

# Association of Optical Tracklets From a Geosynchronous Belt Survey via the Direct Bayesian Admissible Region Approach

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## Abstract

The direct Bayesian admissible region approach is an a priori state free measurement association and initial orbit determination technique for optical tracks. In this paper, we test a hybrid approach that appends a least squares estimator to the direct Bayesian method on measurements taken at the Zimmerwald Observatory of the Astronomical Institute at the University of Bern. Over half of the association pairs agreed with conventional geometric track

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correlation and least squares techniques. The remaining pairs cast light on the fundamental limits of conducting tracklet association based solely on dynamical and geometrical information.

*Keywords:* Space situational awareness; too-short arc problem; optical observations; admissible region; probability density function

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## 1. Introduction

In space situational awareness (SSA), the vast majority of observations of objects beyond low-Earth orbit are made by optical sensors, which measure a time history of angles called “tracklets” for a given object (DeMars et al., 2009; Tommei et al., 2007; Maruskin et al., 2009). The range and range-rate, however, remains largely unconstrained, and thus multiple tracklets must be combined in order to obtain a full 6-dimensional state estimate. For short-arc observations common in survey-type observations, this task is not trivial as a large subset of the state space is consistent with any given tracklet pair. Therefore, traditional initial orbit determination (IOD) techniques often perform poorly giving rise to false correlation results and unrealistic state estimates.

The direct Bayesian admissible region approach proposed by Fujimoto and Scheeres is an *a priori* state free measurement association and IOD technique (Fujimoto and Scheeres, 2012a). Given a tracklet, a compact region in the range / range-rate space is defined based on a set of physical constraints such that all likely and relevant orbits are contained within it. The admissible region (AR) is a uniform probability density function (pdf) whose support is the aforementioned compact set (DeMars and Jah, 2012). Multiple ARs may be propagated to a common epoch and an *a posteriori* pdf computed based on Bayes’ rule. Such a direct approach is feasible because the ARs are well approximated as 2-dimensional manifolds in a 6-dimensional space, making the problem sparse. Furthermore, from the Theory of General Position, two ARs do not intersect generically and thus a non-zero *a posteriori* pdf is, in almost all cases, indicative that their corresponding tracklets are associated. IOD is achieved by examining the support of the *a posteriori* pdf. Therefore, the rationale for the measurement association and the IOD are separate, allowing for both processes to be robust to outliers without the need for excessive parameter tuning.

## 2. Background

In this section, necessary concepts are introduced, such as the too-short arc problem and tracklet association with the direct Bayesian admissible region. Next, observation capabilities at AIUB as well as the current procedure to process observations are discussed.

### 2.1. The Too-Short Arc Problem

Optical observations of resident space objects (RSOs) only contain angular information regarding the observed objects' states; that is, per observation, the range and range-rate remain largely unconstrained. Consequently, orbit determination has traditionally been conducted with some type of batch or sequential estimation algorithm, whose *a priori* information is supplied via geometric techniques known as initial orbit determination (IOD) (Tapley et al., 2004; Vallado, 2007). Here, the association of observations must be assumed initially and then deduced from the quality of the least-squares fit; that is, the association of observations is a direct function of the quality of the orbit estimation and vice versa. This approach becomes problematic especially in a survey-type observation strategy. Usually, only a limited number of observations are available per night per object, each over short observation arcs, or *tracklets*, that span a few minutes (Maruskin et al., 2009). Given such a small window of data, a large subset of the state space remains consistent with each tracklet, leading to poor convergence to the true solution if not divergence. The association of tracklets, therefore, cannot be inferred confidently.

Figure 1 shows the time history of the residuals in the angular variables when fitting two tracklets of a geostationary (GEO) satellite to its true state (Object 1) as well as a fictitious state (Object 2) separated by at least 270 km but still consistent with the observations. Table 1 shows the root mean square (RMS) of the residuals. Each tracklet spans 2 minutes and consists of 5 observations; the two tracklets are separated by 20 minutes. A 2 arcsec  $1\text{-}\sigma$  Gaussian noise is added to the observed angles. The Keplerian orbit elements ( $a$  [km],  $e$ ,  $i$  [deg],  $\Omega$  [deg],  $\omega$  [deg],  $M$  [deg]) for each object at the simulation epoch is as follows:

**Object 1** (42164.154,  $10^{-6}$ , 0.1, 0, 0, 0)

**Object 2** (41079.037, 0.01965, 0.1060,  $-18.29$ ,  $-161.8$ ,  $-179.9$ ).

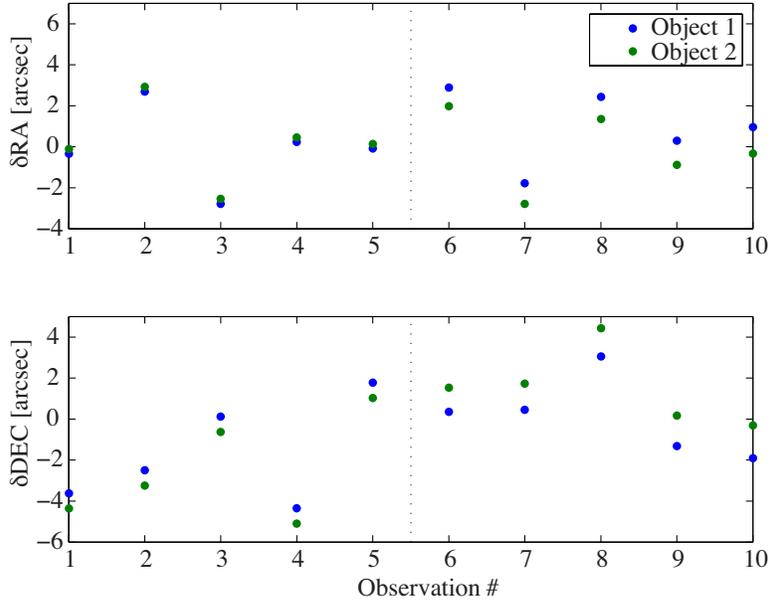


Figure 1: History of residuals in right ascension ( $\delta RA$ ) and declination ( $\delta DEC$ ) for simulated observations of a GEO object fit to its true state (Object 1) and a consistent but false state (Object 2).

We find that, in this situation, the residuals give us no indication which state is more likely.

Furthermore, in order to derive tractable geometric relationships between line-of-sight vectors, a simplistic dynamical model must be incorporated in the IOD. For example, the orbit may be assumed to be circular or the Earth’s gravity field may be considered as a point mass (Früh et al., 2009a). The former fails to incorporate eccentric orbits such as those in a geostationary transfer orbit (GTO) or high area-to-mass ratio (HAMR) objects (Rosengren and Scheeres, 2012). The latter, although valid for celestial bodies that are predominantly influenced by gravity, is less effective for RSOs which experience many perturbing forces including atmospheric drag, irregularities of the central body, and solar radiation pressure, just to name a few.

These difficulties in the association of optical tracklets of RSOs as well as the subsequent orbit determination are referred to as the *too-short arc* (TSA) problem (Tommei et al., 2007). A similar problem, albeit in longer time scales, has been studied for heliocentric orbits; in fact, the method discussed in this paper was originally devised for the astrometry of celestial

Table 1: The RMS of the residuals in right ascension ( $\delta$ RA) and declination ( $\delta$ DEC) for simulated observations of a GEO object when fit to its true state (Object 1) and a consistent but false state (Object 2).

	Object 1	Object 2
$\delta$ RA [arcsec]	1.83405	1.72122
$\delta$ DEC [arcsec]	2.37513	2.86219

bodies (Milani et al., 2004; Milani and Knežević, 2005). The more general problem of multiple target tracking using bearing-only sensors continues to be tackled in the filtering community, but most solutions require a reference state, a Gaussian assumption on the error distribution, or great computational power (Reid, 1979; Gustafsson et al., 2002).

## 2.2. The Direct Bayesian Admissible Region Approach

Various methods applying the *admissible region* (AR) concept to the TSA problem for RSOs have been studied in recent years (Maruskin et al., 2009; Tommei et al., 2007; Farnocchia et al., 2010; DeMars and Jah, 2012). In this paper, we define the AR as a pdf constrained in the range  $\rho$  and range-rate  $\dot{\rho}$  directions via a few physical criteria such as that the orbit is elliptic, the object’s range is within the sensing capabilities, and so on (Fujimoto and Scheeres, 2012a). The angle and angle-rate, nominally in right ascension  $\alpha$  and declination  $\delta$ , at the epoch of a tracklet may be estimated via a least-squares fit of the tracklet data to a polynomial model in time. These variables plus necessary parameters, such as the latitude  $\phi$  and longitude  $\Theta$  of the observation point, are referred to collectively as the *attributable vector* (Maruskin et al., 2009). Thus, each point on the AR combined with the attributable vector corresponds one-to-one with a state that the observed object may have taken. Furthermore, the covariance from the least-squares fit may be incorporated in the AR to represent observational errors.

