

Addendum

# Addendum: A dispersive analysis of $\eta' \to \pi^+\pi^-\gamma$ and $\eta' \to \ell^+\ell^-\gamma$

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**Abstract** In this addendum to Ref. [1] we show that the mismatch between the  $\rho$ - $\omega$  mixing parameter  $\epsilon_{\rho\omega}$  as extracted from  $\eta' \to \pi^+\pi^-\gamma$  and  $e^+e^- \to \pi^+\pi^-$  can be resolved by including higher orders in the expansion in  $e^2$  in the description of the  $\eta' \to \pi^+\pi^-\gamma$  decay. We repeat the analysis in this extended framework and update the numerical results accordingly.

## Addendum to: Eur. Phys. J. C

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### 1 Extended formalism

Following the notation from Ref. [1] throughout, the spectrum for  $P \to \pi^+\pi^-\gamma$  can be expressed as

$$\frac{\mathrm{d}\Gamma(P \to \pi^+ \pi^- \gamma)}{\mathrm{d}s} = 16\pi \alpha \Gamma_0 |F_{\pi}^V(s)|^2 \left| P(s) \left( 1 + \Pi_{\pi}(s) \right) - \frac{e^2 F_{P\gamma\gamma}}{s} - \frac{g_{P\omega\gamma}}{g_{\omega\gamma}} \frac{\epsilon_{\rho\omega} - e^2 g_{\omega\gamma}^2}{M_{\omega}^2 - s - i M_{\omega} \Gamma_{\omega}} \right|^2,$$
(1.1)

generalizing Eq. (D.14) in Ref. [1] by the next order in the expansion in  $e^2$  (the sign convention is such that  $g_{P\omega\gamma} < 0$ ). The most important change, numerically, concerns  $\epsilon_{\rho\omega} \to \epsilon_{\rho\omega} - e^2 g_{\omega\gamma}^2$  in the numerator of the  $\omega$  propagator, corresponding to the photon contribution in  $\epsilon_{\rho\omega}$  as defined in resonance chiral perturbation theory [2–4]. In our formalism,  $\epsilon_{\rho\omega}$ , determined from a fit to the bare cross section for  $e^+e^- \to \pi^+\pi^-$ , does not include this VP effect, in line with the definition in Ref. [5] (numerically, it evaluates to  $e^2 g_{\omega\gamma}^2 = 0.34(1) \times 10^{-3}$ ). This shift removes the tension

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observed between  $\eta' \to \pi^+\pi^-\gamma$  and  $e^+e^- \to \pi^+\pi^-$  in Ref. [1].

The coefficients appearing in Eq. (3.9) of Ref. [1] are generalized according to Eq. (1.1):

$$\mathcal{A}_{2} = -\Gamma(\eta' \to \pi^{+}\pi^{-}\gamma) + 16\pi\alpha \int_{4M_{\pi}^{2}}^{M_{\eta'}^{2}} ds \, \Gamma_{0}|F_{\pi}^{V}(s)|^{2}$$

$$\times \left| \frac{g_{\eta'\omega\gamma}}{g_{\omega\gamma}} \frac{\epsilon_{\rho\omega} - e^{2}g_{\omega\gamma}^{2}}{M_{\omega}^{2} - s - iM_{\omega}\Gamma_{\omega}} + \frac{e^{2}F_{\eta'\gamma\gamma}}{s} \right|^{2},$$

$$\mathcal{A}_{1} = 32\pi\alpha \int_{4M_{\pi}^{2}}^{M_{\eta'}^{2}} ds \, \Gamma_{0}|F_{\pi}^{V}(s)|^{2} \operatorname{Re} \left[ P_{\text{ev}}(s) \left( 1 + \Pi_{\pi}^{*}(s) \right) \right]$$

$$\times \left( \frac{g_{\eta'\omega\gamma}}{g_{\omega\gamma}} \frac{e^{2}g_{\omega\gamma}^{2} - \epsilon_{\rho\omega}}{M_{\omega}^{2} - s - iM_{\omega}\Gamma_{\omega}} - \frac{e^{2}F_{\eta'\gamma\gamma}}{s} \right),$$

$$\mathcal{A}_{0} = 16\pi\alpha \int_{4M_{\pi}^{2}}^{M_{\eta'}^{2}} ds \, \Gamma_{0}|F_{\pi}^{V}(s)|^{2} P_{\text{ev}}^{2}(s) \left| 1 + \Pi_{\pi}(s) \right|^{2}.$$

$$(1.2)$$

In the following, we provide the updated numerical results when including the additional  $e^2$  effects as given in Eq. (1.1), implemented in the fit via Eq. (1.2).

## 2 Numerical results

The updated fit parameters are collected in Table 1, Fig. 1, and Table 2. The main difference to the results presented in Ref. [1] is that the shift  $\epsilon_{\rho\omega} \to \epsilon_{\rho\omega} - e^2 g_{\omega\gamma}^2$  removes the tension between  $e^+e^- \to \pi^+\pi^-$  and the  $\eta' \to \pi^+\pi^-\gamma$  spectrum, markedly improving the quality of the combined fit.



