



# Rejuvenating the ocean: mean ocean radiocarbon, CO<sub>2</sub> release, and radiocarbon budget closure across the last deglaciation

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**Abstract.** Radiocarbon is a tracer that provides unique insights into the ocean's ability to sequester CO<sub>2</sub> from the atmosphere. While spatial patterns of radiocarbon in the ocean interior can indicate the vectors and timescales for carbon transport through the ocean, estimates of the global average ocean–atmosphere radiocarbon age offset (B–Atm) place constraints on the closure of the global carbon cycle. Here, we apply a Bayesian interpolation method to compiled B–Atm data to generate global interpolated fields and mean ocean B–Atm estimates for a suite of time slices across the last deglaciation. The compiled data and interpolations confirm a stepwise and spatially heterogeneous “rejuvenation” of the ocean, suggesting that carbon was released to the atmosphere through two swings of a “ventilation seesaw” operating between the North Atlantic and both the Southern Ocean and the North Pacific. Sensitivity tests using the Bern3D model of intermediate complexity demonstrate that a portion of the reconstructed deglacial B–Atm changes may reflect “phase-attenuation” biases that are unrelated to ocean ventilation and that arise from independent atmospheric radiocarbon dynamics instead. A deglacial minimum in B–Atm offsets during the Bølling–Allerød could partly reflect such a bias. However, the sensitivity tests further demonstrate that when correcting for such biases, ocean “ventilation” could still account for at least one-third of deglacial atmospheric

CO<sub>2</sub> rise. This contribution to CO<sub>2</sub> rise appears to have continued through the Younger Dryas, though much of the impact was likely achieved by the end of the Bølling–Allerød, indicating a key role for marine carbon cycle adjustment early in the deglacial process. Our global average B–Atm estimates place further new constraints on the long-standing mystery of global radiocarbon budget closure across the last deglaciation and suggest that glacial radiocarbon production levels are likely underestimated on average by existing reconstructions.

## 1 Introduction

The evolving spatial distribution of marine radiocarbon provides unique constraints on air–sea gas exchange at the sea surface, and the ocean dynamics that convey surface ocean–atmosphere *p*CO<sub>2</sub> differences into the ocean interior (Koeve et al., 2015). These transports of carbon in turn set constraints on the ability of the ocean to store CO<sub>2</sub>, away from the atmosphere, in the form of either “disequilibrium” or “respired” dissolved inorganic carbon (Galbraith and Skinner, 2020). A slower overturning circulation enhances the accumulation of respired carbon in the ocean interior, in parallel with greater radiocarbon decay (Menviel et al., 2017;

Eggleson and Galbraith, 2018; Tschumi et al., 2011). Reduced air–sea exchange impedes the release of respired carbon from upwelled water to the atmosphere while also impeding radiocarbon input from the atmosphere to the ocean (e.g. Khatiwala et al., 2019; Stocker and Wright, 1996). Both processes enhance carbon sequestration in the ocean (Galbraith and Skinner, 2020; Eggleson and Galbraith, 2018) while depleting the ocean’s average radiocarbon activity relative to the atmosphere (Skinner and Bard, 2022). Therefore, estimates of the global average ocean–atmosphere radiocarbon age offset (B-Atm) may provide powerful constraints on the evolving carbon exchange between the ocean and atmosphere (Siegenthaler et al., 1980). This in turn has implications for the closure of the radiocarbon budget of the atmosphere, given that the atmospheric radiocarbon inventory is largely set by the radiocarbon production rate and ocean–atmosphere radiocarbon exchange (Siegenthaler, 1989).

Few estimates of global average B-Atm prior to the instrumental record currently exist. This largely relates to the difficulty of estimating the “volumetric representativity” of sparse data points, in order to weight their contribution to the global mean. Using a simple unweighted arithmetic mean of existing data from the Last Glacial Maximum (LGM,  $\sim 20$  ka), Sarnthein et al. (2013) demonstrated an increase in the mean ocean–atmosphere radiocarbon age offset (B-Atm) by  $\sim 600$   $^{14}\text{C}$  yr. A subsequent study instead made use of a Bayesian interpolation method, over a  $4^\circ$  resolution global grid, to generate a global field of B-Atm offsets from which to derive a volume-weighted global average (Skinner et al., 2017). This alternative approach again showed that the glacial ocean was significantly more radiocarbon depleted relative to the atmosphere, especially in the deep Atlantic and the Southern Ocean (Sikes et al., 2016a; Skinner et al., 2021, 2010; Gottschalk et al., 2020). This study inferred a mean ocean B-Atm increase of  $689 \pm 53$   $^{14}\text{C}$  yr as compared with the pre-industrial ocean. Using yet another approach, whereby individual data points were weighted according to their location (and their current offset from modern regional and/or basin averages), Rafter et al. (2022) have again confirmed an “ageing” of the deep ocean at the LGM. The study of Rafter et al. (2022) was restricted to the deep ocean  $> \sim 1000$  m water depth and therefore did not yield global mean estimates. Nevertheless, the  $\sim 1050 \pm 150$   $^{14}\text{C}$  yr increase inferred by Rafter et al. (2022) for the deep ocean (approximately equivalent to  $> 2$  km water depth) almost certainly implies a significant increase in the global mean B-Atm at the LGM. Note that Rafter et al. (2022) incorrectly identified the global mean estimate of Skinner et al. (2017) as referring to the deep ocean  $> 3$  km water depth, and that consequently the estimates of Rafter et al. (2022) and Skinner et al. (2017) may be more consistent with each other than initially apparent.

Overall, a consistent picture has therefore emerged of a relatively “aged” dissolved inorganic carbon (DIC) pool in the LGM ocean. A recent review of the existing data has fur-

ther demonstrated coherent patterns in marine radiocarbon evolution since the LGM, across the last deglaciation (Skinner and Bard, 2022). Several key observations have so far emerged from the compiled data: (1) at the onset of deglaciation, during Heinrich Stadial 1 (HS1,  $\sim 17.5$ – $14.7$  ka), and to a lesser extent the Younger Dryas (YD,  $\sim 12.9$ – $11.7$  ka), B-Atm offsets increased in the deep and intermediate North Atlantic, while they decreased in the Southern Ocean and North Pacific (particularly in the intermediate North Pacific, suggesting a “ventilation seesaw”) (e.g. Broecker, 1998; Skinner et al., 2013, 2014; Menviel et al., 2018; Ahagon et al., 2003; Menviel et al., 2014; Freeman et al., 2015; Max et al., 2014; Okazaki et al., 2010; Walczak et al., 2020). (2) B-Atm offsets decreased significantly throughout the global ocean at the Bølling–Allerød (BA,  $\sim 14.7$ – $12.9$  ka), in places suggesting an overshoot to values younger than pre-industrial (e.g. Barker et al., 2010). (3) Surface reservoir ages in the high latitudes of the North Atlantic and the Southern Ocean tend to track B-Atm changes at intermediate depths (Skinner et al., 2019), suggesting limited changes in transport  $< 2$  km water depth (and therefore a significant contribution from restricted gas-exchange efficiency instead). This contrasts with the deep ocean  $> 2$  km, where stronger evidence for transport changes is apparent (Skinner et al., 2021; Marchal and Zhao, 2021a). Encouragingly, similar observations have recently been reported based on a complex approach to data screening, averaging, and weighting (Rafter et al., 2022). The latter study emphasized the influence of transport changes on B-Atm offsets: enhanced mid-depth ventilation was inferred in the North Pacific at the LGM; and reduced transport rates were inferred in the deep Pacific at the LGM, and in the deep North Atlantic during HS1 and the YD.

Several questions arise in light of the existing body of deglacial marine radiocarbon data. First, what do the existing data imply for *global average* (i.e. rather than regional or deep ocean) B-Atm offsets across the last deglaciation? Second, can we determine the degree to which observed changes in the mean ocean B-Atm offset reflect changes in “ventilation” specifically (i.e. circulation rates *and* gas exchange), as opposed to, for example, independent changes in atmospheric radiocarbon (Franke et al., 2008)? Third, can we quantify the likely impact of past radiocarbon ventilation effects (regardless of their origin) on ocean–atmosphere carbon partitioning and, therefore, atmospheric  $\text{CO}_2$  (Skinner and Bard, 2022)? Finally, can the temporal evolution of mean ocean B-Atm offsets be reconciled with past atmospheric radiocarbon activities and past radiocarbon production rates; and do they resolve the long-standing “mystery” of deglacial radiocarbon budget closure (Broecker and Barker, 2007; Köhler et al., 2022, 2006)?

Here, we attempt to address each of these questions in turn. We update a previously deployed Bayesian interpolation technique (Skinner et al., 2017), which we apply to an updated compilation of radiocarbon data spanning the last deglaciation, to derive estimates of global average B-Atm

offsets. These estimates are interpreted against a suite of new sensitivity tests and transient simulations conducted using the Bern3D model of intermediate complexity (Müller et al., 2006; Ritz et al., 2011; Roth et al., 2014). The new simulations are aimed at constraining the potential magnitude of “ventilation” versus “non-ventilation” effects in global mean B-Atm offsets, as well as associated changes in ocean–atmosphere CO<sub>2</sub> partitioning.

## 2 Methods

### 2.1 Radiocarbon data

Our data compilation is based on the work of Zhao et al. (2018) as presented in Skinner and Bard (2022). The compilation has been reconciled with a similar collection of data that has recently emerged (Rafter et al., 2022). A key difference in our compilation is that we include warm water (near surface) coral data and surface reservoir age data (e.g. Skinner et al., 2019) where direct estimates exist. Our compilation adopts revised chronologies and radiocarbon age offsets consistent with the latest *Intcal20* (Reimer et al., 2020) and *Marine20* (Heaton et al., 2020) calibration curves. To update the chronologies of the collected time series, where the only age control derives from planktonic radiocarbon dates, we have used the *Marine20* calibration curve (Heaton et al., 2020), in conjunction with modern local deviations from the global average surface reservoir age ( $\Delta R$  values  $\sim$  local R-age minus 550 <sup>14</sup>C yr) (Heaton et al., 2020), to generate “reservoir-age” corrected calibrated calendar ages, using the R package *Bchron* (Parnell et al., 2008). While this approach has the advantage of remaining as faithful as possible to the original published age scale for each record, it also has the notable drawback of neglecting potential changes in  $\Delta R$ , which are known to have occurred in mid-to-high latitudes (e.g. Skinner et al., 2019) and have been directly estimated in some of the studies included in the compilation. For records updated in this way, where the density and/or down-core variability of planktonic radiocarbon dates in a sediment core is sufficient to result in age reversals (i.e. older dates stratigraphically above younger dates), we have generated new sediment depth-age models using the MCMC approach of *Bchron*, from which maximum likelihood calendar age estimates and 95 % uncertainty intervals ( $\sim 2\sigma$ ) are obtained for each sample depth. For time series that have made use of U-series dating (e.g. Robinson et al., 2005; Burke and Robinson, 2012; Chen et al., 2015; Hines et al., 2015), or stratigraphic alignments (including stratigraphic alignments that have provided time-varying surface reservoir age corrections; e.g. Skinner et al., 2010; Gottschalk et al., 2020; Austin et al., 2011; Peck et al., 2006), as well as a priori assignment of down-core changes in reservoir ages (e.g. Ronge et al., 2020), we adopt the published calendar ages, as these do not depend on the atmospheric radiocarbon calibration curve.

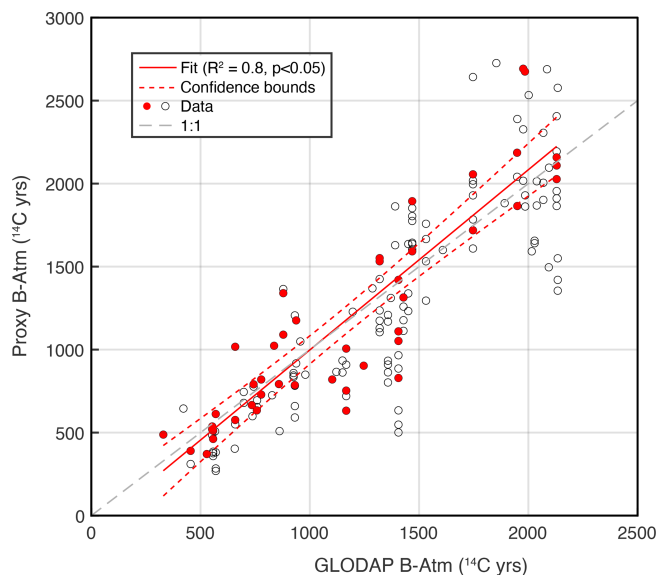
For all compiled data, we have recalculated radiocarbon reservoir age offsets relative to the *Intcal20* (Reimer et al., 2020) atmospheric reference curve using the *radcal* function in the R package *ResAge* (Soulet, 2015). This approach yields 95 % uncertainty limits for the resulting radiocarbon age offsets based on the joint calendar age and radiocarbon age uncertainties. The resulting probabilistic B-Atm estimates differ from simple differences between marine and atmospheric radiocarbon values to some extent (Rafter et al., 2022). Estimates of the relative (radiocarbon) isotopic enrichment of the ocean versus the atmosphere are expressed as radiocarbon age offsets between the ocean and the contemporary atmosphere (i.e. in <sup>14</sup>C yr), which are equivalent to  $-8033 \times \ln(\alpha_{\text{O-Atm}})$ , where  $\alpha_{\text{O-Atm}}$  represents the ratio of marine radiocarbon activity versus the contemporary atmospheric radiocarbon activity at time  $T$ , i.e.  $\text{Fm}_{\text{O}}^T / \text{Fm}_{\text{Atm}}^T$  (Soulet et al., 2016). We do not refer to offsets between marine and atmospheric  $\Delta^{14}\text{C}$  (i.e.  $\Delta^{14}\text{C}_{\text{O}}$  minus  $\Delta^{14}\text{C}_{\text{Atm}}$ ), since this metric does not relate to the isotopic depletion of the marine reservoir relative to the atmosphere in a predictable way without knowledge of the absolute atmospheric and marine radiocarbon activities (e.g. Cook and Keigwin, 2015).