Suppose that, given some set of criteria  $\mathcal{C}$ ,  $A$  is a compact set in state space  $\mathcal{X}$  that meet  $\mathcal{C}$ . Then, the AR  $F_{\mathcal{C}}[\mathbf{X}(t^0); \mathfrak{Y}^0]$  is a pdf over  $\mathcal{X}$  assigned to an attributable vector  $\mathfrak{Y}^0$  such that the probability  $p$  that the observed object exists in region  $B \subset A$  at time  $t^0$  is

$$p[\mathbf{X}(t^0)] = \int_B F_{\mathcal{C}}[\mathbf{X}(t^0); \mathfrak{Y}^0] dX_1^0 dX_2^0 \dots dX_n^0, \quad (1)$$

where  $\mathbf{X}(t^0) \in \mathcal{X}$  and

$$\mathbf{X}(t^i) \equiv \mathbf{X}^i = (X_1^i, X_2^i, \dots, X_n^i). \quad (2)$$

Note that we impose  $\int_A F_{\mathcal{C}}[\mathbf{X}(t^0); \mathfrak{Y}^0] d\mathbf{X}^0 = 1$ . Figure 2 is an example of an AR; here, as well as in the main analysis of this paper, the criteria are

$$\mathcal{C} = \bigcap_{i=1}^4 \mathcal{C}_i, \quad (3)$$

and

$$\mathcal{C}_1 = \{(\rho, \dot{\rho}) : E \leq 0\} \quad (4)$$

$$\mathcal{C}_2 = \{(\rho, \dot{\rho}) : 1.03 \leq \rho \leq 8.53, -5 \leq \dot{\rho} \leq 5\} \quad (5)$$

$$\mathcal{C}_3 = \{(\rho, \dot{\rho}) : 1.03 \leq r_p\} \quad (6)$$

$$\mathcal{C}_4 = \{(\rho, \dot{\rho}) : r_a \leq 15\}. \quad (7)$$

where  $E$  is the specific geocentric energy of the particle, and  $r_a$  and  $r_p$  are the apoapsis and periapsis radii of the orbit, respectively. Units of length are in Earth radii and time in hours. The different shadings represent the different regions which satisfy each criterion in set  $\mathcal{C}$ ; thus, the admissible region is where all types of shading overlap, or the region outlined by the black line. These criteria ensure that the AR encompasses most trackable object relevant to SSA while simultaneously filtering out highly eccentric orbits. Note that changing  $\mathcal{C}$  allows one to be explicit about the types of orbits that are included in the analysis. For example, if the observer is only interested in identifying objects in and near the GEO belt,  $\mathcal{C}$  may be modified to

$$\mathcal{C}_1 = \{(\rho, \dot{\rho}) : E < 0\} \quad (8)$$

$$\mathcal{C}_2 = \{(\rho, \dot{\rho}) : 5 \leq \rho \leq 8, -1 \leq \dot{\rho} \leq 1\} \quad (9)$$

$$\mathcal{C}_3 = \{(\rho, \dot{\rho}) : 5 \leq r_p\} \quad (10)$$

$$\mathcal{C}_4 = \{(\rho, \dot{\rho}) : r_a \leq 8\}. \quad (11)$$

The AR expresses our *limited* knowledge regarding  $\rho$  and  $\dot{\rho}$  which are not directly observed. In conventional filtering, pdfs of the observations only describe the error in the attributable vector and are realized in the state space as likelihoods. For underdetermined systems, the integral of the likelihood

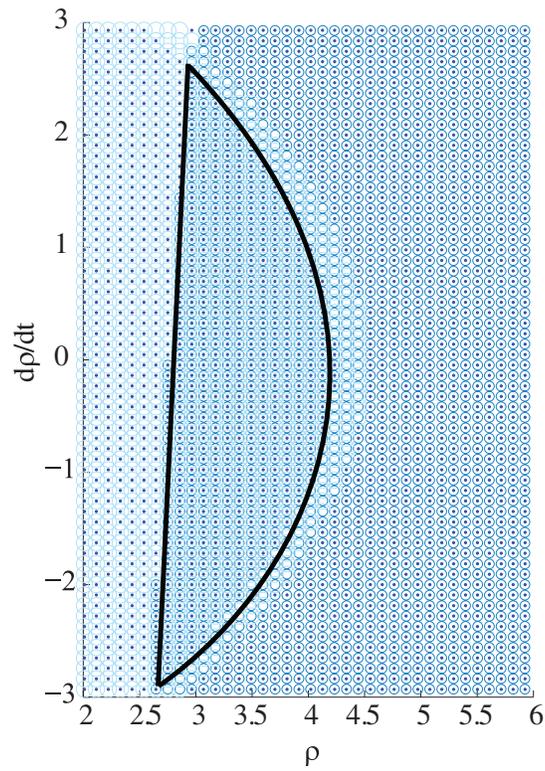


Figure 2: An admissible region for attributable vector  $\mathfrak{A} = (\alpha, \delta, \dot{\alpha}, \dot{\delta}, \phi, \Theta) = (118.26$  deg,  $-13.62$  deg,  $29.06$  arcsec/sec,  $3.75$  arcsec/sec,  $5.73$  deg,  $275.02$  deg).

function over the state space is divergent as we gain no information from the observations in coordinate directions corresponding to the variables not directly observed. We realize, however, that knowledge in these directions is not completely *lacking* for many real-world systems as the likelihood function may suggest. That is, we may add physical constraints  $\mathcal{C}$  to the observed object's state such that we define a compact pdf  $F$  still representative of all relevant states.

As a result, we may apply Bayes' rule directly to ARs in a common state space and at a common epoch  $\tau$ ; no reference state is required. To obtain the posterior pdf  $h[\mathbf{X}(\tau)]$  based on two ARs  $F_{\mathcal{C}}[\mathbf{X}^1; \mathfrak{Y}^1]$  and  $F_{\mathcal{C}}[\mathbf{X}^2; \mathfrak{Y}^2]$ ,

$$h[\mathbf{X}(\tau)] = \frac{\{\mathcal{T}(\tau, t_1) \circ F_{\mathcal{C}}[\mathbf{X}^1; \mathfrak{Y}^1]\} \{\mathcal{T}(\tau, t_2) \circ F_{\mathcal{C}}[\mathbf{X}^2; \mathfrak{Y}^2]\}}{\int \{\mathcal{T}(\tau, t_1) \circ F_{\mathcal{C}}[\mathbf{X}^1; \mathfrak{Y}^1]\} \{\mathcal{T}(\tau, t_2) \circ F_{\mathcal{C}}[\mathbf{X}^2; \mathfrak{Y}^2]\} d\mathbf{X}}, \quad (12)$$

where  $\mathcal{T}(\tau, t^i)$  is a transformation that maps some pdf  $f(\mathbf{X}^i, t^i)$  from time  $t^i$  to  $\tau$ , and  $\mathbf{X}(\tau) \equiv \mathbf{X}$ . The domain of integration is over the entire state space. Note that, in general, any pdf may be used as input, such as density information from debris distribution models (Oswald et al., 2006). This approach is computationally feasible because each AR, ignoring observation errors, has codimension 4, making the problem extremely sparse. Furthermore, the sparseness also ensures that misassociations are highly unlikely unless the association is consistent with both the observation geometry and the dynamics (Carter, 1995). From the Theory of General Position,  $h[\mathbf{X}(\tau)] = 0$  for all  $\mathbf{X}$  generically if

$$\dim \{F_{\mathcal{C}}[\mathbf{X}^1; \mathfrak{Y}^1]\} + \dim \{F_{\mathcal{C}}[\mathbf{X}^2; \mathfrak{Y}^2]\} < \dim(\mathcal{X}), \quad (13)$$

where  $\dim(\mathcal{X})$  is the dimension of the state space. Again, ignoring observation errors,  $\dim \{F_{\mathcal{C}}[\mathbf{X}^i; \mathfrak{Y}^i]\} = 2$  so the inequality holds for  $\dim(\mathcal{X}) > 5$ . The justification of associations is not at all related to the OD quality but rather solely by the geometry of the AR maps; therefore, this method is robust with minimal tuning.

