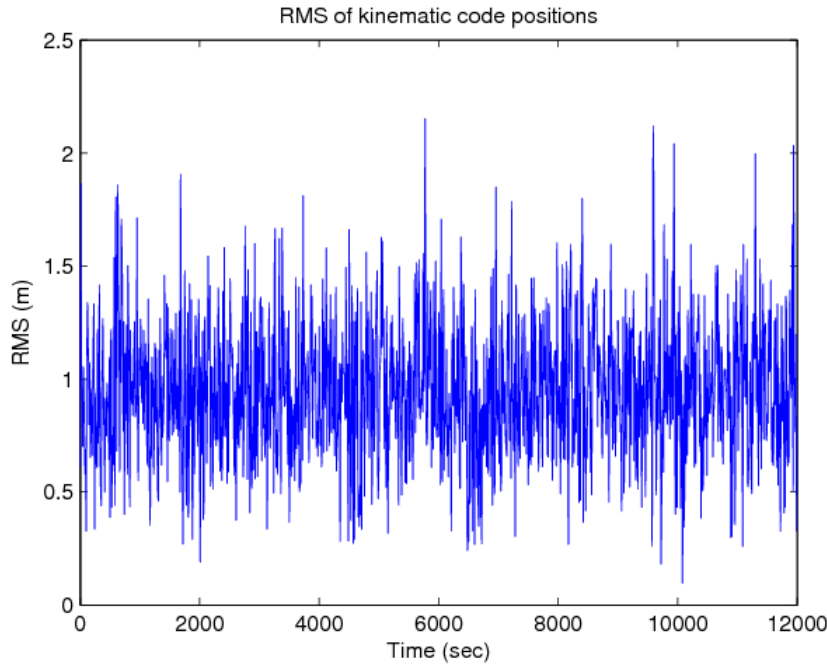
$N_{11}^* = N_{11} - N_{12} N_{22}^{-1} N_{21}$  is the normal equation matrix of  $x_1$

$b_1^* = b_1 - N_{12} N_{22}^{-1} b_2$  is the corresponding right-hand side of the normal equation system



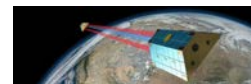


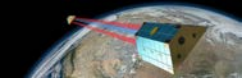
# Main Results from Lab 2 (1)



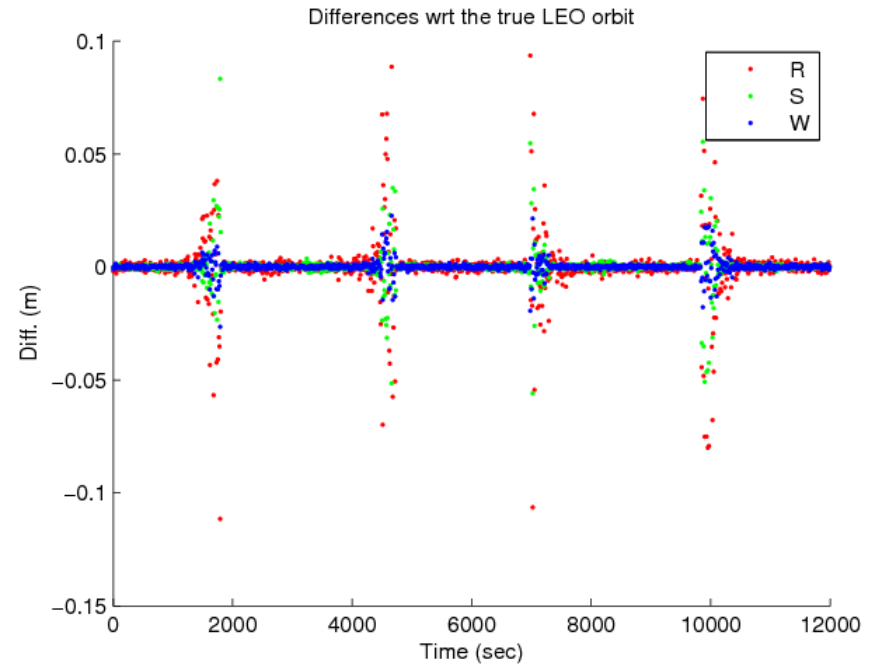
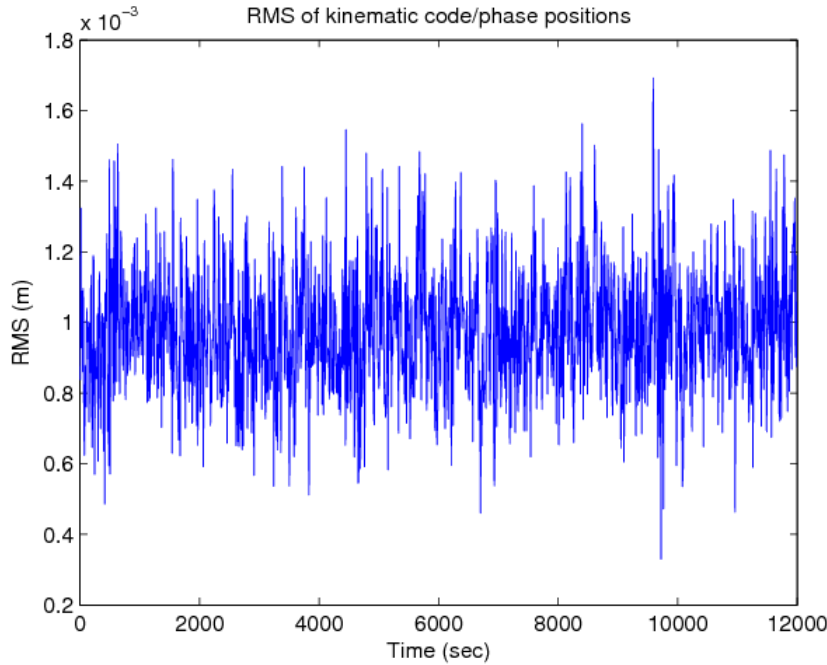
The **a posteriori RMS of unit weight** ( $m_0$ ) of a Code-only kinematic LEO point positioning varies between about 0.5m and 2m (left).

The differences between the estimated kinematic positions and the true LEO orbit show large deviations twice-per-revolution (right).





# Main Results from Lab 2 (2)



The **a posteriori RMS of unit weight** ( $m_0$ ) of a combined Code- and Phase kinematic LEO point positioning varies between about 0.6 and 1.4mm (left). Due to the phase observations the quality of the kinematic positions has reached **sub-cm accuracies** for most epochs. Problematic phases twice-per-revolution, however, are still visible.