Data quality flags are assigned to anomalous values or datasets, including those that yield negative ventilation ages, negative reservoir ages, or negative benthic–planktonic offsets (B–P), or that exhibit significant down-core age reversals (e.g. Rose et al., 2010). In addition, datasets that deviate significantly from regional trends, for example due to alternative age models, are also flagged. The latter category includes records based on “plateau tuned” (PT) chronologies (Sarnthein et al., 2015, 2007, 2020, 2013; Ausín et al., 2021) that differ significantly from alternative reconstructions using the same data, or from reconstructions at proximal sites (Skinner and Bard, 2022). Such differences may relate to identifiable drawbacks of the “plateau tuning” methodology (Bard and Heaton, 2021). Datasets that are interpreted as being influenced by hiatuses are also flagged (e.g. Ronge et al., 2020, 2019), as are data that are interpreted as having been influenced by localized geological or sedimentary carbon sources (Rafter et al., 2019, 2018; Lindsay et al., 2016, 2015; Ronge et al., 2016; Bova et al., 2018; Stott et al., 2009; Marchitto et al., 2007). A similar selection was effectively made in the recent study of Rafter et al. (2022), where all data from water depths less than  $\sim 1000$  m were omitted from consideration, and where individual data points were removed from binned regional water-depth groupings based on their deviation (by  $> 3\sigma$ ) from the group’s weighted arithmetic mean.

Although the data compiled in this study have all been published, including the vast majority in a previous compilation (Zhao et al., 2018), it remains possible that some data have been affected by as yet undetected biases arising from bioturbation (Bard et al., 1987; Dolman et al., 2021) and/or diagenesis (Wycech et al., 2016). Unless obviously anomalous signals are produced (Missiaen et al., 2020), such bi-

ases can be extremely difficult to positively identify but could have profound impacts (Lougheed et al., 2020; Stott, 2023). Although sites with very low sedimentation rates may be particularly prone to bioturbation effects (Bard et al., 1987; Dolman et al., 2021), the avoidance of all data from such locations might also introduce a sampling bias due to their prevalence at greater water depths. This presents a significant challenge. In order to limit the potential impacts of anomalies caused by bioturbation (Dolman et al., 2021), we apply data flags to all sediment cores with average accumulation rates  $< 2 \text{ cm kyr}^{-1}$  (rounded to the nearest digit). Of the remaining data, 75 % derive from sites with accumulation rates  $> 10 \text{ cm kyr}^{-1}$  and 95 % derive from sites with accumulation rates  $> 4 \text{ cm kyr}^{-1}$ . A comparison of late Holocene data ( $< 6 \text{ ka BP}$ ) with pre-industrial observations (Key et al., 2004) serves to demonstrate the general fidelity of B-Atm offsets derived from fossil substrates as compared with modern seawater values ( $R^2 = 0.81$ ; see Fig. 1). However, this comparison also illustrates the significant scatter that exists in the fossil data (RMSE  $\sim 273$  years; or an “inverse prediction interval”, which is more appropriate if considering a “calibration” of fossil data,  $\sim \text{RMSE} \times 1.96 = 535$  years; McClelland et al., 2021). A similar result has recently been demonstrated for sub-modern B-Atm offsets by Rafter et al. (2022). The observed scatter is expected to derive primarily from sedimentary and sampling issues, rather than analytical uncertainty, and calls for caution when interpreting the subtler features that emerge from the collected data. Accordingly, an underlying premise of our work (as for previous compilations) is that the temporal trends and spatial patterns that emerge from a large body of independent data are far less likely to be dominated by sedimentary or diagenetic biases that can more readily influence individual sediment cores (Ronge et al., 2019; Stott, 2020).

In order to assess the implications of applying the data flags described above, three alternative interpolations were performed: (1) the baseline interpolation, where all flagged data were excluded (Table 1, “Baseline”); (2) as for the baseline but with low sedimentation rate sites ( $< 2 \text{ cm kyr}^{-1}$ ) retained (Table 1, “Low sedimentation”); and (3) as for the baseline but retaining only high sedimentation rates  $> 10 \text{ cm kyr}^{-1}$  (mainly affecting sites in the 2500–4000 m water depth range). In keeping with a similar assessment made by Rafter et al. (2022), we find that the inclusion of very low sedimentation rate sites ( $< 2 \text{ cm kyr}^{-1}$ ) has no major effect on our findings, particularly for the global averages (Table 1). The same is true when only retaining sites with very high sedimentation rates  $> 10 \text{ cm kyr}^{-1}$  (Figs. S2 and S3 in the Supplement), though global averages are biased to slightly lower values (Stott, 2023), particularly at the LGM (Table 1, “High sedimentation”). This apparent lack of sensitivity arises here due to the averaging of a relatively large number of data points (thereby reducing the influence of statistical outliers); however, it should be borne in mind that the inclusion or exclusion of individual data points from



**Figure 1.** Comparison of modern seawater (corrected for the effects of atomic bomb testing) B-Atm radiocarbon age offsets (Key et al., 2004) versus proxy-based B-Atm reconstructions using material deposited during the past  $\sim 6000$  years (open black circles) and the past  $\sim 1500$  years (filled red circles). Dashed grey line indicates the 1 : 1 trend. The linear fit to the data is indicated by the red line (dashed red lines show 95 % confidence limits), with the following equation:  $y = (1.1 \pm 0.1)x + (89 \pm 98)$ ,  $R^2 = 0.81$ ,  $p = 1.77 \times 10^{-17}$ .

sparsely sampled regions (e.g. for the Indian Ocean) can influence the detail of inferred spatial patterns.

## 2.2 Time slices

We isolate data associated with six key time slices based on their calendar ages and associated uncertainties: the LGM (19–21.8 ka BP), Heinrich Stadial 1 (15–17.5 ka BP), the Bølling-Allerød (12.8–14.8 ka BP), the Younger Dryas (YD, 11.8–12.7 ka BP), the early Holocene (EHOL, 9–11 ka BP), and the late Holocene (HOL,  $< 6 \text{ ka BP}$ ). Note that the boundaries of these time intervals are informed by, but do not correspond precisely to, their respective ice-core chronozone definitions (Rasmussen et al., 2014). This is due to the generally lower resolution of the marine records and the need to be pragmatic in avoiding, as much as possible, misattributing data to an adjacent chronozone. Where calendar age uncertainties result in an ambiguous assignation of a single time slice, no time slice is assigned.

## 2.3 Radiocarbon “ventilation” metrics

The term “ventilation” is defined here as comprising the processes that influence physical and chemical property exchanges between the ocean interior and the atmosphere (Skinner and Bard, 2022). Accordingly, the term “(radiocar-



	Baseline ( $n = 2579$ , avg. $z = 1407\text{ m}$ )						Low sedimentation ( $n = 2688$ , avg. $z = 1472\text{ m}$ )						High sedimentation ( $n = 2036$ , avg. $z = 1244\text{ m}$ )										
Time slice	Mean ocean B-Atm			$\Delta(\text{B-Atm})$			err	Mean ocean B-Atm			$\Delta(\text{B-Atm})$			err	Mean ocean B-Atm			$\Delta(\text{B-Atm})$			err		
	$^{14}\text{C}_{\text{yr}}$	$^{14}\text{C}_{\text{yr}}$	$^{14}\text{C}_{\text{yr}}$	B-Atm	$^{14}\text{C}_{\text{yr}}$	$^{14}\text{C}_{\text{yr}}$		B-Atm	$^{14}\text{C}_{\text{yr}}$	$^{14}\text{C}_{\text{yr}}$	$^{14}\text{C}_{\text{yr}}$	B-Atm	$^{14}\text{C}_{\text{yr}}$		$^{14}\text{C}_{\text{yr}}$	$^{14}\text{C}_{\text{yr}}$	B-Atm	$^{14}\text{C}_{\text{yr}}$	$^{14}\text{C}_{\text{yr}}$	$^{14}\text{C}_{\text{yr}}$			
HOL	1335	-24	22	1328	-31	24		1350	-9	20		1350	-9	20		1350	-9	20		1350	-9	20	
EHL	1359	0	26	1360	1	39		1329	-30	31		1329	-30	31		1329	-30	31		1329	-30	31	
YD	1484	125	37	1515	156	37		1478	119	41		1478	119	41		1478	119	41		1478	119	41	
BA	1435	76	32	1487	128	37		1432	73	32		1432	73	32		1432	73	32		1432	73	32	
HS	1877	518	40	1981	621	55		1793	434	55		1793	434	55		1793	434	55		1793	434	55	
LGM	2123	764	32	2185	825	36		1993	634	46		1993	634	46		1993	634	46		1993	634	46	

In summary, B-Atm should not be used to directly imply transit times or ideal ages, which are not equivalent to “radiocarbon ventilation ages”. Furthermore, metrics such as B-Atm offsets are influenced by a wider variety of processes than just ventilation (i.e. more than just gas exchange and transport), including, for example, atmospheric radiocarbon production changes (Heaton et al., 2021). Therefore, B-Atm offsets should only be used to refer to ventilation (i.e. gas exchange and transport) with caution, as we explore in this study.

We use the interpolation method described in Skinner et al. (2017), updated to weight input data according to their reported uncertainties. The method interpolates the observed ocean–atmosphere radiocarbon age offset anomalies onto the grid of an ocean model with 24 vertical levels and  $2^\circ \times 2^\circ$  horizontal resolution. The anomalies are defined relative to the modern GLODAP radiocarbon ages (Key et al., 2004). The interpolating function is constructed using a weighted superposition of radial basis functions centred at the sediment-core locations. These basis functions are built here using an exponential basic function,  $\varphi(r) = e^{-\varepsilon r}$ , with shape parameter,  $\varepsilon$ . The distance  $r$  is defined as the time for tracer impulses to spread out from the core locations to the rest of the grid, based on the modern transport and density stratification. The basis-function weights, the shape parameter, and two hyper-parameters that scale the relative size of the precisions in the prior and likelihood functions are inferred using a three-level Bayesian procedure. (Further details are given in Skinner et al., 2017.) The interpolation process is repeated independently for each time slice so that the data from one

time slice do not influence the interpolation for the other time slices.

This interpolation method is somewhat conservative because it is based on the distribution of relative tracer transport and diffusion timescales in the modern ocean. This choice is premised on the need to obtain an interpolated solution in three dimensions that is physically sensible and informed by large-scale oceanographic requirements, including vertical stratification, basin margins/topographic boundaries, and dominant transport pathways (e.g. the Antarctic Circumpolar Current, Deep Western Boundary currents, etc.). It is important to note that the interpolation method does not impose the modern transport on the resulting interpolated fields, nor does it assume strict adherence to the modern density field. Rather, it represents a physically guided and data-constrained anomaly relative to the modern state. The inferred anomalies relax to zero in the absence of data constraints, though some locations in the ocean can have far reaching influence. This method represents a counterpoint to the approach recently taken by Rafter et al. (2022), for example, where spatial interpolations were illustrated for data projected onto 2D sections. In their study, averages were also produced for time series grouped and weighted according to their position within the modern density- and B-Atm distributions (i.e. locations closest to the modern mean for a given basin and density class were weighted to dominate the mean in the past). With this approach, the spatial distribution of density classes is assumed invariant over time, despite the inference of circulation changes. This mirrors to some extent our use of a modern transport field as a guide to the volumetric representativity of individual sample locations. Interpolating a sparsely sampled 3D field is a difficult problem, and a diversity of approaches is surely useful; however, it should be noted that Rafter et al. (2022) did not consider the upper  $\sim 1000$  m of the ocean and did not produce or discuss global averages.

