

inversion and is self-sustaining. Upon gelation, a dramatic colour change from clear red to dark blue was observed. The effect of solvent on gelation was tested. The critical gel concentration of **1** in different solvents when combined with CuCl₂ in ethanol is as follows: chloroform, 0.02 M; dichloromethane, 0.04 M; chlorobenzene, 0.03 M; and *o*-dichlorobenzene, 0.04 M. The gelation was not effective in non-halogenated solvents.

To further explore the role of metal salts, different metal cations (Table S1, ESI†) and counterions (Table S2, ESI†) were chosen for the gelation studies. The result was that MX₂ with M(II) = Cu, Cd, Co (X = Cl⁻, Br⁻) can form dark blue gels, but not in the cases of Zn(II), Mn(II) and Fe(II). The anion clearly plays a role in the gelation, as fluoride and perchlorate copper salts do not form gels. The halide ions show high propensity to form gels, very probably ascribed to their ability to form metal clusters *via* halide bridges,¹⁰ and here we focus on these salts.

It has to be noted that the metallogels exhibit high stability as evidenced by the fact that the sol-gel transformation cannot be tuned by external chemical stimuli. For example, differential scanning calorimetry (DSC) measurements of the CuCl₂ xerogel (Fig. S1, ESI†) indicate three irreversible endothermic transitions, very probably due to the loss of solvents and the solid-state-like phase transition. It appears that the gel is thermo-irreversible as it remains intact when heated up to 70 °C while it is gradually fragmented upon increasing the temperature, finally leading to a dark blue precipitate in solution. It is neither responsive to a metal trapping reagent – such as ethylene diamine tetraacetic acid (EDTA) – nor to oxidation. Additionally, the reversibility of the gelation process did not occur through mechanical stimuli either. When shaking breaks a gel, the appearance of the system becomes one that is composed of fibres dispersed in a solvent. This effect is irreversible. All these results demonstrate that metallogels possess a rigid rather than flexible suprastructure.

The UV-visible absorption spectra of **1** in a CHCl₃ solution and as a film as well as of a CuCl₂ gel are shown in Fig. 2. According to our previous results of analogues,¹¹ the intense absorption band peak at 488 nm in solution corresponds to an intramolecular charge-transfer transition from the HOMO localized on the TTF core to the LUMO localized on the BTM moiety. A slight

red-shift of this optical absorbance in the film is attributed to strong intermolecular interactions. The copper chloride gel shows an even larger red-shift extending into the near-IR region giving rise to a low bandgap (<1.2 eV), indicative of the coordination of Cu(II) by the nitrogen atoms on the BTM portion of **1**. These results are consistent with the ¹H NMR data of the CdBr₂ gel (Fig. S2, ESI†). Upon gelation, the signals due to the benzene ring and propyl groups of **1** become broad and are shifted upfield. Additionally, FTIR spectra of **1** and its MCl₂ xerogels (Fig. S3, ESI†) indicate slight shifts in the vibration frequencies for the BTM aromatic ring in the range of 1200 cm⁻¹ and 1500 cm⁻¹ upon gelation. All these observations imply that the coordination bonds are primary driving forces to form the fibers while the π-π stacking between BTM-fused TTF units and non-covalent van der Waals interactions play an important role in the entanglement of fibers and the capture of the solvent.

Powder X-ray diffraction patterns of a xerogel from CuCl₂ and **1** as well as of monomer **1** on its own are shown in Fig. S4 (ESI†). The gel shows some amorphous as well as crystalline character, as evidenced by the presence of both broad and sharp peaks. The gel spectrum exhibits some similarity to **1** but also new features that do not correspond to either pure CuCl₂ or **1**, indicating that the self-organization in the gel is not simply a precipitated monomer, but a new hybrid of both metal ions and **1**.

Magnetic susceptibility (SQUID) measurements (Fig. S5, ESI†) also determined that there was 8.3% of copper(II) in the xerogel, which is close to the expected value of 11% for a 1:1 ratio of **1** to copper(II) in the bulk gel. Moreover, the magnetic susceptibility data show that there is a pronounced antiferromagnetic coupling between neighbouring Cu(II) ions. These results indicate that Cu(II) is involved in the gelation and not just a precipitate dispersed in the gel matrix.