Finally, transformation  $\mathcal{T}(\tau, t^i)$  is expressed analytically by means of a special solution to the Fokker-Planck equations valid for all deterministic dynamical models. Given solution flow  $\mathbf{X}(t) = \phi(t; \mathbf{X}^i, t^i)$  to the dynamics for initial conditions  $\mathbf{X}^i$ , the pdf  $\mathcal{T}(\tau, t^i) \circ f(\mathbf{X}^i, t^i) = f(\mathbf{X}, \tau)$  is expressed as

$$f(\mathbf{X}, \tau) = f[\phi(\tau; \mathbf{X}^i, t^i), \tau] = f(\mathbf{X}^i, t^i) \left| \frac{\partial \mathbf{X}(\tau)}{\partial \mathbf{X}^i} \right|^{-1}, \quad (14)$$

where  $|\cdot|$  indicates the determinant operator.



Figure 3: Current setup of ZimSMART (Herzog et al., 2011).

### 2.3. Observation Capabilities And Processing at AIUB

The Zimmerwald observatory, located about 10km south of Bern, Switzerland, consists of several optical telescopes (Herzog et al., 2011, 2010; Früh et al., 2009b). One of them, the Zimmerwald SMAll Aperture Robotic Telescope (ZimSMART), is best suited for surveying the sky searching for RSOs. ZimSMART is used to develop an orbital elements catalogue; a photograph is given in Figure 3 its specs are listed in Table 2. Two different orbital regions are surveyed: the GEO ring and the Medium Earth Orbit region (MEO). The aim of the surveys of the GEO ring is maximum coverage of the region around the celestial equator which can be observed from Zimmerwald.

Images taken with ZimSMART are analyzed as follows. First, the right ascension and declination (RA/DEC) of each RSO in the image files are automatically extracted. For each star, which appear streak-like, the center of mass is calculated. RSOs, on the other hand, appear point-like; their centers of mass are calculated as well. As each image has a finite exposure time, the epoch of the coordinates is chosen to be the mid-exposure time. The RA/DEC of the RSOs are calculated relative to celestial bodies whose physical coordinates are known and cataloged. If the same object is detected on at least 3 images, a *tracklet* will be produced; i.e., a text file containing the observing epoch, position in right ascension and declination, and apparent

Table 2: Instrument specs for the ZimSMART telescope.

<b>Spec</b>	<b>Value</b>
Lat., Long., Alt. [deg, deg, m]	46.8772 N, 7.4652 E, 951.2
Cartesian [m]	4331306.2000, 567553.9900, 4633121.6600
System	WGS-84
Telescope mount	ASA DDM85
Telescope tube	Takahashi $\varepsilon$ -180
Aperture diameter	180 mm
Focal length	500 mm
Detector type	CCD
Detector size	4096 $\times$ 4096 pixel
Field of view	4 $^\circ$ $\times$ 4 $^\circ$
Typical readout time	7 s
Wavelengths	White light
Typical exposure time	10 s
Sensitivity	Magnitude 13.5 for 10s exposure time (1m objects in GEO)

magnitude of the object for each image.

After extracting tracklets, one has to identify the observed objects. We perform this process in three steps. First, we correlate each tracklet with the JSpOC two-line element (TLE) catalogue and an internal AIUB catalogue via positions and velocities. The complete procedures are described in detail in Fröh, et al (Fröh et al., 2009b). In the second step, the leftover tracklets are tested pairwise to check if some of them belong to the same object; if so, they are stored as combined tracklets. Tracklets, for which no other fitting tracklet could be found, remain single. This procedure reduces the amount of computations in the following step. In the last step of the object identification process, the orbital elements of objects in the AIUB internal catalogue are compared with those of the new combined and single tracklets. This method is very effective for newly detected objects with observations from only one night. More details are described in Herzog, et al (Herzog et al., 2010). The identifications via positions and velocities as well as those via orbital elements have to be confirmed by a statistical orbit determination (OD). The new tracklet is associated with an internal catalog object only if the OD is successful; i.e., if the RMS of the residuals of a least squares

batch filter is below 1.5 arcsec. Due to the tracklets being *too-short arc* and lacking dynamical information, especially when the tracklet pairs span a single night, not every Keplerian element is included in the RMS, but rather only the semi-major axis, inclination, and right ascension of the ascending node.

Surveys of the geostationary ring are executed by scanning declination stripes with fixed right ascension. These observations are taken without a priori information of any catalogue objects. For the survey from which data processed in this paper is extracted, 24 stripes are taken separated by 1 hour in right ascension. These stripes are at 0 hr, 1hr,  $\dots$ , 23 hr. Each stripe contains five fields separated in declination by the field of view, and similarly, five images are taken for each field. The declination of the lowest field depends on the known density of RSOs. The advantage of this method is that the observations can be acquired in a fully automated fashion with no human interaction. The telescope software chooses the visible fields automatically. Again, a tracklet contains a minimum of three images and a maximum of five, corresponding to the number of images taken per field. Depending on the exposure time and the number of images, a tracklet thus spans anywhere between 1  $\sim$  2 minutes.

### 3. Method

The direct Bayesian approach is applied to optical tracks taken with the ZimSMART telescope in order to ascertain its validity in real-world observing scenarios. In this section, the methodology from the observation to the orbit estimate end-product is discussed. As such, two potential difficulties are anticipated with the nominal assumption that the observation errors are small enough to be ignored. The first is that the zero-error assumption causes missed associations especially when the state space discretization is refined as in (Fujimoto and Scheeres, 2011). The extension of the two-dimensional linear map extrapolation discussed in (Fujimoto and Scheeres, 2012b) is determined to be too computationally expensive for this particular problem. The second is that the ambiguity in the number of revolutions the observed object potentially made between two tracklet pairs leads to a large number of false associations. A theoretical explanation of these fictitious solutions, or *multi-rev* solutions, is given in (Fujimoto and Scheeres, 2012a).

We begin by discussing the strategy with which tracklets are obtained; especially how the RA-DEC space and time are discretized. Next, simulation

studies on the detrimental effects of measurement error on the association process are shown. As a solution to the two difficulties mentioned previously, a “direct Bayesian + least squares” hybrid algorithm is presented.

### 3.1. Assessing The Effect of Measurement Error

The addition of observational errors would make the AR, in principle, 6-dimensional. Previous work showed that the uncertainty in the attributable vector has a small effect on the discretized admissible region map and that the 2-D assumption is justified for several days but only if the discretization of the state space is coarse ( $\sim 100$  km resolution in semi-major axis,  $\sim 5^\circ$  in mean anomaly) (Fujimoto and Scheeres, 2012a). We are interested in how appropriate the 2-D assumption is for finer state space discretizations.

We consider 27000 objects distributed uniformly in the semi-major axis ( $19134.3 \leq a$  [km]  $\leq 51024.8$ ), eccentricity ( $0 \leq e \leq 0.8$ ), and inclination ( $-90^\circ \leq i \leq 90^\circ$ ) spaces but randomly in the other Keplerian orbit elements. The observation point is also randomly chosen for each object so that the observation takes place at least  $11.5^\circ$  above the horizon. This approach ensures more control over the sample set than choosing objects directly from ARs. Each state is then displaced in the 6 coordinate directions of the inertial topocentric spherical coordinate frame one by one:  $\Delta\rho = 6.059$  km,  $\Delta\dot{\rho} = 13.71$  km/hr,  $\Delta\alpha = \Delta\delta = 1.62$  arcsec, and  $\Delta\dot{\alpha} = \Delta\dot{\delta} = 3.6 \cdot 10^{-3}$  arcsec/sec (Maruskin et al., 2009). These displacements correspond to the size of one admissible region sample subset when the sample size is  $\sim 10^5$  and when Gaussian observation errors of  $\sigma = 2$  arcsec in both  $\alpha$  and  $\delta$  are considered up to  $3\text{-}\sigma$ . The displacements are mapped statically to the Poincaré orbit element space  $(\mathcal{L}, \mathbf{l}, \mathfrak{G}, \mathbf{g}, \mathfrak{H}, \mathbf{h})$  via a linear transformation  $\Phi_S = \Phi_3 \circ \Phi_2 \circ \Phi_1$ , where

$$\Phi_1 = \frac{\partial(x, y, z, \dot{x}, \dot{y}, \dot{z})}{\partial(\rho, \dot{\rho}, \alpha, \delta, \dot{\alpha}, \dot{\delta})}, \Phi_2 = \frac{\partial(a, e, i, \Omega, \omega, M)}{\partial(x, y, z, \dot{x}, \dot{y}, \dot{z})} \quad (15)$$

$$\Phi_3 = \frac{\partial(\mathcal{L}, \mathbf{l}, \mathfrak{G}, \mathbf{g}, \mathfrak{H}, \mathbf{h})}{\partial(a, e, i, \Omega, \omega, M)}, \quad (16)$$

and  $(x, y, z, \dot{x}, \dot{y}, \dot{z})$  are the Cartesian states. The Poincaré orbit elements are a canonical counterpart to equinoctial elements. They are defined with

respect to the Keplerian orbit elements as follows:

$$\begin{aligned}
\mathfrak{l} &= \Omega + \omega + M \\
\mathfrak{L} &= \sqrt{\mu a} \\
\mathfrak{g} &= \sqrt{2\mathfrak{L} \left(1 - \sqrt{1 - e^2}\right)} \cos(\omega + \Omega) \\
\mathfrak{G} &= -\mathfrak{g} \tan(\omega + \Omega) \\
\mathfrak{h} &= \sqrt{2\mathfrak{L}\sqrt{1 - e^2} (1 - \cos i)} \cos \Omega \\
\mathfrak{H} &= -\mathfrak{h} \tan \Omega,
\end{aligned} \tag{17}$$

where  $\mu$  is the standard gravitational parameter (Vallado, 2007). Finally, the dynamics from the initial epoch  $t^0$  to some time  $\tau$  are added also linearly with a state transition matrix  $\Phi_4(\tau, t^0)$  such that the complete linear map is (Fujimoto and Scheeres, 2012a).