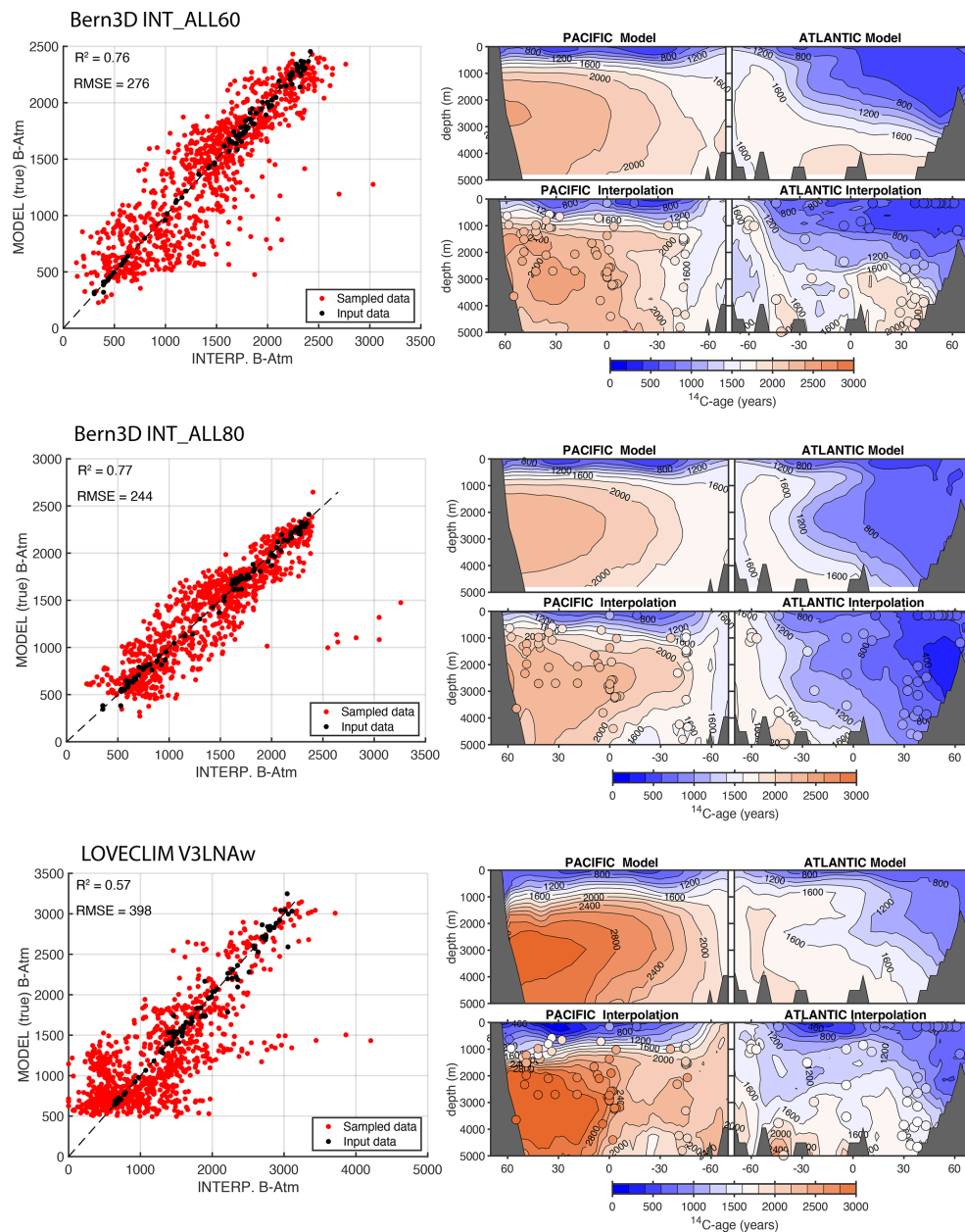
In order to assess the accuracy of the interpolation method of Skinner et al. (2017), and its ability to represent different circulation states, we apply it to a range of modelled radiocarbon fields. For this we use new outputs from the Bern3D model (see below) and pre-existing outputs from the LOVECLIM model (Menviel et al., 2017, 2018). We extract data from the modelled fields at the same locations as the available proxy data (e.g. for the LGM) and attempt to reproduce the modelled global field using the interpolation method. As shown in Fig. 2, the interpolation method is generally successful in reproducing the simulated radiocarbon fields for altered circulation states. The interpolation does best for circulation states characterized by altered wind, diffusivity, and gas exchange, yielding average errors for 1000 randomly selected individual grid cell estimates of RMSE  $\sim 276$  years (INT\_ALL60; this study), RMSE  $\sim 244$  years (INT\_ALL80; this study), RMSE  $\sim 398$  years (V3LNAw; Menviel et al., 2017), and RMSE  $\sim 511$  years (V3LNAwSOWSHWw; Menviel et al., 2017) (see Figs. 2 and 3). The interpolation performs slightly less well for extremely different circulation

states, e.g. yielding RMSE  $\sim 457$  years for the HS1 simulation of Menviel et al. (2018) and RMSE  $\sim 577$  years for the collapsed Atlantic Meridional Overturning Circulation (AMOC) simulation of Menviel et al. (2017), both of which yielded reversed radiocarbon gradients in the Atlantic as compared with today (Fig. 3). Encouragingly, the interpolation method reproduces the modelled fields more accurately than model simulations of the LGM are typically able to match observations (Muglia et al., 2018; Menviel et al., 2017).

In the present study, the interpolation is primarily intended for calculating global average B-Atm offsets using a relatively sparse dataset ( $N = 124\text{--}476$ ). It is therefore notable that the interpolation reproduces modelled global averages even more accurately than the spatial distributions, with RMSE  $\sim 53$  years (Fig. 4, circles). In contrast, simple arithmetic means of B-Atm values drawn from a sub-set of locations in the modelled fields consistently underestimate the true global mean by  $\sim 274$  years on average (Fig. 4, triangles). This comparison underlines the importance of applying a spatially resolved volumetric weighting to the observations for the derivation of accurate global averages. While our interpolation approach leaves room for improvement, it complements alternative data-constrained modelling approaches that have been applied to, for example, the LGM and the YD (Pöppelmeier et al., 2023), and it represents a first step towards addressing the problem of deriving 3D global fields, and appropriately weighted global averages, from sparsely sampled proxy data, without (as yet) the use of forward or inverse models.

## 2.5 Modelling

The intermediate complexity Bern3D v2.0 Earth system model (Ritz et al., 2011; Roth et al., 2014) has been used to explore a series of idealized scenarios that aim to probe the parallel sensitivities of atmospheric CO<sub>2</sub> and marine radiocarbon, subject to changes in global vertical diffusivity, Southern Ocean winds, and/or Southern Ocean gas-exchange efficiency. The implementation of radiocarbon in the Bern3D model is described elsewhere (Müller et al., 2006; Müller et al., 2008; Dinauer et al., 2020), and the simulations presented here extend those performed by Jeltsch-Thömmes et al. (2019) and Dinauer et al. (2020), as well as those performed using an earlier version of the Bern3D model by Tschumi et al. (2011). The applied version of the Bern3D model comprises a single-layer energy-moisture balance atmosphere with a thermodynamic sea-ice component (Ritz et al., 2011), coupled to a 3D geostrophic-frictional balance ocean model (Edwards et al., 1998; Müller et al., 2006), and a four-box representation of the land biosphere that simulates the dilution of an atmospheric isotopic perturbation by the land biosphere but here does not address changes in land carbon stocks (Siegenthaler and Oeschger, 1987). Furthermore, marine sediments were not implemented in the model set-up

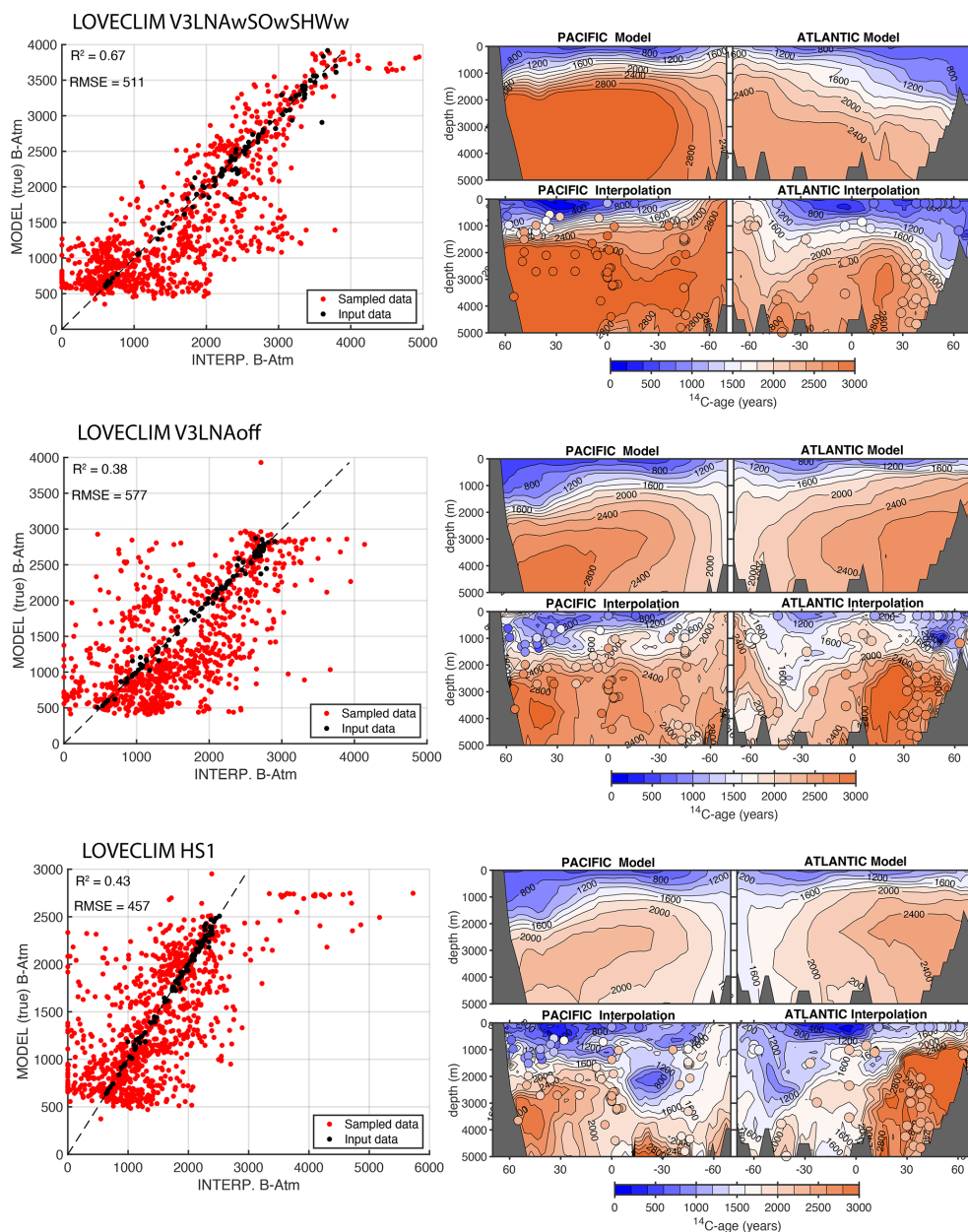


**Figure 2.** Testing the interpolation method using numerical model outputs. Interpolations for the Bern3D and LOVECLIM models, from top to bottom: INT\_ALL60 (this study), INT\_ALL80 (this study), and V3LNAw (Menviel et al., 2017). Cross plots at left show “true” versus interpolated B-Atm offsets (black circles for input data, red circles for 1000 randomly sampled locations in the ocean interior), with  $R^2$  and RMSE shown. Contoured panels show zonally averaged B-Atm offsets (i.e.  $^{14}\text{C}$ -age relative to the atmosphere), for the Pacific (left) and Atlantic (right), for both the model output (upper panels) and the interpolated reconstructions (lower panels), with input data used for the interpolations indicated by the filled circles.

used here, as we sought to *isolate* the impacts of ventilation processes (i.e. ocean gas exchange and transport).

A first set of sensitivity tests (PI-), performed under “pre-industrial” conditions, and run out for 10 000 years with a fully interactive climate-carbon cycle system (minus sediments), consisted of step changes in (1) global vertical diffusivity in the ocean ( $K_v$ ), (2) Southern Ocean wind stress

(SW), (3) Southern Ocean gas-exchange efficiency or piston velocity (SG), or (4) a combination of  $K_v$ , SW, and SG (ALL). In each case the relevant parameter was reduced by 20 %, 40 %, 60 %, or 80 % relative to its baseline value, and impacts were evaluated after 2000 years. (Impacts after 500 and 5000 years are illustrated in Fig. S1.)

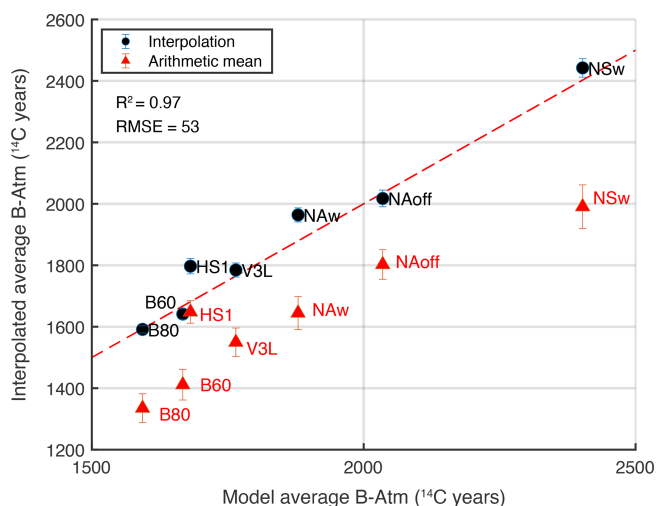


**Figure 3.** Testing the interpolation method using numerical model outputs, as for Fig. 2. Interpolations for simulations using the LOVECLIM model, from top to bottom: V3LNAwSOwSHWw (Menviel et al., 2017), V3LNAoff (Menviel et al., 2017), and HS1 (Menviel et al., 2018). Cross plots at left show “true” versus interpolated B-Atm offsets (black circles for input data, red circles for 1000 randomly sampled locations in the ocean interior), with  $R^2$  and RMSE shown. Contoured panels show zonally averaged B-Atm offsets (i.e.  $^{14}\text{C}$ -age relative to the atmosphere), for the Pacific (left) and Atlantic (right), for both the model output (upper panels) and the interpolated reconstructions (lower panels), with input data used for the interpolations indicated by the filled circles.