Scanning electron microscopy (SEM) was used to visualize the morphology of the xerogels formed under different conditions. As depicted in Fig. S6 (ESI†), the gels have a wavy suprastructure in less polar solvents, such as chloroform and chlorobenzene, and more flat structures in more polar solvents, such as *o*-dichlorobenzene and dichloromethane. The 3-D network has a woven appearance with belt-like fibers of about 50 nm in diameter, which shows some alignment over several tens of microns. Similar features, albeit with some subtle differences, were observed in the SEM images of other metallogels (Fig. S7, ESI†).

The gels show an intriguing response to washing (Fig. S8, ESI†). The material washed away depends upon the washing solvent. The dark blue xerogel is initially rinsed with ethanol, which washes away the metal salt and changes the colour to that of **1**, red. This observation can be confirmed by energy-dispersive X-ray spectroscopy (EDX). For instance, the atomic ratio of S:Cu changes from 70:10 for the as synthesized gel to 96:2 for the gel washed with ethanol. Also SEM images reveal the variation of the fibrous nature of the gel by washing (Fig. S9, ESI†). The as synthesized gel shows a columnar structure with overlying parallel wavy features. When the gel is washed with ethanol, however, what remains appears very similar to pure **1**, suggesting the removal of CuCl₂. When the gel is further rinsed with dichloromethane, the layers of **1** are removed from the gel surface, and the gel turns dark blue. This process of selectively

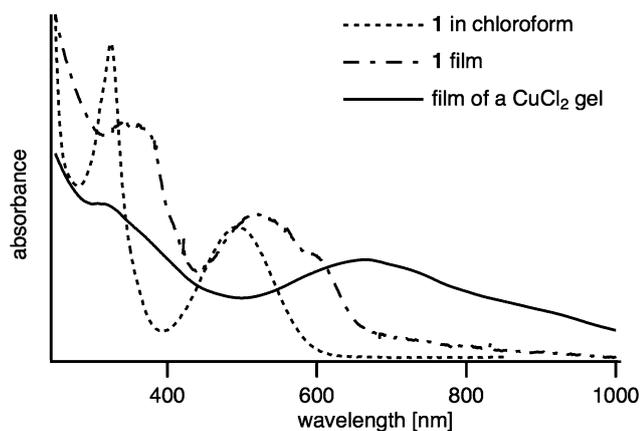


Fig. 2 UV-vis spectra of **1** in solution, film, and as a CuCl₂ gel.



washing away portions of the gel by alternating solvents is indicative of the formation of a stable layered suprastructure (Fig. S10, ESI[†]) with a lamellar structural sequence during the gelation process. This structural feature is not beneficial to the penetration of the solid EDTA into the gel, which accounts for the aforementioned failure of EDTA to uptake Cu(II) from the gel.

While bulk gels are interesting, electronically active materials require easy processability into thin films for useful device manufacture. To that end effort was focused on creating homogeneous thin films. Ideally, spin coating or drying over a thermal gradient would be a favourable method for it, however, due to the gel-like consistency of the material we probed another route. Loose gels, made from a mixture of a solution of **1** in *o*-dichlorobenzene and a solution of CuCl₂ in ethanol, were sandwiched between two microscope slides, which were then pressed together to blend the green gelatinous particulates (visible under an optical microscope) into an amorphous green film. The glass slides were then slid laterally apart. The resulting films were allowed to dry, whereupon they darkened. When viewed under a polarizing optical microscope, a homogeneous film of uniform alignment responded to the polarized light, indicating the parallel alignment of the chromophoric units. A SEM image of the resulting film (Fig. S11, ESI[†]) shows that there are features on the scale of 500 nm to 1 μm.