$$\Phi(\tau, t^0) = \Phi_4(\tau, t^0) \circ \Phi_S(t^0). \tag{18}$$

For simplicity, we assume two-body dynamics because the STM deviates from an identity matrix only by the  $\partial \mathfrak{l}(\tau) / \partial \mathfrak{L}^0$  component:

$$\frac{\partial \mathfrak{l}(\tau)}{\partial \mathfrak{L}^0} = \frac{-3\mu^2(\tau - t^0)}{\mathfrak{L}(t^0)^4}. \tag{19}$$

Figure 4 shows the magnitude of the mapped displacements for all simulation cases. The different colors represent different coordinate directions in the spherical coordinates. The simulations are indexed in lexicographical order with  $a$ ,  $e$ , and  $i$  as bases; thus, as the index increases,  $a$  monotonically increases. We find that the errors in the angles, on the average, have an order 2 to 3 smaller effect on the map compared to the other variables. Therefore, it is possible to ignore them and reduce the dimensionality of the problem from 6 to 4 in many situations. We also notice that the mapped displacement in the  $\dot{\alpha}$  and  $\dot{\delta}$  reach the same order of magnitude as those in the  $\rho$  and  $\dot{\rho}$  as either semi-major axis or propagation time increases. The map in the  $\rho$  direction, however, is nearly parallel to the map in either  $\dot{\alpha}$  or  $\dot{\delta}$  especially when propagation times are on the order of days; thus, its effect on the AR is limited.

Conversely, when propagation times are short, the angle-rate errors can displace the admissible region map in directions perpendicular to the  $\rho$  and

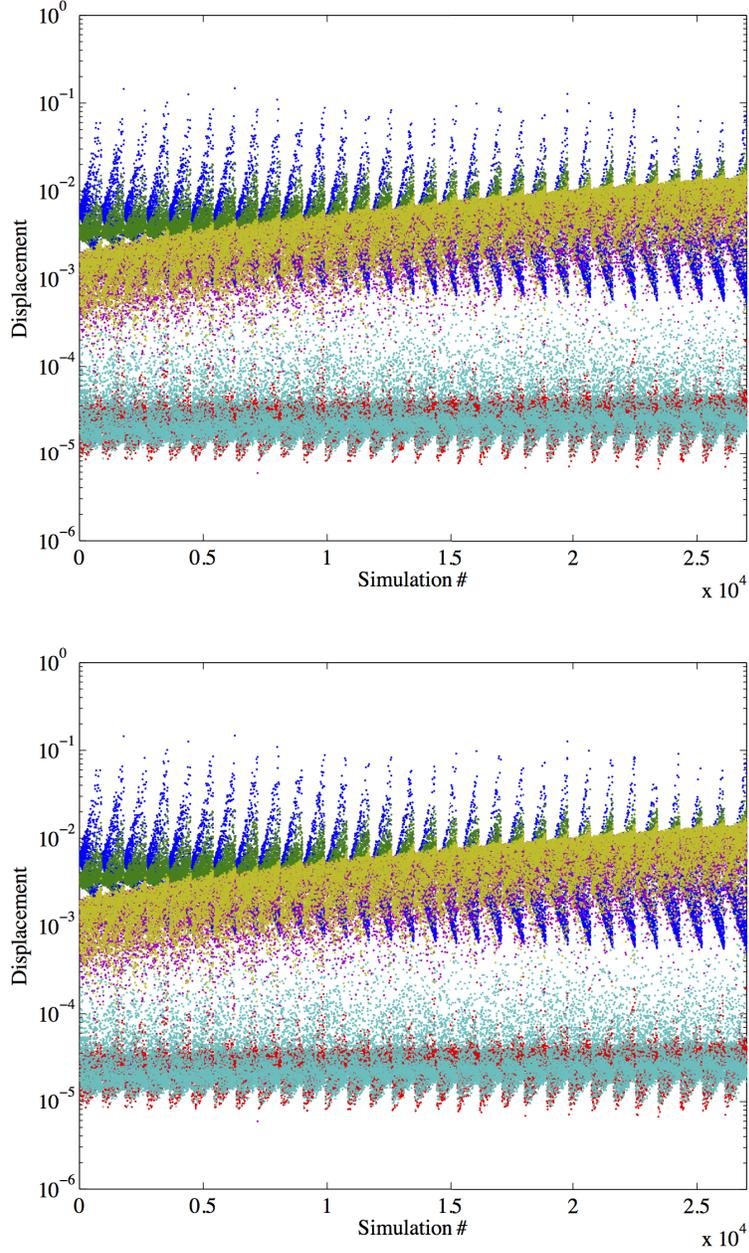


Figure 4: Norm of displacement in the  $\rho$  (blue),  $\dot{\rho}$  (green),  $\alpha$  (red),  $\delta$  (cyan),  $\dot{\alpha}$  (purple), and  $\dot{\delta}$  (gold) directions when mapped to the Poincaré orbit element space. Time propagated for 0 (top) and 120 hours (bottom). Units in Earth radii - kg - hr.

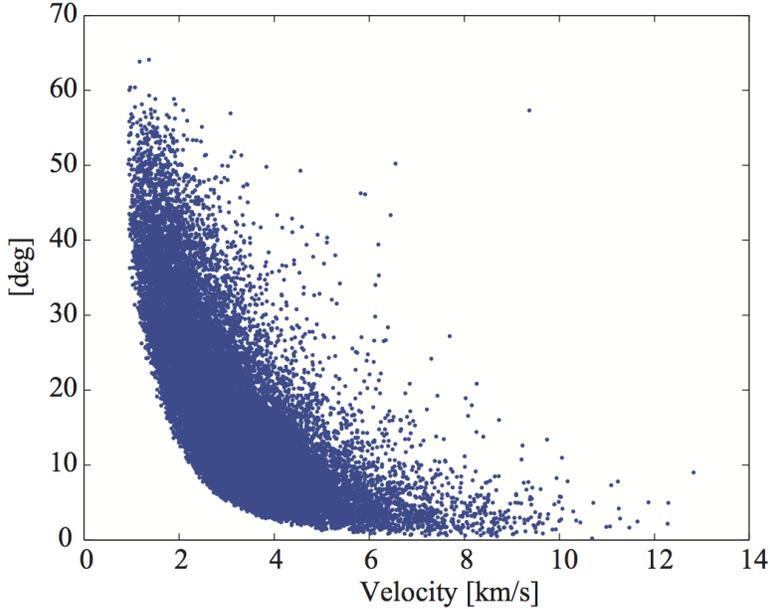


Figure 5: Scatter plot of the smaller angle between the map of the displacement in the range and either angle-rate versus the velocity of the observed object. The propagation time is 0.

$\dot{\rho}$  maps, meaning that ignoring these errors may not be justified for fine discretizations of the state space. This phenomenon is most prevalent when the velocity of the observed object is slow ( $\sim 3$  km/s and lower) as can be inferred from Figure 5. The reason why the map of the errors in  $\rho$  and the angle-rates align is because they both act to change the orbit energy and thus  $\mathcal{L}$  (equivalent to semi-major axis); the former through the orbit radius and the latter through velocity. When the velocity is low, however, even a small error in the angle-rates can strongly influence not only the orbit energy but also the orbit plane orientation.