In an additional set of simulations, we employed a range of glacial/deglacial scenarios, similar to those of Jeltsch-Thömmes et al. (2019), where each is again defined by idealized adjustments to vertical diffusivity, wind stress, Southern Ocean gas exchange, or a combination of these, as described for the above PI- sensitivity tests. For these simulations, the model was “spun up” into equilibrium over 35 000

model years under pre-industrial (1700 CE) boundary conditions. The model was then ramped linearly over 5000 years from the resulting PI equilibrium state into “glacial” boundary conditions (i.e. representative of 50 ka BP) for ice sheet albedo, greenhouse gas radiative forcing, and insolation, as well as prescribed values for diffusivity, wind stress, and Southern Ocean gas-exchange efficiency as described above.





**Figure 4.** Comparison of global average B-Atm offsets based on interpolated fields (circles, with estimated interpolation error), and geometric averages of the input data used for the interpolations (triangles, with standard error), versus the true values for the model simulations shown in Figs. 2 and 3 plus an additional simulation (NAw) from (Menviel et al., 2017). The  $R^2$  correlation coefficient, RMSE, and 1 : 1 line (dashed red line) are indicated.

After this, ice-sheet albedo, greenhouse gas radiative forcing, and insolation were varied based on observations since 50 ka BP, with diffusivity, Southern Ocean wind stress, and Southern Ocean gas-exchange efficiency remaining constant for 32 000 years and then relaxing linearly back to PI control values from 18 to 14.6 ka BP, after which PI control values were maintained.

In one set of “glacial/deglacial” model simulations (INT-), atmospheric radiocarbon was prescribed from 50 ka BP based on the *Intcal20* reference curve (Reimer et al., 2020). Radiocarbon concentrations in the model were therefore required to be consistent with the observed atmospheric radiocarbon concentration, via changes in the global radiocarbon inventory (which are equivalent to ad hoc changes in radiocarbon production). In a second set of simulations (FIX-) atmospheric radiocarbon activity was held constant at 140‰, while in a third set of simulations (CONST-) atmospheric radiocarbon production rates were held constant at the PI value. In each case a control simulation was performed with evolving boundary conditions as described above. A series of additional simulations were then performed with altered global vertical diffusivity ( $K_v$ ), Southern Ocean wind stress (SW), Southern Ocean gas-exchange efficiency (SG), or a combination of these (see Table S1 in the Supplement). The goal of these transient simulations is not to reproduce the last deglaciation but to assess the sensitivity of both marine radiocarbon and atmospheric  $\text{CO}_2$  to a variety of changes in ocean ventilation (in terms of both the type of forcing and its magnitude), specifically where concurrent radiocarbon pro-

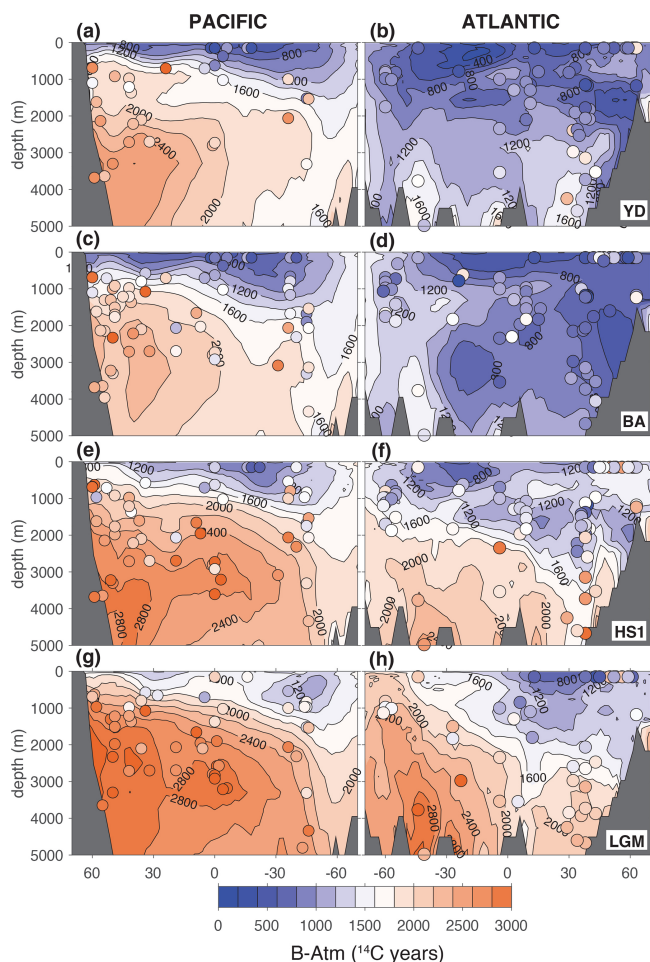
duction rate changes are either minimized (CONST-/FIX-) or maximized (INT-).

### 3 Results

Figure 5 illustrates zonally averaged radiocarbon age offsets (B-Atm) for the LGM, HS1, BA, and YD, in the Atlantic and Pacific basins, based on global interpolations of the compiled data (using the “baseline” data flags). Differences between each successive time-slice interpolation are shown in Fig. 6. Zonal averages and time-slice offsets for the Indian Ocean are shown in Fig. 7. Zonal averages for the EHOL and HOL, which differ only slightly from the modern field (Key et al., 2004), are shown in Fig. 8 and compared with the interpolated fields for the BA and LGM. As noted elsewhere (Skinner and Bard, 2022; Rafter et al., 2022), the current paucity of data from the Indian Ocean makes interpolations for this basin somewhat tentative and more strongly dependent on individual data points (Fig. 7). This highlights the Indian basin as an important target for future work reconstructing past B-Atm changes. Across all the time-slice interpolations, correlations between observed and interpolated values for the same locations yield  $R^2$  values in the range 0.59–0.86. Because the interpolation is not performed on a 2D zonal plane, local B-Atm estimates may deviate from the zonal average where zonal gradients exist. Furthermore, the interpolations necessarily provide an average “snapshot” for an entire time slice and therefore will mask variability within the time slice. These aspects appear to be particularly relevant for HS1 in the Atlantic, as discussed below.

The global interpolations illustrated in Fig. 5 capture the main features of deglacial B-Atm evolution that have previously been identified in individual time series and in compiled time series that were grouped by basin and region, or water depth (Skinner and Bard, 2022). Broadly similar patterns have also been identified in averaged time series and 2D zonal projection contour plots of compiled data (Rafter et al., 2022). The main features include the following:

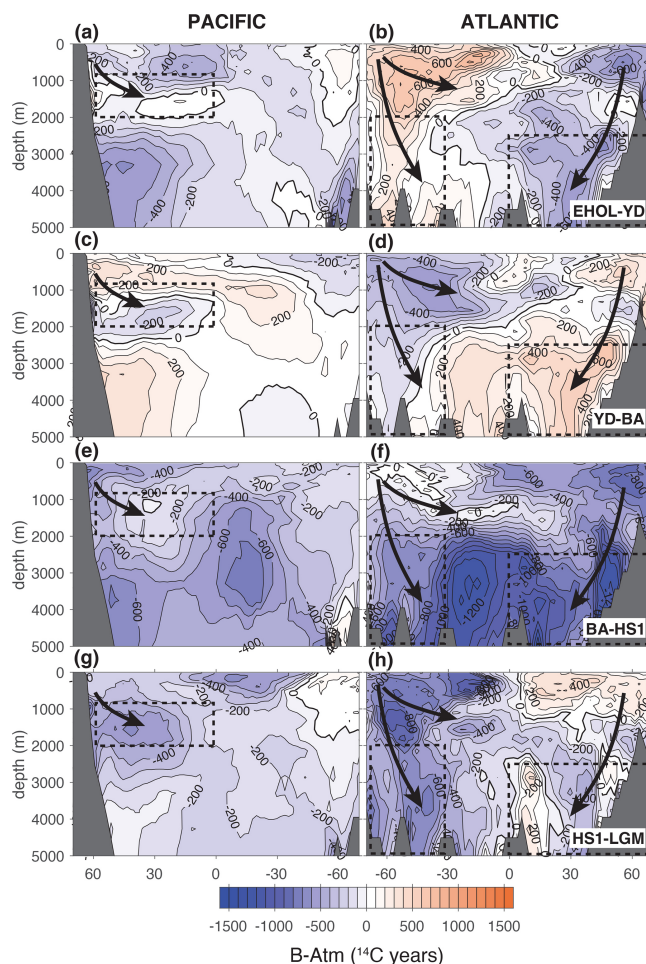
1. increased B-Atm offsets throughout the global ocean at the LGM as compared with the modern state (Fig. 8a, b, g, h), in particular > 2000 m water depth as demonstrated previously (Skinner et al., 2017) but with a slightly larger anomaly  $\sim 760$   $^{14}\text{C}$  yr;
2. a stepwise “rejuvenation” of the ocean interior across the last deglaciation;
3. evidence for positive B-Atm anomalies in the deep North Atlantic, from the LGM to HS1 and, again, from the BA to the YD, occurring in parallel with changes of broadly opposite sign in the Southern Ocean and the intermediate depth North Pacific (boxed areas in Fig. 6c, d, g, h);



**Figure 5.** Zonally averaged interpolated B-Atm radiocarbon age offsets for the LGM, HS1, BA, and YD (Pacific zonal averages at left, Atlantic at right). Filled circles and shading indicate input data and values.

4. a marked “rebound” to lower B-Atm offsets throughout the ocean, from HS1 to the BA (Fig. 6e, f); and
5. a further rebound to lower B-Atm offsets in the deep North Atlantic, from the YD to the EHOL, again with evidence for changes of opposite sign in the Southern Ocean and intermediate depth North Pacific (boxed areas in Fig. 6a, b).

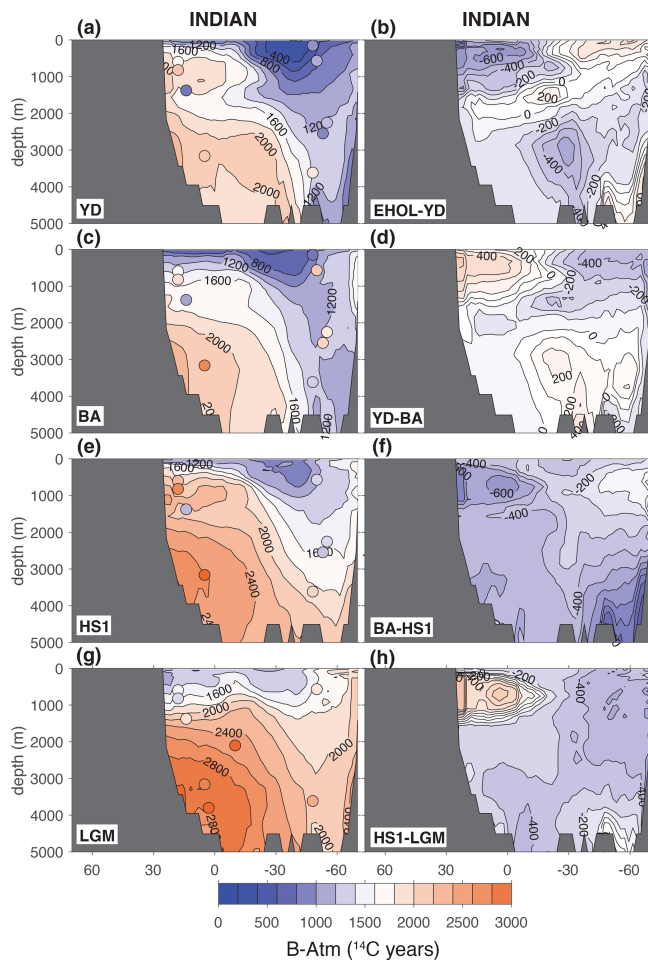
Perhaps the most striking aspect of the successive time-slice reconstructions is the marked drop in B-Atm offsets that coincides with the transition from HS1 to the BA (Fig. 6e, f), which involved positively correlated changes in B-Atm in the deep Southern Ocean and deep North Atlantic (Skinner et al., 2013), and which has been linked to an “overshoot” in B-Atm offsets at some locations (Barker et al., 2010; Hines et al., 2015). Notably, as shown in Fig. 8, the step change at the BA resulted in a global B-Atm distribution very similar to that



**Figure 6.** Offsets between spatial B-Atm interpolations for successive time-slice reconstructions. Boxed areas highlight regions of the deep North Atlantic (> 0° N, > 2.5 km), deep Southern Ocean (> 30° S, > 2 km), and intermediate North Pacific (> 0° N, 0.9–2 km) and deep North Pacific (> 0° N, > 2 km), for which regional time-series splines are illustrated in Fig. 9. Broadly antiphased anomalies are indicated between the North Atlantic and Southern Ocean (especially the Atlantic sector), and between the North Atlantic and intermediate North Pacific.