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**Table 1** Comparison of the fit outcome of the differential decay width in Eq. (1.1) to the BESIII  $\eta' \to \pi^+\pi^-\nu$  spectrum [6] of the binned maximum likelihood and minimum  $\chi^2$  strategies. The  $\chi^2$ /dof is 1.30 and 1.31, respectively, with the one of the Likelihood method extracted by means of the approximation described in App. C of Ref. [7]

Quantity	Likelihood	χ <sup>2</sup>
$A [ \text{GeV}^{-3} ]$	17.12(35)	17.09(32)
$\beta$ [ GeV $^{-2}$ ]	0.714(55)	0.723(45)
$\gamma$ [ GeV $^{-4}$ ]	-0.412(55)	-0.420(45)
$\epsilon_{\rho\omega} \times 10^3$	1.998(67)	1.997(54)
$M_{\omega}$ [ MeV]	782.99(33)	783.00(27)

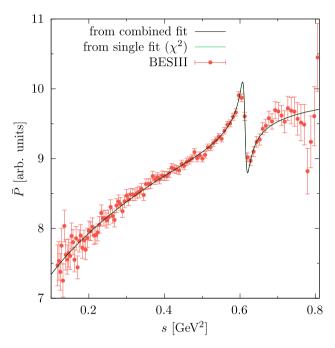


Fig. 1 Fit to the differential decay rate of  $\eta' \to \pi^+\pi^-\gamma$  (individually or combined with the VFF). To highlight potential differences in the  $\rho$ - $\omega$  region, we show the associated function  $\bar{P}$ , as defined in Eq. (3.11) of Ref. [1], compared to the experimental data from BESIII [6]. The two fits cannot be distinguished on this scale

The updated results for the TFF are shown in Fig. 2 and Table 3. In particular, the prediction for the slope parameter

$$b_{n'} = 1.431(23) \,\text{GeV}^{-2}$$
 (2.1)

is reduced by about  $1\sigma$ , which traces back not to the change in  $\epsilon_{\rho\omega}$  (which is marginal given the fact that the fit is dominated by  $e^+e^-\to\pi^+\pi^-$ ), but to a stronger curvature in the polynomial P(s) (the coefficient  $\gamma$  of the quadratic term increases by a factor 3).

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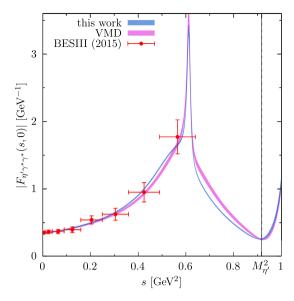
**Table 2** Combined fit to several pion VFF data sets (BaBar, KLOE, CMD-2, SND) and  $\eta' \to \pi^+\pi^-\gamma$  spectrum (BESIII) with overall  $\chi^2/\text{dof} = 1.46$ . In the row for KLOE, the three values for  $M_\omega$  refer to

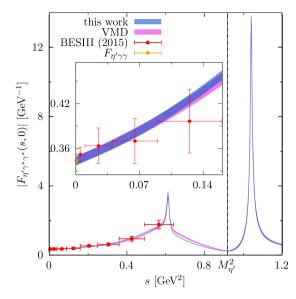
the combinations of the global KLOE  $\omega$  mass and the corresponding mass shifts of the three underlying data sets from 2008, 2010, 2012, respectively

	$\chi^2/{ m dof}$	$M_{\omega} \; [ \mathrm{MeV}]$	$A [GeV^{-3}]$	$\beta  [ \mathrm{GeV}^{-2}]$	$\gamma  [ \mathrm{GeV}^{-4}]$	$\alpha_{\pi} \times 10^2 \; [ \mathrm{GeV^{-2}}]$	$\epsilon_{\rho\omega} \times 10^3$
BaBar	1.26	781.875(82)				)	)
		781.65(12)					
KLOE	1.61	782.10(17)					
		781.84(27)				5.74(14)	2.007(10)
CMD-2	2.18	782.131(68)					
SND	2.16	781.457(97)					
BESIII	1.31	783.00(28)	17.10(32)	0.720(46)	-0.418(46)	J	J



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**Fig. 2** Determination of the  $\eta'$  TFF in comparison to data from BESIII [8] (statistical and systematic errors added in quadrature) scaled with  $F_{\eta'\gamma\gamma}$  and the VMD model from Ref. [1] for the  $\phi$  resonance; for

the kinematic range accessible in  $\eta'$  decays (left) and a larger time-like region including the  $\phi$  resonance with inset magnifying the low-s region (right)

Table 3 Contributions from the various components of the TFF to the sum rules of the normalization and the slope parameter

	$(I=1)_{\epsilon_{\rho\omega}=0}$	$\Delta(I=1)_{\epsilon_{\rho\omega}\neq 0}$	$(I=0)_{\epsilon_{\rho\omega}=0}^{\omega}$	$\Delta(I=0)^{\omega}_{\epsilon_{\rho\omega}\neq 0}$	$(I=0)^{\phi}$	Total
Norm [%]	69.18(86)	-0.1388(19)	7.06(22)	-0.1397(47)	15.85(61)	91.9(1.1)
$b_{\eta'}$ [ ${ m GeV}^{-2}$ ]	1.160(23)	0	0.1176(32)	0	0.1526(53)	1.431(23)

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