### 3.2. Proposed “Direct Bayesian + Least Squares” Hybrid Approach

Here, we propose a hybrid approach that takes the tracklet association and initial orbit determination results of the direct Bayesian method and passes them to a least squares estimator. Although the steps in this new process are similar to those in a traditional IOD (Beutler, 2005), the justification of the association and the estimation are separated, thus improving robustness. The least squares step ensures good estimate precision without having to

use a fine discretization of the state space, minimizing negative effects of measurement error on tracklet association using ARs. Furthermore, in order to better exclude multi-rev solutions, a minimum limit  $p_{\min}$  is set to the  $p$ -value associated with the model utility test of the observed minus computed ( $O - C$ ) residuals. Suppose that for  $n$  pairs of samples  $y_1, \dots, y_n$  each associated respectively to independent variables  $x_1, \dots, x_n$ , the samples are modeled with a simple linear regression model as

$$\hat{y}_i = \hat{\beta}_0 + \hat{\beta}_1 x_i, \quad (20)$$

for all integers  $1 \leq i \leq n$ , where the hat symbolizes that it is a model estimate. Then, for the hypothesis test regarding slope parameter  $\beta_1$

$$\begin{cases} H_0 : \beta_1 = 0 \\ H_1 : \beta_1 \neq 0, \end{cases} \quad (21)$$

where  $H_0$  is the null and  $H_1$  the alternative hypothesis, the probability of falsely rejecting  $H_0$  is set to be  $p_{\min}$ . Through this step, a maximum bound is effectively set for  $\beta_1$  itself, meaning the residuals must be unrelated to time in a linear sense for a tracklet pair to be associated to a state estimate. The best estimate for the slope parameter  $\hat{\beta}_1$  is

$$\hat{\beta}_1 = \frac{n \sum x_i y_i - (\sum x_i)(\sum y_i)}{n \sum x_i^2 - (\sum x_i)^2}, \quad (22)$$

where the summation is from  $i = 1$  to  $n$ . Since the tracklets arcs are so short, we can assume that the residual bias is linear enough for us to use slope parameter  $\beta_1$  in this context. The test statistic is Student's  $t$ -distribution

$$t = \hat{\beta}_1 \sqrt{\frac{n(n-2)}{n \sum y_i^2 - (\sum y_i)^2 + \hat{\beta}_1 (\sum x_i)(\sum y_i) - n \hat{\beta}_1 \sum x_i y_i}}. \quad (23)$$

Note that a normalized statistic of the goodness of fit to a linear model is given by the coefficient of determination  $r^2$

$$r^2 = 1 - \frac{n \sum y_i^2 - (\sum y_i - \hat{\beta}_1 \sum x_i)(\sum y_i) - n \hat{\beta}_1 \sum x_i y_i}{n \sum y_i^2 - (\sum y_i)^2} \quad (24)$$

$r^2 = 1$  iff all sample pairs lie on a straight line.

We now present an outline of the hybrid algorithm. First, the time history of right ascension and declination must be converted into an attributable vector at the tracklet epoch; i.e. a single set of angles and angle-rates. The measured angles are fit to a polynomial kinematic model in time, such as for the right ascension

$$\alpha(t) = \alpha^0 + \dot{\alpha}^0(t - t^0) + \frac{1}{2}\ddot{\alpha}^0(t - t^0)^2, \quad (25)$$

where superscript 0 denotes the state at the tracklet epoch (Maruskin et al., 2009). Next, admissible regions are computed for each attributable vector in the Poincaré orbit element space. As discussed in (Fujimoto and Scheeres, 2012a), the admissible region is divided into 375,000 subsets (750 units of discretization in the range-direction  $\times$  500 units in the range-rate) and each subset linearly extrapolated. The Poincaré space, and consequently the ARs, are discretized such that the bounds of the state space are

$$\mathbf{X}_{\min} = (4.5285, 0, -3, -3, -4, -4) \quad (26)$$

$$\mathbf{X}_{\max} = (14.110, 6.2832, 3, 3, 4, 4), \quad (27)$$

where the units are in Earth radii - kg - hr. The bin size is set such that the sides are  $1.1052 \cdot 10^{-2}$  ( $\mathfrak{L}$ ),  $1.7453 \cdot 10^{-2}$  ( $\mathfrak{I}$ ),  $1.6667 \cdot 10^{-2}$  ( $\mathfrak{G}$ ,  $\mathfrak{g}$ ), and  $2.2222 \cdot 10^{-2}$  ( $\mathfrak{J}$ ,  $\mathfrak{h}$ ) for a total of  $5.2424 \times 10^{15}$  bins over the entire space. This resolution corresponds to approximately 100 km in the semi-major axis direction and 1 degree in the mean anomaly direction. The admissible regions are propagated to a common epoch, which is chosen to be the tracklet epoch of the first tracklet, under two-body dynamics. The two-body assumption is made only to simplify the problem and is not central to the direct Bayesian technique.

To avoid the high computational cost of all-on-all association, the posterior pdf  $h[\mathbf{X}(\tau)]$  based on the admissible regions is computed for tracklet pairs in reverse chronological order (i.e., for a set of  $N$  tracklets ordered by epoch, Tracklet 1 + Tracklet  $N$ , Tracklet 1 + Tracklet  $N-1$ , ...) until we find a pair for which  $h[\mathbf{X}(\tau)] > 0$  (Fujimoto and Scheeres, 2012a). We then temporarily claim these tracklets as associated and run a bank of least squares filters simultaneously to refine the fit of the measurements to the state estimate. Note that if an object catalog exists, then one should first correlate tracklets with these objects first. Also, only tracklet pairs whose epochs are separated by at least 24 hours are considered so that enough dynamical information is available.

The reference state of each filter is the centroid of each bin where  $h[\mathbf{X}(\tau)] > 0$  transformed into the J2000 cartesian space. In this paper, we assume that no *a priori* information exists; if desired, the *a priori* covariance may be set to approximate  $h[\mathbf{X}(\tau)]$ . The assumed observation error is set to 2 arcsec  $1-\sigma$ . The observation-state relationship and corresponding linear partials matrix assume a spherical Earth

$$x = \rho \cos \alpha \cos \delta \tag{28}$$

$$y = \rho \sin \alpha \cos \delta \tag{29}$$

$$z = \rho \sin \delta. \tag{30}$$

For the set of filters that converge, if

1. the RMS of the  $O - C$  residuals for both the right ascension and declination over all tracklets processed is less than some maximum  $\text{RMS}_{\text{max}}$   
AND
2. the  $p$ -value of the model utility test for both the right ascension and declination for each individual tracklet is greater than some minimum  $p_{\text{min}}$ ,

then the tracklets are confirmed to be associated and the state estimate with the smallest  $O - C$  residual RMS is added to the object catalog. In this paper,  $\text{RMS}_{\text{max}} = 0.7$  arcsec and  $p_{\text{min}} = 0.1$ ; these values are chosen to best describe the observational capabilities of ZimSMART. Finally, the next tracklet in the set is paired with other tracklets as before, and the process is repeated until all tracklets are processed. Figure 6 is an overall flowchart.

As an example of how the linear regression slope parameter can help identify multi-rev solutions, Table 3 shows the association result of two tracklets correlated with the current AIUB approach to object 98006B in the JSpOC TLE catalog. For Solution 1, which is most likely a multi-rev solution, even though the RMS of the residuals in RA over the first tracklet is smaller than the gating criterion, the coefficient of determination value suggests a strong linear relationship between the residuals over time. In addition, the small  $p$ -value indicates that it is highly unlikely to falsely infer  $\beta_1 \neq 0$  due to random chance.