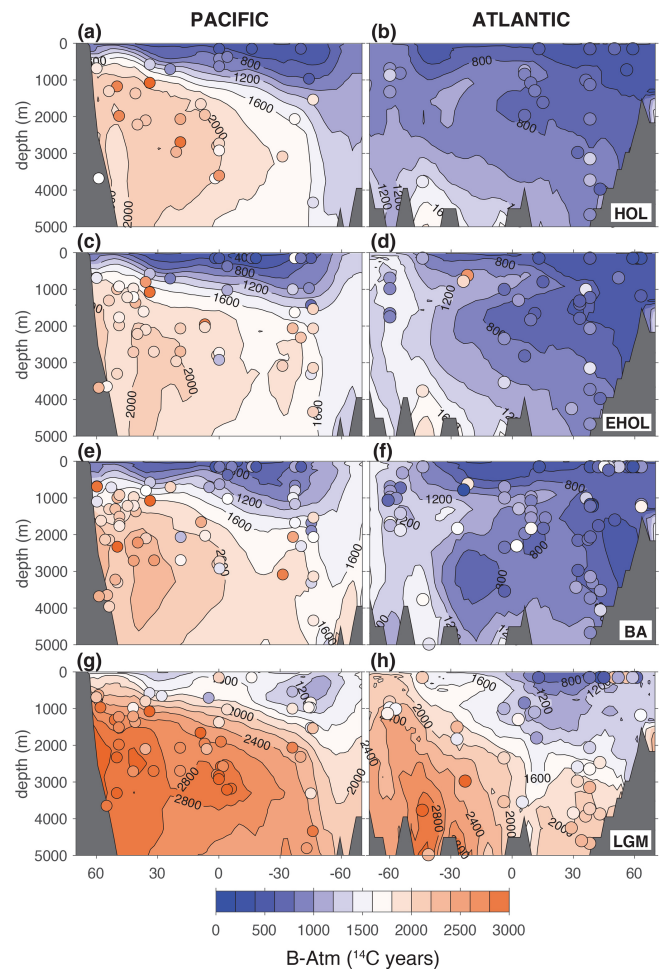
of the late Holocene, despite representing the approximate temporal midpoint of deglaciation.

The features identified above can also be discerned in regional time-series averages (Skinner and Bard, 2022; Rafter et al., 2022). We show this here using cubic splines for North Atlantic, Southern Ocean, and North Pacific data grouped according to the boxed areas highlighted in Fig. 6. The resulting regional average splines are illustrated in Fig. 9b–d compared with global average B-Atm values derived for each successive time slice (Fig. 9d, filled circles; Table 1). During HS1 and the YD, the regional splines exhibit broadly antiphased trends in B-Atm between the deep North Atlantic and the deep Southern Ocean and intermediate North Pacific (Fig. 9, shaded vertical bars). The collected time series thus support



**Figure 7.** Zonally averaged interpolated B-Atm radiocarbon age offsets in the Indian Ocean for the LGM, HS1, BA, and YD. Sparse data coverage is notable. Left: interpolated time-slice reconstructions (as for Fig. 5). Right: offsets between spatial B-Atm interpolations for successive time-slice reconstructions (as for Fig. 6).

the antiphase patterns apparent in the zonal averages shown in Fig. 6 (boxed areas). The averaged time series only tentatively support the suggestion of a “flipped” vertical radiocarbon gradient between the intermediate- and deep North Pacific at the LGM (Rafter et al., 2022), and more emphatically demonstrate this for the deglacial period that includes the YD, BA, and HS1. Indeed, during the BA, the North Pacific, North Atlantic, and Southern Ocean all exhibit B-Atm offsets at least as low as modern, resulting in global average B-Atm that also approaches the modern value and slightly “overshoots” relative to the subsequent YD (see Fig. 9d). Consistent with the fact that the Pacific accounts for  $\sim 50\%$  of the global ocean volume and the Southern Ocean ventilates  $\sim 58\%$  of the ocean interior (Primeau, 2005), Fig. 9 also shows that global average B-Atm estimates generally track the deep North Pacific and the deep Southern Ocean. Different patterns of variability are expressed in the North Atlantic



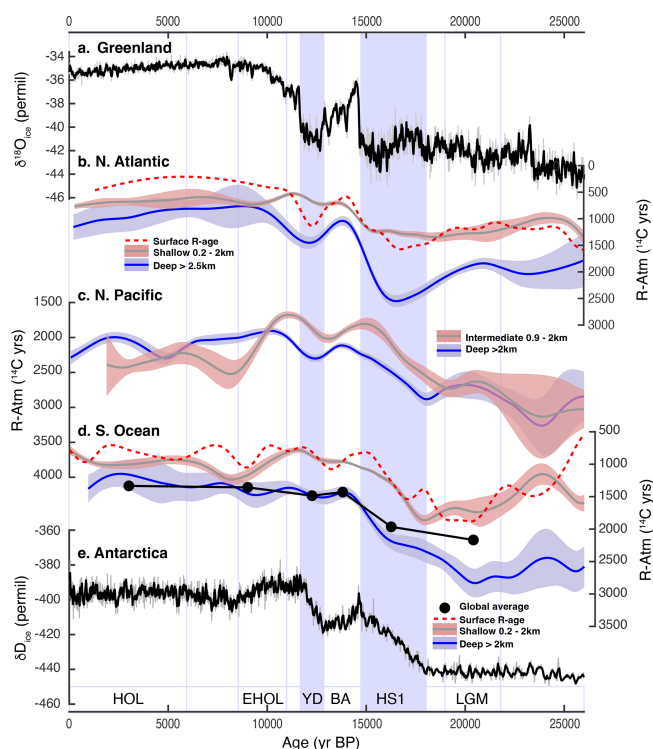
**Figure 8.** Zonally averaged interpolated B-Atm offsets for the LGM, BA, EHOL, and HOL (Pacific zonal averages at left, Atlantic at right). Filled circles and shading indicate input data and values.

and the intermediate North Pacific, where more rapid and regionally important fluctuations are apparent (Freeman et al., 2015).

## 4 Discussion

In principle, evolving large-scale patterns of ocean–atmosphere radiocarbon age offsets (B-Atm) will primarily reflect the combined influences of (1) radiocarbon production, (2) ocean transports, (3) air–sea gas-exchange efficiency (especially in the regions of deep-water export), and (4) the changing contributions of different source regions to locations in the ocean interior. While the influence of atmospheric radiocarbon production changes on evolving B-Atm offsets is often ignored, it may in principle influence deep ocean B-Atm offsets through relatively rapid (i.e. sub-millennial) changes in atmospheric radiocarbon activity that are only conveyed to the deep ocean on the millennial timescale of ocean turnover (Adkins and Boyle, 1997). This can produce





**Figure 9.** “Ventilation seesaws” across the last deglaciation based on cubic spline fits to compiled ocean–atmosphere radiocarbon age offsets (Skinner and Bard, 2022). Cubic splines with a stiffness factor of  $10^{-9}$  are fit to all existing data, using updated age models and taking into account B-Atm uncertainties, and applying the same “baseline” data flags as for the time-slice interpolations. (a) Greenland temperature proxy (Svensson et al., 2008). (b) northeast Atlantic shallow sub-surface reservoir ages (dashed red line; Skinner et al., 2019); B-Atm from the North Atlantic 0.2–2 km (grey line, shaded red area); B-Atm from the deep North Atlantic > 2.5 km (blue line, shaded blue area). (c) B-Atm from the intermediate North Pacific 0.9–2 km (grey line, shaded red area); B-Atm from the deep North Pacific > 2 km (blue line, shaded blue area). (d) Mean ocean B-Atm estimates (black line and circles; see Table 1 “baseline” data flags) and compiled shallow sub-surface reservoir ages from the Southern Ocean (dashed red line; Skinner et al., 2019); B-Atm from the Southern Ocean < 2 km (grey line and shaded red area); B-Atm from the deep Southern Ocean > 2 km (blue line and shaded blue area). (e) Antarctic temperature proxy (EPICA Community Members, 2004; Lemieux-Dudon et al., 2010). Vertical lines and shaded bars indicate the timing of the LGM, HS1, BA, YD, EHOL, and HOL time slices.

a convergence or divergence of marine and atmospheric radiocarbon ages, due to atmospheric radiocarbon changing quickly and the deep ocean remaining relatively invariant (Franke et al., 2008; Heaton et al., 2021).

Below, we discuss how all these processes have influenced marine radiocarbon cycling across the last deglaciation. Accordingly, we address (1) evidence for the operation of a “ventilation seesaw” that we emphasize was linked to *both* gas-exchange and transport anomalies emanating from the

main regions of deep- and intermediate water formation (i.e. the North Atlantic, Southern Ocean, and North Pacific); and (2) the potential for atmospheric radiocarbon dynamics (independent of ocean ventilation) to bias B-Atm offsets by  $\sim$  hundreds of  $^{14}\text{C}$  yr, in particular during the apparent BA “overshoot”. We further seek to quantify the likely carbon cycle impacts associated with the observed global average B-Atm changes, and to reconcile these impacts with the evolution of the global radiocarbon budget since the last glacial period.

#### 4.1 Ventilation “seesaws”: gas-exchange and transport effects

A survey of the modern ocean’s transport pathways indicates that  $\sim 86\%$  of the ocean interior is sourced by water that last made contact with the atmosphere in three key regions: the Southern Ocean (contributing  $\sim 58\%$ ), the high latitude North Atlantic (contributing  $\sim 21\%$ ), and the North Pacific (contributing  $\sim 7\%$ ) (Primeau, 2005). These three regions account for  $\sim 31\%$  of the ocean’s surface and “ventilate”  $\sim 86\%$  of the ocean’s interior (Primeau, 2005). Although the long equilibration time for dissolved  $\Delta^{14}\text{C}$  (DIC) means that a water parcel’s point of last contact with the atmosphere will not necessarily be equivalent to its point of “radiocarbon equilibration” (high latitude sources are typically characterized by distinct levels of disequilibrium (Bard, 1988; Matsumoto, 2007), changes in the three main regions of deep-water export will exert a strong influence on temporal variations in the ocean interior’s radiocarbon distribution.

This expectation is clearly borne out in the global interpolations shown in Figs. 5 and 6, and in the regional stacks illustrated in Fig. 9. Deglacial changes in marine radiocarbon distribution were apparently dominated by anomalies extending from the North Atlantic (affecting the deep Atlantic in particular), coordinated with anomalies of broadly opposite sign originating in the Southern Ocean and North Pacific (Figs. 6 and 9). The time-slice interpolations and regional averages therefore cohere with numerous previous proposals for “ventilation seesaws” operating between the North Atlantic and the Southern Ocean (Broecker, 1998; Skinner et al., 2013, 2014; Menviel et al., 2018), and between the North Atlantic and North Pacific (Menviel et al., 2014; Freeman et al., 2015; Max et al., 2014; Okazaki et al., 2010; Walczak et al., 2020).

As suggested previously (Skinner et al., 2019; Skinner and Bard, 2022), Fig. 9 also shows that B-Atm offsets in the deep (> 2 km) North Atlantic (Fig. 9b, blue line and shaded area) exhibit a similar pattern of variability to that of the upper ocean (0.2–2 km) and surface “reservoir ages” in the region (Fig. 9b grey lines with shaded area, and dashed red lines). The same is apparent in the Southern Ocean (Fig. 9d). These relationships, and their further link to polar climate variability (Fig. 9a, e), suggest a mechanistic link between the observed deglacial B-Atm variability and the “thermal bipolar seesaw” (Stocker and Johnsen, 2003; EPICA Community



Members, 2006). This association most likely operated via coordinated changes in North Atlantic and Southern Ocean convection and advection (Broecker, 1998; Menviel et al., 2015; Skinner et al., 2014, 2020) associated with regional changes in sea ice (Skinner et al., 2019; Rae et al., 2018), winds (Sikes et al., 2016b; Menviel et al., 2018), and/or buoyancy forcing (Ferrari et al., 2014; Hines et al., 2019; Watson et al., 2015). The further coupling between the Southern Ocean and North Pacific has been proposed to relate to freshwater balance in the Pacific basin (Menviel et al., 2014), possibly influenced by changing moisture transports across the Isthmus of Panama (Leduc et al., 2007), and/or Cordilleran ice mass balance (Walczak et al., 2020).

While the “ventilation seesaws” noted above may reflect changes in ocean circulation to some extent, it is important to note that they also reflect changes in gas-exchange efficiency. This is demonstrated by the fact that

the amplitude of B-Atm variability in the shallow Atlantic 0.2–2 km water depth (e.g. Freeman et al., 2015; Chen et al., 2015) differs very little from that of surface “reservoir ages” in the northeast Atlantic (Fig. 9b, grey line and dashed red line). This implies a muted contribution from flow speed changes in the upper Atlantic < 2 km (Bradtmeier et al., 2014) and a dominant influence on radiocarbon signatures from gas-exchange (i.e. “pre-formed ages”) instead. Again, the same is true for the Southern Ocean, where B-Atm offsets from the shallow ocean (0.2–2 km) generally overlap with surface reservoir age reconstructions (Fig. 9d, grey line and dashed red line). Accordingly, B-Atm values at the LGM in mid-depths may indeed have been consistent with the modern transport (Marchal and Zhao, 2021a).

In contrast, however, “pre-formed ages” cannot account for the amplitude of B-Atm changes observed in the deep Atlantic (Fig. 9b, blue line) and the deep Southern Ocean (Fig. 9d, blue line), indicating a more significant contribution from water sourcing and/or transport changes > 2 km, and perhaps between 2 and 3 km water depth in particular (Skinner et al., 2021; Lund et al., 2015; Rafter et al., 2022). The patchy response seen in the zonal average anomaly for the North Atlantic between HS1 and the LGM (Fig. 6h) contrasts with the clearer signal seen in individual and collected time series (Fig. 9b). In part, this likely reflects a spatially heterogeneous hydrographic response, both in the depth domain (Skinner et al., 2021; Lund et al., 2015) and in the eastern versus western North Atlantic (Gherardi et al., 2005; Ng et al., 2018). This interpretation is supported by the less ambiguous positive B-Atm anomaly seen in the deep North Atlantic regional time-series spline during HS1 (Fig. 9b, blue line and shaded area; see also Rafter et al., 2022).

The patchy spatial interpolation outcome for HS1 may also reflect a more complex pattern of ventilation during HS1 than is captured by a “collapsed AMOC” scenario that lasted until the onset of the BA, as indicated by a recent multi-proxy and modelling study (Pöppelmeier et al., 2023). Furthermore, in addition to tentative evidence for a “mid-HS1 B-Atm mini-

mum” (Skinner and Bard, 2022), there is evidence for a drop in B-Atm offsets at intermediate depths (< 2.5 km) in the western Atlantic late in HS1 (Robinson et al., 2005; Thiagarajan et al., 2014), and evidence for declining B-Atm offsets in the Nordic Seas ~ 400 years prior to the onset of the BA (Muschitiello et al., 2019). Any such changes during HS1 will have been averaged out in our time-slice interpolation for 15–17.8 ka BP. These observations underline the need for further detailed reconstruction of the spatial expression and temporal evolution of ocean–atmosphere radiocarbon offsets across the North Atlantic during HS1, particularly with a view to disentangling transport and gas-exchange impacts.

Overall, the message that emerges from the collected data is that deglacial marine B-Atm changes throughout the global ocean were significantly influenced by *both* air–sea gas-exchange effects and mass transport effects (Koeve et al., 2015), with the latter primarily affecting parts of the deep ocean > 2 km, but with evidence for a complex response during HS1. The inferred spatial expression and coordination of both gas-exchange and transport effects in the Atlantic may have been related to changes in the spatial extent of different water masses in the deep ocean. This would be consistent with previous work where B-Atm offsets have been found to place stronger constraints on water mass “geometry” than, for example, AMOC strength per se (Muglia and Schmittner, 2021; Pöppelmeier et al., 2023). The influence of further non-ventilation effects at each time-slice is taken up in the following section.

#### 4.2 “Attenuation biases” in B-Atm offsets at the BA

As noted above, global average B-Atm changes can occur independently of ventilation effects, due to rapid atmospheric radiocarbon variability to which the deep ocean is too slow to respond. We refer to such effects as “attenuation biases”, as they reflect the phase and attenuation response of a slowly adjusting reservoir (the global ocean interior), subject to continuous exchange with a more rapidly changing reservoir (the atmosphere) (Maier-Reimer and Hasselmann, 1987). The primary external (i.e. non-marine) driver for rapid atmospheric radiocarbon variability is likely to be changes in radiocarbon production (Köhler et al., 2022), though transient terrestrial carbon sources (e.g. from permafrost) might also be hypothesized (Köhler et al., 2014; Wu et al., 2022).

For attenuation biases in mean ocean B-Atm offsets to occur, atmospheric radiocarbon activity would have to change rapidly, that is, on a timescale that is shorter than the mean mixing time of the ocean (i.e. < 1000 years). Accordingly, gradual long-term trends in radiocarbon production are unlikely to produce significant attenuation biases, as is confirmed by model simulations where atmospheric radiocarbon production is based on relatively smooth trends in mean relative (geomagnetic) palaeointensity (RPI) (Dinauer et al., 2020). However, given the scatter amongst existing reconstructions of past radiocarbon production (e.g.

Laj et al., 2004; Adolphi et al., 2018; Nowaczyk et al., 2013; Channell et al., 2018), particularly on (sub-) millennial timescales, it remains unclear to what extent rapid (centennial or millennial) radiocarbon production variability and/or other non-marine carbon sources might have biased B-Atm offsets via their impact on the atmosphere. Therefore, to explore the maximum possible contribution of externally driven atmospheric radiocarbon variability to mean ocean B-Atm changes, we compare transient simulations using the Bern3D model where radiocarbon production rates or atmospheric radiocarbon are held constant (i.e. FIX- and CONST-, respectively) with simulations where nearly all variability in atmospheric radiocarbon is assumed to have occurred independently of ocean ventilation change (i.e. INT-, where atmospheric radiocarbon is prescribed according to *Intcal20*). The difference between the INT- simulations and their CONST/FIX- counterparts yields the *maximum* amplitude of B-Atm changes that could be produced independently of ocean ventilation change, for each idealized scenario. Indeed, as discussed below, these estimates are likely overestimates as they are premised on the assumption that all atmospheric radiocarbon variability occurred independently of ocean–atmosphere radiocarbon exchange, which is unlikely.

As illustrated in Fig. 10 (and summarized in Table S1), our idealized scenarios demonstrate three key points regarding the emergence of “attenuation biases” in B-Atm offsets:

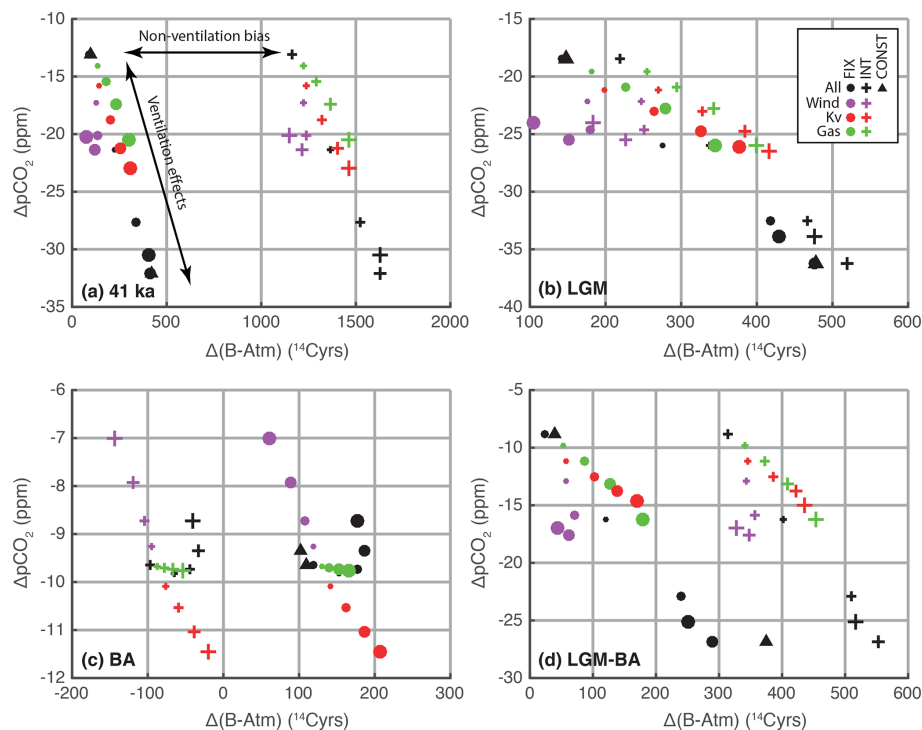
1. These biases depend on the occurrence of rapid atmospheric variability that is not forced by ocean ventilation change, but they can in principle result in persistent long-term biases via an accumulation of centennial or millennial perturbations (e.g. as for a “random walk” process).
2. Such biases are time varying and may be positive relative to a given reference time (e.g. at the Laschamps event; Fig. 10a), negative (e.g. notably, at the BA onset; Fig. 10c), or nil (e.g. at the LGM; Fig. 10b), thus enhancing, diminishing, or not affecting the B-Atm changes that are expressed between time periods (Fig. 10d).
3. The relative contribution of such biases will be diminished and potentially eliminated to the extent that they coincide with large abrupt ocean ventilation changes (not included in our idealized model simulations), though it is apparent that their magnitude depends mainly on the amplitude of atmospheric radiocarbon production changes, as they yield a relatively invariant offset between the INT- and CONST/FIX- simulations for any given time slice; Fig. 10a–d).

It is important to stress that the true magnitude of such attenuation biases cannot be determined without prior detailed knowledge of the history and magnitude of both ocean ventilation changes and non-marine carbon and/or radiocarbon

inputs to the atmosphere. Nevertheless, two observations are worth underlining: the first is that maximum estimates of the attenuation biases that could hypothetically affect deglacial B-Atm offsets would only result in relatively minor biases in the incremental B-Atm changes between time slices (ranging from  $-227$  to  $+25$   $^{14}\text{C}$  yr). Even if maximal attenuation biases are hypothesized, the observed deglacial mean ocean B-Atm trends must include a significant ventilation contribution. Indeed, such changes are directly attested to by proxy evidence for ocean transport and sea-ice change (McManus et al., 2004; Schüpbach et al., 2018), as well as the spatially heterogeneous patterns in marine B-Atm offsets that we present here (e.g. Figs. 6 and 9).

A second key observation is that correcting for such biases would tend to diminish the apparent “ventilation surge” that might be inferred from the change in B-Atm between HS and the BA (and more so between the LGM and the BA). Thus, if the rapid atmospheric radiocarbon decline across HS1 and at the onset of the BA were entirely driven by radiocarbon production changes and/or non-marine carbon inputs to the atmosphere, then the change in global average B-Atm between the LGM and the BA would be biased by at most  $\sim 227$   $^{14}\text{C}$  yr. In this case, radiocarbon ventilation ages during the BA would be  $\sim 227$   $^{14}\text{C}$  yr *older* than inferred from observed B-Atm offsets, resulting in a smaller radiocarbon ventilation change between the LGM and the BA. Non-ventilation biases affecting global average B-Atm differences between the LGM and the BA could therefore imply a smaller contribution from gas-exchange and transport rate changes to the apparent radiocarbon “ventilation surge” suggested for the BA in Fig. 6e and f. In this case, the “rejuvenation” of the marine radiocarbon pool may not in fact have been completed midway through deglaciation as initially apparent (Rafter et al., 2022). This inference would be more consistent with stable carbon isotope ( $^{13}\text{C}/^{12}\text{C}$ ) evidence (Sikes et al., 2016b), which suggests an ongoing contribution to ocean–atmosphere carbon exchange well beyond the BA. Indeed, a significant portion of the convergence between marine and atmospheric radiocarbon observed at the BA may have been driven by “old” carbon release to the atmosphere, e.g. from melting permafrost (Köhler et al., 2014; Wu et al., 2022).

Although they remain difficult to accurately quantify, “attenuation biases” unrelated to ventilation are important to acknowledge, as these may exert a subtle yet potentially significant influence on our interpretation of transient changes in B-Atm offsets and their quantification, as illustrated above for the BA. These effects will need to be explored in greater detail in the future, a task that will inevitably require the use of numerical models and also require accurate knowledge of past radiocarbon production changes (Köhler et al., 2022; Dincau et al., 2020). The accuracy of our existing radiocarbon production records is discussed further in Sect. 4.4.



**Figure 10.** Model outputs from idealized “deglacial” scenarios illustrating hypothetical attenuation biases that are unrelated to ocean “ventilation” effects (indicated by horizontal offsets between similar symbol colours; horizontal arrow in panel (a)). Panel (a) shows 41 ka BP (coinciding with the Laschamps geomagnetic excursion), (b) the LGM, (c) the BA, and (d) the offset between the LGM and the BA. Data are expressed as 1000-years average anomalies relative to the pre-industrial period 4–5 ka BP, when atmospheric radiocarbon was relatively stable and therefore plausibly approached a pseudo-equilibrium state. In all panels, crosses indicate model experiments carried out with prescribed atmospheric radiocarbon as given by *Intcal20* (Reimer et al., 2020) (INT), circles are for atmospheric radiocarbon fixed at 140‰ (FIX), and triangles are for constant pre-industrial radiocarbon production rates (CONST). Symbol colours distinguish experiments carried out with altered Southern Ocean winds (purple), vertical diffusivity (red), Southern Ocean gas-exchange efficiency (green), or all combined (black). Four symbol sizes indicate the extent of parameter reduction: 0 % (control, smallest symbols), 20 %, 40 %, 60 %, and 80 % (largest symbols). For CONST, only two sets of experiments were run: for control conditions and for all tuning variables reduced by 60 %. Note that the impact on B-Atm due to reduced Southern Ocean winds (purple circles) is attenuated in well-equilibrated states (e.g. 41 ka and LGM), because of, for example, slow adjustments in the distribution of North Atlantic sourced deep water (see text).

### 4.3 Ventilation-related atmospheric CO<sub>2</sub> change

In theory, a broadly linear relationship between atmospheric CO<sub>2</sub> change and global average marine radiocarbon age anomalies is to be expected when these are driven by gas-exchange and/or transport rates (Skinner and Bard, 2022; Skinner et al., 2017). This is because (1) longer residence times in the ocean interior result in greater respired carbon accumulation along with greater radiocarbon decay, and (2) restricted gas exchange in regions of “upwelling” and/or deep mixing impede the conversion of respired and/or disequilibrium carbon to equilibrium carbon (Eggleston and Galbraith, 2018) while also impeding the invasion of radiocarbon into the ocean.