#### 4. Results

In this section, the outcome of observation association via the direct Bayesian only approach and hybrid approach are compared. Then, the orbits

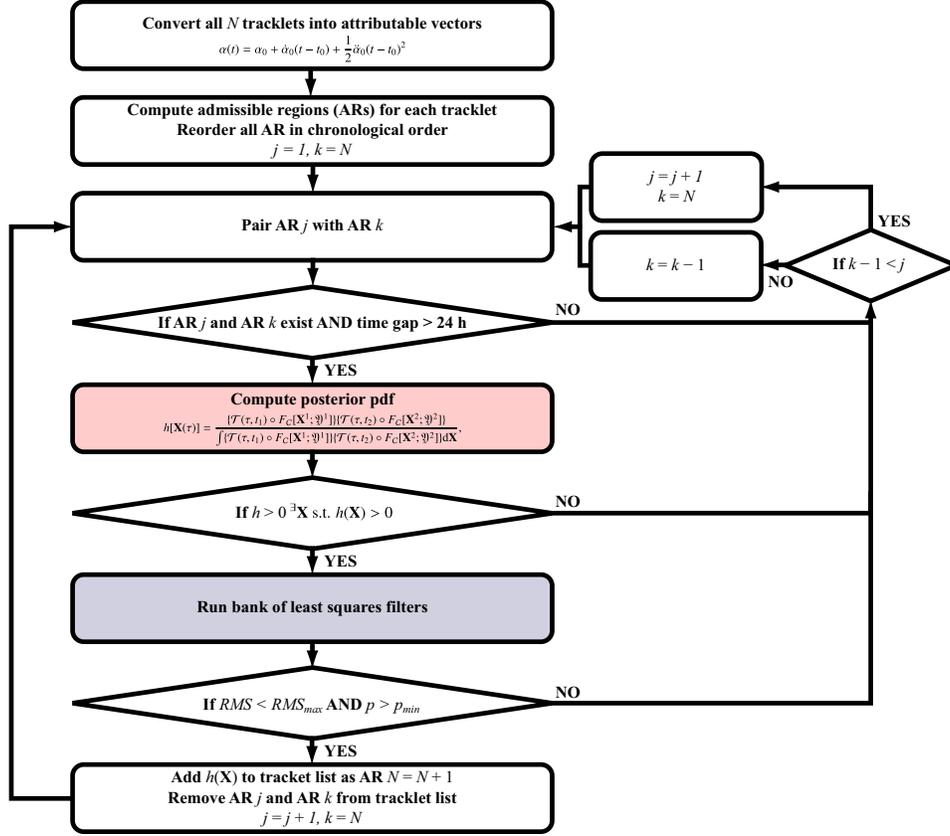


Figure 6: A flowchart of the “direct Bayesian + least squares” hybrid approach.

Table 3: For two orbit estimates computed based on the hybrid approach, the semi-major axis ( $a$ ), eccentricity ( $e$ ) and inclination ( $i$ ) are listed along with the RMS of the residuals for the first tracklet in RA ( $\text{RMS}_{\text{RA}}$ ) and the coefficient of determination ( $r_{1,\text{RA}}^2$ ) with its corresponding  $p$ -value.

	$a$ [km]	$e$	$i$ [deg]	$\text{RMS}_{\text{RA}}$ [arcsec]	$r_{1,\text{RA}}^2$	$p$
<b>Solution 1</b>	20285	0.57013	2.2261	0.24841	0.97512	0.00469
<b>Solution 2</b>	42166	0.00466	0.2886	0.05599	-0.13651	0.82673

Table 4: Parameters for the data set used in this example. # of objects detected is based on AIUB correlation results.

Parameter	Value
Epoch of Initial Field	Aug 18, 2012 22:59:08.64 UTC
Epoch of Final Field	Aug 20, 2012 02:01:32.69 UTC
Stripe Right Ascension	23 hr
Total # of Fields	55
Total # of Tracklets	212
Total # of Objects Detected	48
# of Objects Detected Twice w/ 24h Gap	19

of objects detected by the hybrid approach are examined in further detail. The observation strategy is described in Section 2.3. We process a set of tracklets taken with the ZimSMART telescope over one RA stripe; detailed measurement parameters are given in Table 4.

Table 5 is a table of all 20 associated tracklet pairs detected. The solutions can be categorized into three types.

**Type I** The solution is most likely a true solution; associates tracklets that are also similarly correlated with the AIUB code (11/20 = 55% of solutions).

**Type II** The solution is most likely a false positive (multi-rev) solution; associates tracklets that are correlated to two separate objects with the AIUB code AND the solution does not exist at or near ( $\pm 200$  km) GEO altitude (7/20 = 35%).

**Type III** The solution is most likely a new true solution; at least one of the associated tracklets are not correlated with the AIUB code AND the solution exists at or near ( $\pm 200$  km) GEO altitude (2/20 = 10%).

Table 6 compares the distribution of solution types for the direct Bayesian only and hybrid methods. Note that, for the direct Bayesian only case, the discretization is refined dynamically over the support of the posterior pdf to  $6.1449 \cdot 10^{-4}$  ( $\mathfrak{L}$ ),  $1.7453 \cdot 10^{-2}$  ( $\mathfrak{I}$ ),  $3.3333 \cdot 10^{-3}$  ( $\mathfrak{O}$ ,  $\mathfrak{g}$ ), and  $4.4444 \cdot 10^{-3}$  ( $\mathfrak{J}$ ,  $\mathfrak{h}$ ) for a total of  $5.8928 \times 10^{19}$  bins over the entire space. As expected, for the direct Bayesian only case, over 60% of the solutions found are multi-rev solutions. This ratio improves to 35% for the hybrid case; furthermore,

Table 5: Summary of association results ordered by solution type. Tracklets correlated to the JSpOC TLE catalog objects are indicated by the object’s 6 letter international designator. Tracklets associated with objects in AIUB’s internal catalog are indicated by the object’s 7 letter designator starting with “Z.” Tracklets newly associated are indicated by a bracketed number assigned by tracklet epoch.

Type	Object ID	Tracklet #1	Tracklet #2	$\delta t$ [days]
<b>Type I</b>	2	'94022A'	'94022A'	1.01325
	3	'93078B'	'93078B'	1.00508
	5	'00081A'	'00081A'	1.00508
	11	'91075A'	'91075A'	1.00387
	12	'02015B'	'02015B'	1.00172
	13	'98006B'	'98006B'	1.00172
	14	'10025A'	'10025A'	1.00172
	15	'08034B'	'08034B'	1.00172
	16	'98057A'	'98057A'	1.00172
	17	'85015B'	'85015B'	1.01285
19	'Z11003C'	'Z11003C'	1.00932	
<b>Type II</b>	4	'10032B'	'98050A'	1.00635
	6	'98050A'	'09008B'	1.08071
	7	'00054A'	'10025A'	1.00635
	9	'08065B'	'10021A'	1.01453
	10	'11041A'	'98057A'	1.00603
	18	'04008A'	'98024A'	1.03621
	20	'01042A'	'Z12230G'	1.00963
<b>Type III</b>	1	[3]	'Z12230C'	1.00635
	8	[13]	[120]	1.00635

2 additional “true” solutions are detected. If an object is indeed observed twice or more, the fact that more multi-rev solutions are rejected means that their observations are not “used up” before being properly associated. All solutions obtained with the hybrid method are listed in Appendix A. Based on these results, in a real-world scenario where hundreds of tracklets that contain measurement errors are to be associated, we recommend modifying the direct Bayesian admissible region approach so that it explicitly accounts for errors as well as reduces multi-rev solutions.

#### 4.1. Type II Solutions

We now examine in detail the 45% of hybrid method solutions that do not agree with existing techniques. Table 7 lists the semi-major axis, eccentricity,

Table 6: Orbit solutions found by type (I, II, or III) and association method.  $\Delta$  is the difference between methods, and bold numbers indicate an improvement.

	<b>I</b>	<b>II</b>	<b>III</b>	<b>Total</b>
<b>Direct Bayesian only</b>	10	18	1	29
<b>Hybrid</b>	11	7	2	20
<b><math>\Delta</math></b>	<b>+1</b>	<b>-11</b>	<b>+1</b>	<b>-9</b>

Table 7: For Type II solutions, listed here are the semi-major axis  $a$ , eccentricity  $e$ , inclination  $i$ , and ratio between the observation time gap  $\delta t$  and the orbit period  $T$ .

<b>Object ID</b>	<b><math>a</math> [km]</b>	<b><math>e</math></b>	<b><math>i</math> [deg]</b>	<b><math>\delta t/T</math></b>
4	20370.27	0.574014	2.36898	3.001
6	44103.76	0.033165	0.15711	1.013
7	20367.00	0.601097	4.21171	3.002
9	20489.57	0.566496	2.53346	2.999
10	26695.54	0.385897	1.64413	2.000
18	27240.55	0.344950	0.95749	2.001

and inclination of the Type II solutions. All objects except for Object 6 are in altitudes where it is unlikely that any objects exist. As can be inferred from the strict gating required for an association to be detected ( $\text{RMS}_{\text{max}} = 0.7$  arcsec), however, these are not degenerate solutions but rather likely to be multi-rev solutions that arise from the ambiguity of the number of revolutions made by the observed object during the observation gap. Indeed, if we are to look at ratio of the time gap between tracklets to the orbital period of the solution, they are all nearly integers between 1 and 3. It should be stressed that, without *a priori* knowledge regarding the distribution of RSOs, for instance, these Type II solutions are dynamically just as viable as the solutions given by the AIUB code. That multi-rev solutions may be solved for without the need for parameter tuning is a potential strength of the direct Bayesian AR method; if RSOs indeed exists on these orbits, they would be difficult to detect with conventional geometric techniques.