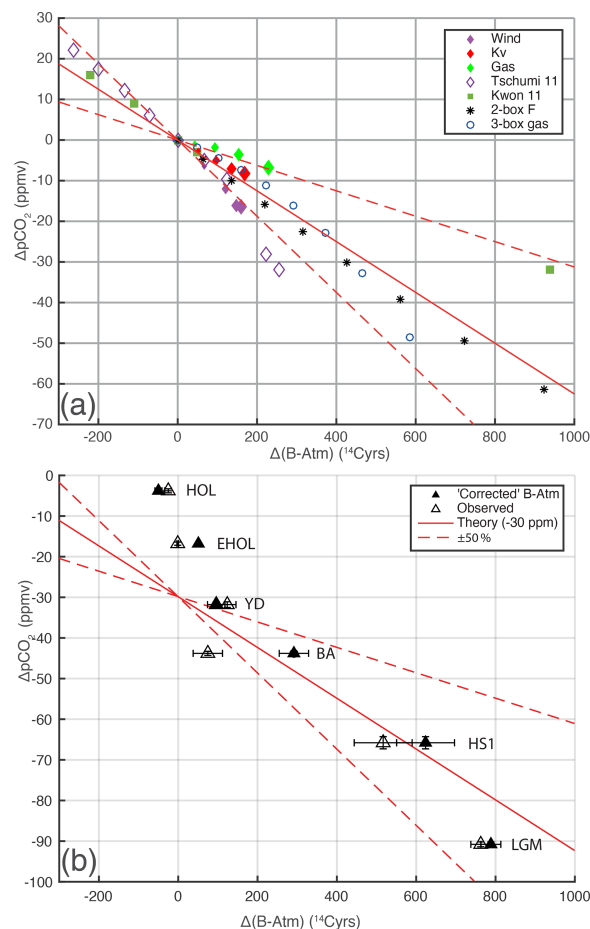
Our model sensitivity tests, involving shifts in vertical diffusivity, Southern Ocean winds, and/or gas exchange, cohere with a number of existing simulations using box models and Earth system models of intermediate complexity (e.g. Tschumi et al., 2011; Kwon et al., 2011; Skinner and

Bard, 2022), confirming a broad relationship between B-Atm anomalies and associated atmospheric CO<sub>2</sub> change (Fig. 11a, Table 2), despite large differences in model experiment set-up. These experiments include the effects of *p*CO<sub>2</sub> changes on air–sea <sup>14</sup>C exchange (Galbraith et al., 2015; Bard, 1988). Our sensitivity tests indicate consistent sensitivities for individual suites of experiments, e.g. for varying Southern Ocean winds, vertical diffusivity, or Southern Ocean gas-exchange rates. However, depending on the processes responsible for altering deep ocean ventilation, modelled sensitivities span a range of approximately ±50 % for a given magnitude of B-Atm change. Furthermore, the sensitivities may vary with perturbation timescale, with different impacts on millennial timescales than on glacial-interglacial or longer timescales (Jeltsch-Thömmes and Joos, 2020). For example, a reduction in Southern Ocean winds draws down atmospheric CO<sub>2</sub> and increases global average B-Atm on a timescale of ∼2000 years via impacts on gas-exchange

rates and upper ocean turnover (Fig. 11a). However, after  $\sim 5000$  years (see Fig. S1) the mean ocean B-Atm response is diminished, primarily due to the spin up of North Atlantic deep water, which replaces southern sourced water in the deep Atlantic. A muted B-Atm response can therefore be seen for stronger Southern Ocean wind forcing in the 50 ka simulations illustrated in Fig. 10a (purple symbols). Despite their potential role in transient millennial scale pulses, Southern Ocean winds are therefore unlikely to account for long-term glacial-interglacial B-Atm changes, at least in the absence of further forcing to suppress North Atlantic ventilation (Menviel et al., 2018).

The broadly linear scaling apparent in the suite of PI model sensitivity tests is consistent with basic theory (Skinner et al., 2017; Skinner and Bard, 2022), based on a two-box ocean connected to an atmosphere to form a closed system (Fig. 11a, solid line). This simple inventory theory predicts a higher sensitivity for higher global export productivity and/or a higher “Revelle buffer factor” (i.e. higher background atmospheric  $p\text{CO}_2$ ), all else being equal. Clearly, the marine carbon cycle response to the variety of ventilation processes that can affect radiocarbon cannot be reduced to a single linear scaling. However, the degree of consistency in the parallel sensitivities of atmospheric  $\text{CO}_2$  and marine B-Atm offsets implies that the wide range of modelled sensitivities may be approximated by a theoretical prediction using an arbitrary  $\pm 50\%$  range in export productivity as a tuning parameter in the two-box ocean model (Fig. 11a, dashed lines). This theoretical scaling, arbitrarily tuned to the range of more complex model outputs, would suggest  $-6.3 \pm 3.2$  ppm  $\text{CO}_2$  change per  $100$   $^{14}\text{C}$  yr of global mean B-Atm change. Such a sensitivity would tentatively imply a drawdown of atmospheric  $\text{CO}_2$  by  $\sim 50 \pm 27$  ppm associated with an increase in global average B-Atm by  $\sim 789 \pm 32$   $^{14}\text{C}$  yr, as reconstructed for our “baseline” data flag scenario at the LGM and assuming a maximum “attenuation bias correction” (Table 1). Estimates based on uncorrected B-Atm offsets and/or on alternative data flagging scenarios (Table 1) differ by up to  $\sim 8$  ppm. These estimates should be interpreted as indicating a non-negligible contribution to deglacial  $\text{CO}_2$  rise, perhaps equivalent to over a third of the total glacial-interglacial atmospheric  $\text{CO}_2$  change.

Interestingly, observed global average B-Atm estimates also loosely correlate with observed atmospheric  $\text{CO}_2$  changes across the last deglaciation (Fig. 11b, open triangles). The correlation with observed  $p\text{CO}_2$  anomalies, particularly for the BA, is improved when global average B-Atm is “corrected” for maximal possible attenuation biases as discussed above (Fig. 11b, filled triangles). The similarity of the observed correlation and the modelled  $\text{CO}_2$  sensitivity (Fig. 11b, dashed and solid red lines) again merely suggests that a significant portion of the incremental changes in atmospheric  $\text{CO}_2$ , stepping through the deglaciation from the LGM to the early Holocene, could in principle be accounted for by ocean ventilation changes that influenced



**Figure 11.** Panel (a) shows theoretical and modelled sensitivity of atmospheric  $\text{CO}_2$  anomalies and mean ocean radiocarbon ventilation age (B-Atm) anomalies to air–sea gas-exchange and transport changes. Filled diamonds indicate model sensitivity tests from this study, based on step changes under PI conditions and 2000 years of equilibration time (purple: Southern Ocean wind, red: vertical diffusivity, green: Southern Ocean gas-exchange efficiency). Open diamonds are Southern Ocean wind experiments (Tschumi et al., 2011). Filled squares are wind and/or diffusivity experiments using an idealized radiocarbon-like tracer (Kwon et al., 2011). Asterisks and open circles represent two- and three-box model experiments, with varying overturning rates ( $F$ ) and “high latitude” gas-exchange rates (gas), respectively. A median theoretical sensitivity is indicated by the solid and dashed red lines, derived for the two-box model results (Skinner and Bard, 2022), equivalent to  $-6.3 \pm 3.2$  ppm  $\text{CO}_2$  change per  $100$   $^{14}\text{C}$  yr of global radiocarbon ventilation age change. Panel (b) shows observed values of coincident global average B-Atm and atmospheric  $p\text{CO}_2$ , expressed as time-slice anomalies versus pre-industrial values. Open triangles show paired observations of mean ocean B-Atm (this study) and atmospheric  $\text{CO}_2$  (Monnin et al., 2001; Marcott et al., 2014). Filled triangles show observed B-Atm age offsets corrected for hypothetical “non-ventilation” biases (see text). Solid and dashed red lines in (b) show the same theoretical sensitivities as for (a), offset by  $-30$  ppm to aid comparison with observed trends.



**Table 2.** B-Atm and atmospheric CO<sub>2</sub> results for sensitivity experiments carried out using the Bern3D model with varying Southern Ocean (SO) wind, gas exchange, or global vertical diffusivity (illustrated in Fig. 10). Results are evaluated 2000 years after forcing is applied (see text).

Scenario	Ventilation parameter changed	Reduction (%)	B-Atm	$p\text{CO}_2$	$\Delta(\text{B-Atm})$	$\Delta(p\text{CO}_2)$
PI (ctrl)	none	0	1428	276.5857	0	0
PI	SO wind	20	1495	270.2083	67	−6
		40	1548	264.4817	120	−12
		60	1577	260.3277	150	−16
		80	1588	259.9954	160	−17
PI	SO gas exchange	20	1472.1	275.7578	44	−1
		40	1523.3	274.6609	96	−2
		60	1585.5	272.8658	158	−4
		80	1662.4	269.6315	235	−7
PI	Vertical diffusivity	20	1480.3	273.7565	53	−3
		40	1523.6	271.5337	96	−5
		60	1564.6	269.6152	137	−7
		80	1597.5	268.211	170	−8

global average B-Atm offsets. Although the steeper dashed line in Fig. 11b would imply a maximal  $\sim 80$  ppm contribution to deglacial atmospheric CO<sub>2</sub> rise due to ocean ventilation alone, we believe this is unlikely. Rather, the observed correlation between atmospheric CO<sub>2</sub> and global average B-Atm offsets more likely implies a mixture of direct and indirect causal connections that may have been coordinated by the thermal bipolar seesaw. Direct impacts of ocean ventilation on atmospheric CO<sub>2</sub> (e.g. Fig. 11a) would therefore have coincided with contributions from linked processes, such as ocean temperature change, export productivity anomalies, and others (Marchal et al., 1998; Menviel et al., 2008; Jochum et al., 2022; Gottschalk et al., 2019; Menviel et al., 2012), or indeed long-term Southern Ocean wind changes that have had little effect on mean ocean B-Atm.

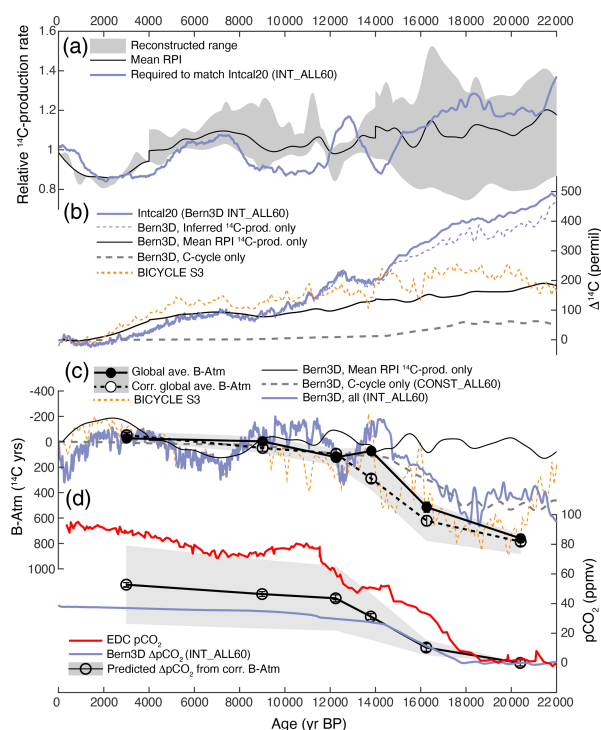
In any event, if mean ocean B-Atm changes scale in a consistent manner with marine CO<sub>2</sub> sequestration changes (as suggested by simple inventory theory and the relationships in Fig. 11a), then the ocean ventilation contribution to deglacial atmospheric CO<sub>2</sub> rise would have been primarily associated with HS1, the BA, and the YD. The magnitude of ocean ventilation impacts on CO<sub>2</sub> after  $\sim 15$  ka BP appear to hinge crucially on attenuation biases that may have affected B-Atm offsets during the BA. Nevertheless, the increase in atmospheric CO<sub>2</sub> observed across the Holocene coincided with rather muted changes in global average B-Atm and low potential for significant attenuation biases. This suggests a minor role for ocean ventilation in CO<sub>2</sub> rise from the onset of the Holocene, which would be consistent with the inference that CO<sub>2</sub> rise from  $\sim 6$  ka BP was primarily linked to changes in ocean temperature, the terrestrial biosphere, coral reef formation and carbonate preservation (i.e. ocean alkalinity), and/or the solid earth (i.e. volcanism) (Joos et al., 2004; Broecker and Clark, 2007). This observation might also res-

onate with the speculative proposal of a “natural tendency” for atmospheric CO<sub>2</sub> decline across an interglacial due to ocean ventilation processes (Barker et al., 2019).

#### 4.4 Towards a closure of the global carbon and radiocarbon cycles since the LGM

The reconciliation of past radiocarbon production changes with records of atmospheric  $p\text{CO}_2$  and  $\Delta^{14}\text{C}$  (hereafter  $\Delta^{14}\text{C}_{\text{atm}}$ ) across the last deglaciation represents a long-standing puzzle that remains unresolved (Bard, 1998). This puzzle has a direct bearing on our understanding of past atmospheric  $p\text{CO}_2$  change, as well as our understanding of geomagnetic and solar variability (Heaton et al., 2021). The record of atmospheric radiocarbon variability indicates a significant decrease across the last deglaciation, equivalent to a change in  $\Delta^{14}\text{C}_{\text{atm}}$  of just over  $\sim 400$ ‰ (Reimer et al., 2020) (Fig. 12b, blue line). If there were no change in the steady state global radiocarbon inventory, the 90 ppm increase in atmospheric  $p\text{CO}_2$  that occurred since the last glacial period would alone account for only  $\sim 25$ ‰ of this change (Bard, 1998; Siegenthaler et al., 1980). This implies a dominant role for changes in the global radiocarbon budget (i.e. radiocarbon production) and/or the distribution of radiocarbon between atmosphere and other carbon reservoirs (i.e. the carbon cycle).

Model simulations of deglacial radiocarbon production and atmospheric radiocarbon and  $p\text{CO}_2$  consistently indicate that past  $\Delta^{14}\text{C}_{\text{atm}}$  cannot be accounted for by existing reconstructions of changing radiocarbon production alone (e.g. Hain et al., 2014; Köhler et al., 2006; Dinauer et al., 2020; Köhler et al., 2022). Earth system model simulations applying mean relative palaeomagnetic intensity- (RPI) based radiocarbon production rate changes (Dinauer et al., 2020),



**Figure 12.** Observed and modelled atmospheric and marine radiocarbon, compared with radiocarbon production rates, and atmospheric  $\text{CO}_2$ . Panel (a) shows relative radiocarbon production rates. Shaded area: full range of reconstructed values (Laj et al., 2000, 2004; Adolphi et al., 2018; Channell et al., 2018; Nowaczyk et al., 2013); solid black line: mean RPI (Laj et al., 2000, 2004; Channell et al., 2018; Nowaczyk et al., 2013); thick blue line: as inferred from an idealized model scenario to match *Intcal20*, INT\_ALL60 (1 kyr moving average; this study). Panel (b) shows atmospheric  $\Delta^{14}\text{C}$  anomalies normalized to modern values: prescribed in the INT\_ALL60 model scenario (i.e. *Intcal20*, Reimer et al., 2020; solid dark blue line; this study); driven only by inferred production rates that match *Intcal20* (INT\_ALL60-CONST\_ALL60; dashed blue line; this study); driven only by mean RPI production rates (solid black line; Dincauer et al., 2020); driven only by carbon cycle and ventilation changes in the CONST-ALL60 model scenario (dashed grey line; this study); simulated with the BICYCLE box model (Köhler et al., 2006) using  $^{10}\text{Be}$ -based (Muscheler et al., 2005) radiocarbon production estimates and a full carbon cycle scenario (dashed orange line). Panel (c) shows B-Atm radiocarbon age offset anomalies relative to modern values: Bern3D and BICYCLE model outputs as for (b); and inferred from time-slice interpolations (filled black circles and line; this study), including a maximal correction for “attenuation biases” (open black circles and dashed line), with the full range of reconstructed values (shaded area) due to uncertainties, corrections, and alternative data flagging scenarios (see Table 1). Panel (d) shows atmospheric  $\text{CO}_2$  normalized to LGM values: from EDC (red line; Monnin et al., 2001); inferred from observed mean “corrected” ocean B-Atm using a sensitivity of  $-6.3 \pm 3.2$  ppm per 100  $^{14}\text{C}$  yr (open black circles and line with shaded range; this study); and simulated in the Bern3D model for the INT\_ALL60 scenario that is also illustrated in (a)–(c).

yield only a small increase ( $\sim 150\%$ ) in  $\Delta^{14}\text{C}_{\text{atm}}$  at the LGM relative to the late Holocene (Fig. 12b, black line). Similar  $\Delta^{14}\text{C}_{\text{atm}}$  changes are obtained using alternative radiocarbon production histories (Hain et al., 2014; Köhler et al., 2006; Dincauer et al., 2020; Köhler et al., 2022). The widely recognized implication of these results is that additional carbon cycle changes (i.e. altered rates of carbon exchange with other reservoirs) and/or different radiocarbon production changes are required in order to account for the observed  $\Delta^{14}\text{C}_{\text{atm}}$  amplitude.

Turning to carbon cycle changes first, given the tight coupling of the marine and atmospheric carbon pools, altered exchange rates between the ocean and atmosphere are likely to have played a leading role in deglacial  $\Delta^{14}\text{C}_{\text{atm}}$  variability (Muscheler et al., 2004). Arguably for the first time, our global average B-Atm estimates confirm such a role. Indeed, our mean ocean B-Atm estimates indicate a lower average exchange rate of radiocarbon (and likely  $\text{CO}_2$ ) between the ocean and atmosphere during the last glacial period, and an increase in this exchange rate across the deglaciation (Fig. 12c, black line and circles). All else being equal, an increase in ocean–atmosphere radiocarbon exchange would result in a drop in  $\Delta^{14}\text{C}_{\text{atm}}$  during deglaciation, in parallel with a decrease in marine B-Atm offsets, as observed in Fig. 12b and c. Such changes would have also contributed to deglacial atmospheric  $\text{CO}_2$  rise (e.g. Muglia et al., 2018; Khatiwala et al., 2019; Brovkin et al., 2012; Ganopolski and Brovkin, 2017). As discussed above, a tentative quantification of this contribution to atmospheric  $\text{CO}_2$  rise can be derived from the observed mean ocean B-Atm (Table 1) and modelled sensitivities (Fig. 11a), as illustrated in Fig. 12d (open circles and shaded region). This tentative contribution compares well with simulated  $\text{CO}_2$  effects in the Bern3D model (INT\_ALL60), bearing in mind that the simulated ageing of the global ocean is only  $\sim 50\%$  of that observed (Fig. 12c, blue line and dashed grey line) and that  $\sim 50\%$  of the simulated  $\text{CO}_2$  signal derives from ocean ventilation impacts alone.

While the observed mean ocean B-Atm estimates suggest a significant impact on the carbon cycle, the observed changes are still too small to account for the observed  $\sim 400\%$  drop in  $\Delta^{14}\text{C}_{\text{atm}}$  across the deglaciation. This is demonstrated by model simulations that produce a mean ocean ageing of  $\sim 500$   $^{14}\text{C}$  yr ( $\sim 63\%$  of the observed value; Table 1) and yield a  $\Delta^{14}\text{C}_{\text{atm}}$  increase of only  $\sim 56\%$ , or  $\sim 14\%$  of observed (Fig. 12b, dashed grey line). Similar results have been obtained using the BICYCLE box model (Köhler et al., 2006), which produced a  $\Delta^{14}\text{C}_{\text{atm}}$  increase of only  $\sim 200\%$  at the LGM ( $\sim 50\%$  of observed), despite yielding mean ocean B-Atm offsets that match our observed values (Fig. 12b, c, dashed orange line). A mismatch was also obtained for the LGM using the CLIMBER intermediate complexity model (Ganopolski and Brovkin, 2017), where a 20% higher radiocarbon production rate, combined with reduced ocean ventilation causing a global mean B-Atm in-

crease of  $\sim 800$  years (again in line with our LGM estimate), coincided with a  $\Delta^{14}\text{C}_{\text{atm}}$  increase of only  $\sim 280\text{‰}$  ( $\sim 70\%$  of observed).

Therefore, existing radiocarbon production records cannot account for past atmospheric radiocarbon variability, either alone or in conjunction with ocean ventilation changes that are consistent with global mean B-Atm estimates (Fig. 12b, c). While there remain uncertainties in atmospheric radiocarbon reconstructions, these are likely on the order of  $\sim 1\%$  (Reimer et al., 2020), and it seems highly unlikely that  $\Delta^{14}\text{C}_{\text{atm}}$  variability over the last glacial cycle has been significantly overestimated. Similarly, it seems implausible (though clearly not impossible) that existing marine radiocarbon data significantly underestimate the magnitude of mean ocean B-Atm change since the last glacial period. The problem of closing the global radiocarbon budget since the last glacial period becomes even more difficult if volcanic/metamorphic  $\text{CO}_2$  inputs are invoked as a significant contributor to the global carbon pool during the last glacial period (Stott et al., 2019, 2009) or if no change in ocean ventilation is inferred (Stott et al., 2021).

Given the wide range of existing radiocarbon production estimates (e.g. as compiled by Dinauer et al., 2020; Fig. 12a, shaded area), it seems reasonable to postulate that existing production reconstructions might, on average, underestimate the amplitude of radiocarbon production rate change between the last glacial and the late Holocene, as suggested by a recent box-model study (Köhler et al., 2022). A “polar bias” in radiocarbon production records derived from ice-core  $^{10}\text{Be}$  fluxes has indeed been recently quantified separately for the geomagnetic and heliomagnetic modulations of cosmogenic production (Adolphi et al., 2023). Nevertheless, it has also been noted that correcting for the identified long-term geomagnetic bias in  $^{10}\text{Be}$  would not eliminate the mismatch between observed and modelled  $\Delta^{14}\text{C}_{\text{atm}}$  at the LGM (Adolphi et al., 2023). Our idealized simulations with prescribed  $\Delta^{14}\text{C}_{\text{atm}}$  allow us to infer the radiocarbon production changes that would be needed to reconcile imposed ventilation and carbon cycle changes with observed atmospheric  $\Delta^{14}\text{C}_{\text{atm}}$  (Fig. 12b, blue line. The rapid deglacial fluctuations in production that are inferred in this way (e.g. across HS1, the BA, and YD) almost certainly reflect biases due to transient carbon cycle and ocean ventilation changes that have not been implemented in our idealized simulations (see Sect. 2.5). However, the longer-term trend in inferred production rates indicates levels at the LGM that are close to the high end of the existing range of estimates (Fig. 12a, shaded area). A similar result has recently been obtained using the BICYCLE box model (Köhler et al., 2022). The higher radiocarbon production rates that are inferred at the LGM would have had a significant impact, possibly accounting for the majority of the deglacial atmospheric radiocarbon signal (Fig. 12b, dashed blue line). The implication of these results is that a parallel closure of the radiocarbon and carbon cycles since the last glacial period might yet be obtained by

exploiting the plausible range of reconstructed radiocarbon production rates (Köhler et al., 2022). Our global average B-Atm estimates provide a useful new constraint for achieving this goal.

## 5 Conclusion

We present spatial interpolations of compiled radiocarbon data for a suite of time slices spanning the last deglaciation. The primary purpose of these interpolations is to derive global average B-Atm estimates. A clear trend in global average B-Atm offsets is apparent since the LGM (when B-Atm was  $\sim 760\text{ }^{14}\text{C yr}$  higher than today on average), demonstrating unambiguous changes in the partitioning of radiocarbon between the ocean and atmosphere across the last deglaciation.

The spatial interpolations cohere with previous studies in indicating a stepwise and spatially heterogeneous rejuvenation of the ocean interior across the last deglaciation, and in suggesting the operation of a “ventilation seesaw” between the North Atlantic and both the North Pacific and Southern Ocean, especially during HS1 the YD.

A comparison of surface-, shallow-, and deep-water B-Atm trends indicates that transport changes in the upper ocean across the last deglaciation were likely modest (Marchal and Zhao, 2021a) and that B-Atm changes in the upper ocean ( $< 2\text{ km}$ ) were more strongly influenced by gas-exchange efficiency changes at high latitudes. In contrast, a more significant contribution from transport and/or water mass geometry is apparent in the deeper ocean ( $> 2\text{ km}$ ).

The time-slice reconstructions emphasize a widespread drop in B-Atm at the onset of the BA, resulting in a global average B-Atm within  $\sim 100\text{ }^{14}\text{C yr}$  of modern global average. However, model sensitivity tests indicate that a portion of this B-Atm drop could have resulted from atmospheric radiocarbon dynamics that were independent of ocean ventilation (radiocarbon production, terrestrial carbon release, etc.). The exact magnitude of this effect cannot yet be quantified, but a maximum bias of  $\sim -190\text{ }^{14}\text{C yr}$  relative to the LGM is estimated. Such a bias would imply that mean ocean B-Atm at the BA underestimates the true “ventilation age” of the ocean.

Model sensitivity tests further suggest a direct relationship between global average B-Atm anomalies and atmospheric  $\text{CO}_2$  change, with a tentative average sensitivity of  $\sim -6.3\text{ ppm CO}_2$  per  $100\text{ }^{14}\text{C yr}$ . On this basis, our global average B-Atm estimates would imply a non-negligible contribution to atmospheric  $\text{CO}_2$  change across the last deglaciation (perhaps as much as  $> 30\%$  of the total observed). Global average B-Atm estimates also suggest that any ventilation contribution to atmospheric  $\text{CO}_2$  change was concentrated during HS1, the BA, and the YD, and was largely exhausted by the onset of the Holocene.

While our results serve to underline, and tentatively to quantify, the ocean's role in deglacial carbon cycle change, they also demonstrate that a complete closure and reconciliation of the radiocarbon and carbon cycles since the last glacial remains to be achieved. Our results point to the possibility that, on average, existing reconstructions may tend to underestimate radiocarbon production rates during the last glacial period. Further work to improve the accuracy of past radiocarbon production rate estimates therefore emerges as a priority.

**Data availability.** Data lodged with PANGAEA include the compiled radiocarbon data and regional splines (<https://doi.org/10.1594/PANGAEA.960693>, Skinner et al., 2023a), netCDF files of zonally averaged time-slice interpolations (<https://doi.pangaea.de/10.1594/PANGAEA.962525>, Skinner et al., 2023b), and global average B-Atm estimates derived for each time slice from the global interpolations (<https://doi.pangaea.de/10.1594/PANGAEA.962526>, Skinner et al., 2023c).

**Video supplement.** Animation of zonal average time-slices spanning the last deglaciation, derived from global interpolations of compiled radiocarbon data (<https://doi.org/10.5446/63113>, Skinner, 2023).

**Supplement.** The supplement related to this article is available online at: <https://doi.org/10.5194/cp-19-2177-2023-supplement>.

**Author contributions.** LS designed the study, compiled and processed the radiocarbon data, and performed the interpolations with the assistance of FP. FP developed the interpolation code, and LS and FP analysed the interpolation outputs. Numerical model runs using Bern3D were performed by AJT and analysed by AJT, FJ, and LS. LS wrote the manuscript with input from all co-authors.

**Competing interests.** The contact author has declared that none of the authors has any competing interests.

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