Figures 7 and 8 are graphical representations of Object 9 as well as the two catalogued objects to which the AIUB code correlated. As expected, Object 9 and the catalogued objects align along the observation direction at

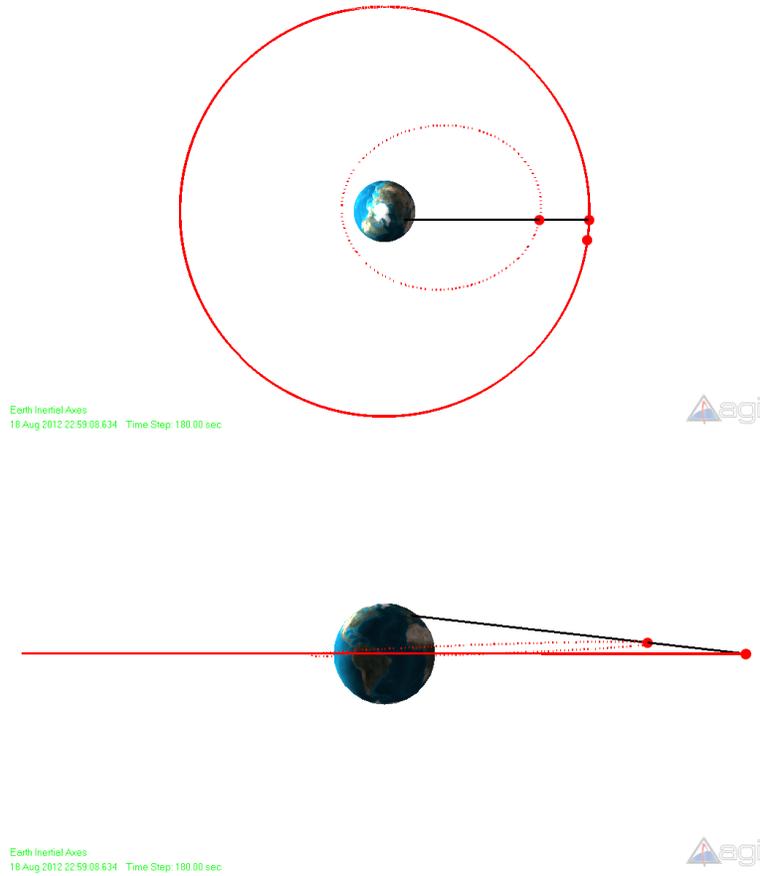


Figure 7: The orbit (dotted red line) and position of Object 9 at the first tracklet epoch along with orbits and positions of catalogued objects 08065B and 10021A (solid red lines) as well as the observation direction (black). Figures generated with AGI’s STK.

each tracklet epoch. We find that the slight inclination common to Type II solutions is necessary so that the solution appears at the same declination as an object in the GEO ring.

#### 4.2. Type III Solutions

Table 8 lists the Keplerian orbital elements of the Type III solutions found and Figure 9 is a plot of their groundtracks. Unlike the Type II solutions, both objects are very near circular at geosynchronous altitude, making it likely that they are uncatalogued objects. Here, the long time gap between

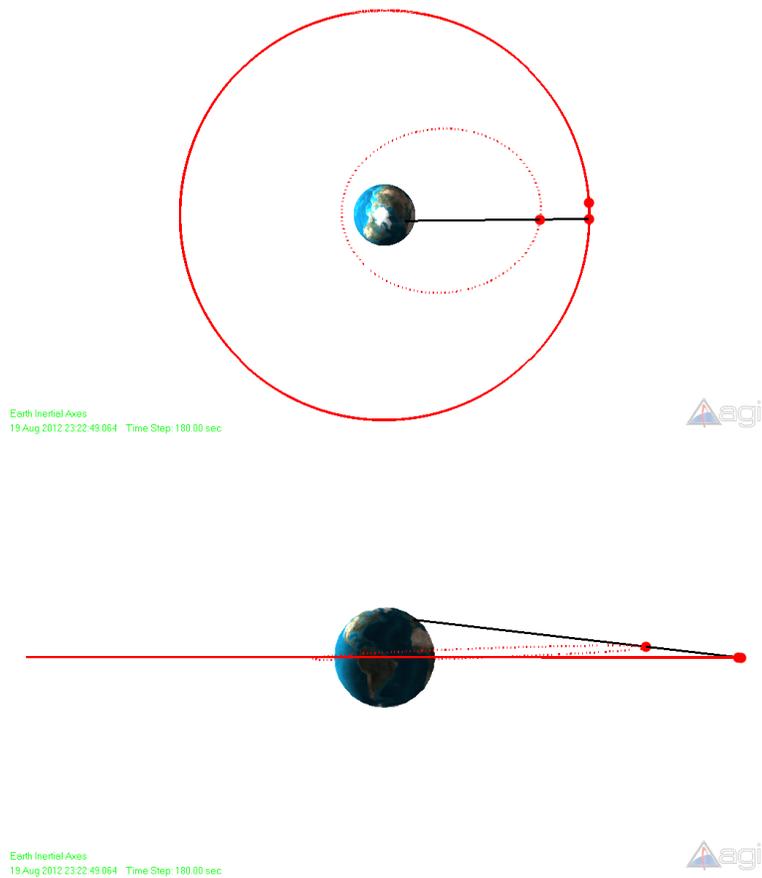


Figure 8: Similar to Figure 7 but for the second tracklet epoch. Figures generated with AGI's STK.

Table 8: For Type III solutions, listed here are all 6 Keplerian orbital elements: semi-major axis  $a$ , eccentricity  $e$ , inclination  $i$ , right ascension of the ascending node  $\Omega$ , argument of periapsis  $\omega$ , and mean anomaly  $M$ .

Object ID	$a$ [km]	$e$	$i$ [deg]	$\Omega$ [deg]	$\omega$ [deg]	$M$ [deg]
1	42167.94	0.005040	8.74794	38.5298	38.449	264.763
8	42166.74	0.000314	0.09691	56.8309	143.201	140.752

tracklets allowed the admissible regions to dynamically evolve so that previously uncorrelated tracks could be linked together.

#### 4.3. Potential Improvements of The Observation Strategy

The proposed idea of applying a least squares batch filter to the direct Bayesian probabilistic output is effective but nonetheless can still be improved. Table 9 sorts the 19 objects expected to be detected based on the AIUB correlation results by whether it matches the results from the hybrid method or not. Type II solutions are not completely rejected; consequently, about 15% of the expected objects are missed due to one or more of their tracklets being associated to a multi-rev solution. Ruling out apparent multi-rev solutions as false associations given just the two tracklets and dynamical system flow may be difficult, especially when the measurement residuals are so well behaved. Note that multi-rev solutions are not a problem in the AIUB code as the tracklets are never associated beyond a single night. This approach is not ideal either; as discussed in Section 2.1, a lack of dynamical information can also lead to poor association solutions. Indeed, new objects within the GEO belt are detected with the hybrid approach where the associated tracklets are separated by at least 24 hours.

The easiest way to reject multi-rev solutions is to conduct follow-on observations based on the estimated state. If the multi-rev solution is indeed truth, then its short orbital period relative to GEO objects should allow it to be observed multiple times per night. Alternatively, one can make better use of the information already available in the observations. Geometric correlation to JSpOC catalogued objects as implemented in the current AIUB code is, in effect, one example where *a priori* information is fused into the association process. Preconditioning the ARs with pdfs derived from debris catalogues or density models may similarly be effective; the prior has been implemented in previous work to greatly improve computational speed (Fuji-

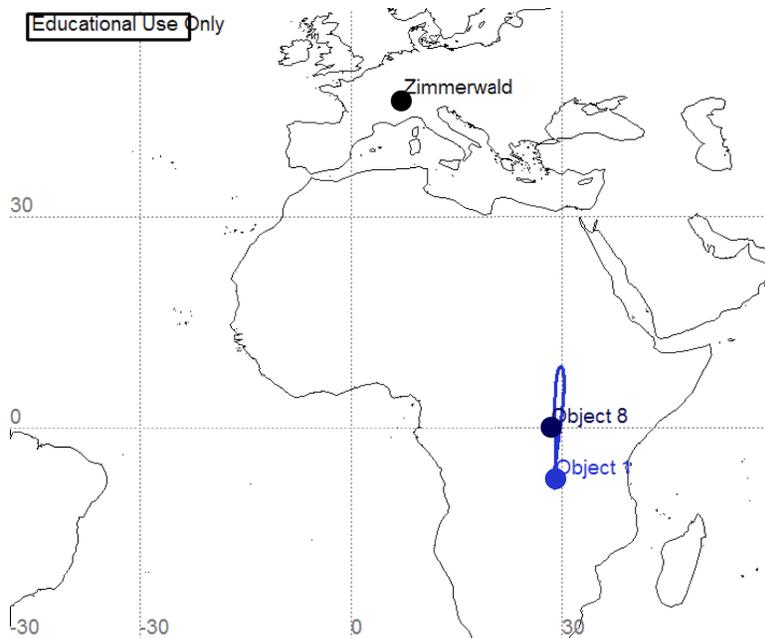


Figure 9: Ground tracks of Type III solutions for a full orbit from the tracklet epoch: Objects 1 (light blue) and 8 (dark blue). The position of each object at the epoch of the first tracklet is indicated by a dot. The position of the Zimmerwald observatory is also plotted with a black dot. Figure generated with AGI’s STK.

[moto and Scheeres, 2012a](#)). Information which would allow one to distinguish between tracklets also exists outside of the realm of dynamics, such as photometry and spectroscopy from the CCD image files ([Scott and Wallace, 2009](#); [Schildknecht et al., 2009, 2010](#)). Finally, because the evaluation of Bayes’ rule when associating ARs is an embarrassingly parallel problem, all-on-all association may be possible within a reasonable time if many computation cores are available.

Additionally, about 25% of the 19 expected objects are completely missed by the hybrid approach. The solution corresponding to object 82082A is rejected due to a particularly large  $O - C$  residual value (-2.445 arcsec) in the declination direction for one angle measurement. If this particular measurement is excluded, the RMS of the declination residuals improves from 1.0113 arcsec to 0.59865 arcsec: within the maximum RMS gate for this paper. For all of the other solutions, the  $p$ -value limit for the model utility test of the residuals is triggered most probably by mistake. Figure 10 is one such example: the residuals in the declination direction for the first tracklet

Table 9: Objects where multiple tracklets with at least a 24 hour time gap are correlated based on the AIUB algorithm. Tracklets of objects under “Agreement” are associated similarly with the hybrid approach, “multi-rev” associated with another object, and “missed” not associated at all. Number of objects in each category in parenthesis.

00081A		
02015B		
08034B		
10025A		
85015B		
91075A		
93078B		82082A
94022A		84028A
98006B	00054A	93015A
98057A	11041A	95067A
Z11003C	98050A	98056A
<b>Agreement (11)</b>	<b>Multi-Rev (3)</b>	<b>Missed (5)</b>

is “linear enough” such that  $p = 0.0432 < p_{\min} = 0.1$ . As such, there exist observation scenarios where reliably evaluating the “no linear relationship” null-hypothesis can be difficult due to the small number of individual angle measurements included in a tracklet. Increasing measurements per tracklet not only will shed better light on any biases present in the residuals but also has the added benefit of improving the angle-rate estimate in the attributable vector.

Because the theory discussed in this paper addresses the TSA problem in a much more probabilistically straightforward way than other IOD techniques, it allows one to reevaluate future observational strategies so that they minimize false positive / negative association solutions. Any changes to the current strategy, such as the time gap between tracklets, directly affects the *a posteriori* pdf in the state space without the need to assume an observation geometry, dynamical system, or type of errors accounted for. Furthermore, it is not necessary to parametrically account for multi-rev solutions, which explained all false positive solutions encountered in this paper; rather, they naturally appear in the *a posteriori* pdf as long as they are dynamically viable.

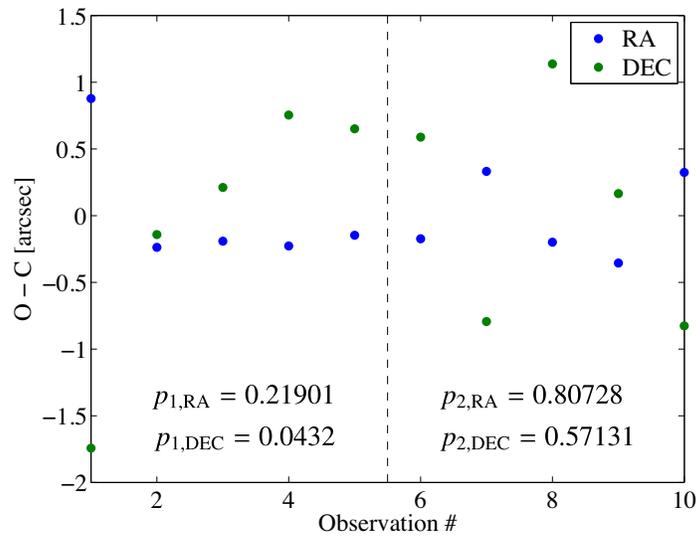


Figure 10: History of RA/DEC  $O - C$  residuals using the hybrid method for two tracklets correlated to object 84028A with current AIUB code. Plot points to the left of the dotted line are for the first tracklet and to the right the second.  $p$ -values for the model utility test are indicated per tracklet and angle.

## 5. Conclusions

In this paper, the direct Bayesian admissible region approach to short-arc association and initial orbit determination is applied to optical observations taken at the Astronomical Institute of the University of Bern. Traditional methods rely on the quality of the orbit determination to conduct observation association, which is often unreliable. The direct Bayesian approach improves robustness by leveraging the sparseness of probability distributions that describe range and range-rate ambiguity given a single optical track. Furthermore, a hybrid approach that appends a least squares batch filter is found to efficiently incorporate measurement error and reduce false positives due to multi-rev solutions. Processing a set of 212 tracklets results in 20 objects detected; 2 of which are newly detected by the proposed method. Nonetheless, about 35% of the solutions are still deemed to be false positive solutions, and thus ideas to further reject multi-rev solutions are proposed. Future work is to implement these ideas, such as increasing the number of angle measurements per tracklet, as well as further testing of the hybrid approach with more data sets.

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## Appendix A: State Estimates of Objects Detected With Hybrid Approach

Table 10: Keplerian orbital elements of all objects detected with the hybrid approach. “Tr 1” and “Tr 2” indicate the objects that the associated tracklets are correlated to by the AIUB code. “Type” is the solution type: I, II, or III.

Obj #	$a$ [km]	$e$	$i$ [rad]	$\Omega$ [rad]	$\omega$ [rad]	$M$ [rad]	Tr 1	Tr 2	Type
1	42167.9389	0.00504009	0.15268036	0.67247208	0.67106167	4.62098602	[ 3]	Z12230C	III
2	42553.9802	0.00220195	0.14845483	1.02426359	-2.9210019	1.58577527	94022A	94022A	I
3	42474.7078	0.00104508	0.05504329	1.20373515	0.46646274	4.33420972	93078B	93078B	I
4	20370.2703	0.57401358	0.04134642	-1.4298436	-2.0797513	3.30584577	10032B	98050A	II
5	42167.8135	0.00485164	0.00833985	1.41736384	-0.1245722	4.65738998	00081A	00081A	I
6	44103.7561	0.03316462	0.00274204	1.53057116	-1.8511525	6.26491311	98050A	09008B	II
7	20367.0045	0.60109671	0.07350824	-2.8500348	-0.3679198	2.27664978	00054A	10025A	II
8	42166.739	0.00031368	0.00169132	0.99188551	2.49933109	2.45658423	[ 13]	[ 120]	III
9	20489.5672	0.56649625	0.04421714	-2.3842792	-0.9701911	2.77169499	08065B	10021A	II
10	26695.5358	0.38589704	0.02869554	-1.0047144	-2.714267	3.61614499	11041A	98057A	II
11	42358.4502	0.00639731	0.12502374	1.02260295	-2.8554889	1.46486536	91075A	91075A	I
12	42166.1451	0.00504257	0.00645396	1.49216639	2.92711687	1.43573959	02015B	02015B	I
13	42165.7048	0.00465978	0.00503632	-0.1067257	-1.8389239	1.53631869	98006B	98006B	I
14	42166.3074	0.00442969	0.00214473	2.01311262	2.3699257	1.51251347	10025A	10025A	I
15	42165.9793	0.00151075	0.00235757	1.3716077	3.06449675	1.4651209	08034B	08034B	I
16	42166.089	0.00093822	0.00177328	1.16793164	0.5203132	4.20931434	98057A	98057A	I
17	42369.2386	0.00313633	0.23807355	0.57337329	0.69360244	4.30588617	85015B	85015B	I
18	27240.5471	0.34495023	0.01671134	-1.7512588	-1.6627614	2.29904061	04008A	98024A	II
19	42359.0222	0.01589573	0.18605515	0.58086998	-2.3777622	0.80333633	Z11003C	Z11003C	I
20	26755.2517	0.38458146	0.03241271	-0.879943	-2.7948581	2.30041338	01042A	Z12230G	II