A doped thin film of the CuCl₂ xerogel was prepared after exposure to iodine vapors for two minutes and further investigated as a conductor. Electron paramagnetic resonance (EPR) spectroscopy is useful for exploring the electronic characteristics of this kind of material.^{6d} As illustrated in Fig. S12 (ESI[†]), the un-doped xerogel is EPR signal silent whereas the doped one exhibits a clear EPR signal with a peak-to-peak linewidth (ΔH_{pp}) of 30 G. This EPR signal corresponds to the formation of the cation radical of **1**.¹² Surprisingly, when the doped xerogels were subjected to an annealing cycle, an irreversible decrease of the peak-to-peak linewidth was observed from 30 G to 6 G. During the annealing process (Fig. S13, ESI[†]), the samples were heated up to 200 °C and subsequently cooled down to room temperature. This change in ΔH_{pp} can be attributed to an irreversible solid-state-like transition to a new structural phase, as reported previously.^{6d} This phase transition was observed when doped samples were heated up to 200 °C. Interestingly, a peak-to-peak linewidth of 6 G is characteristic of oxidized TTF-based organic conductors.¹³ Therefore, four-probe DC resistance measurements on doped CuCl₂ xerogels were carried out to study the electrical properties of the two phases and, furthermore, to control the change in the electrical properties during the phase transition. In a typical experiment, a freshly doped xerogel sample was introduced into an oven that was previously set at 200 °C while a four-probe DC measurement of the sample resistance was conducted at the same time. Clearly, we observed typical semiconductor behaviour in the doped xerogel sample, as the resistance decreased with increasing temperature (Fig. 3).^{6d,14} However, around 150 °C, a change in the slope of the resistance *versus* temperature curve was observed, hence underpinning the phase transition studied using EPR spectroscopy. Also, the low angle X-ray powder diffraction patterns of undoped and doped Cu xerogels as a function of temperature

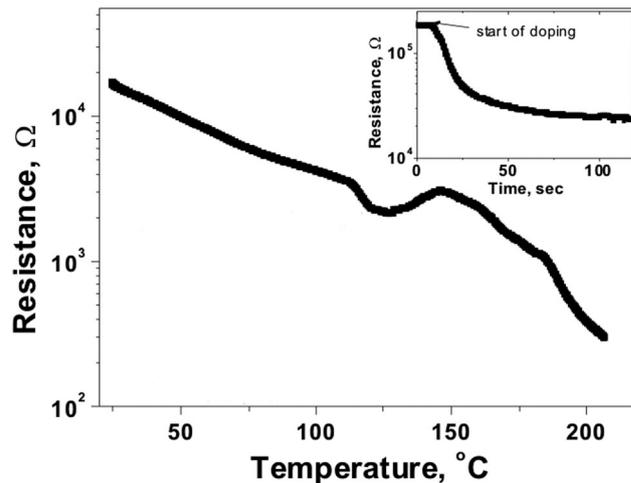


Fig. 3 (a) Temperature dependence of the resistance for a doped CuCl₂ xerogel of **1** during annealing (main panel). Top inset shows typical dependence of resistance vs. time during doping by iodine vapor at room temperature.

(Fig. S14, ESI[†]) clearly indicate different phase transitions between 100 °C and 175 °C. Even though the room temperature resistivity is practically the same for the two phases (around $1.8 \pm 0.2 \Omega \text{ cm}$), their topography changed considerably indicating the irreversible nature of the phase transition (Fig. S15, ESI[†]). While smooth wire-like morphologies were observed for freshly doped xerogel samples, needle-like structures appeared after an annealing cycle was conducted. This change in morphology clearly emphasizes the effect of temperature on the metallo-supramolecular organization.

The electrical properties of the doped film of the CuCl₂ xerogel of **1** were also studied by current sensing atomic force microscopy (CS-AFM). Typical images show regions of high conductivity that have an acicular shape (Fig. 4). These features do not necessarily coincide with the points of greatest height in the film of the xerogel, a feature that rules out any topographic anomaly. In general, scaly features can be appreciated in the topography and the conducting regions are approximately 300 nm wide and up to a micron long, although it is clear that the bulk material is a conductor, otherwise no imaging would be possible. (The contact is made through the graphite substrate.) The *I/V* sweeps at different locations show the clear presence of a gap, characteristic of a semiconducting material.

In summary, we have investigated gelation behaviour of a simple BTDF-fused TTF primarily through coordination bonds of MX₂ (M(II) = Cu, Cd, Co; X = Cl⁻ and Br⁻) with the BTDF moiety. Among the tested transition metal ions, Cu(II) is most likely to undergo self-sustaining gelation. The morphologies of the metallo-gels depend upon the solvent used to dissolve **1**. Furthermore, EPR and four-probe DC resistance measurements confirm that these doped CuCl₂ xerogels can undergo a high-temperature phase transition to a new phase with similar electronic properties to the as-prepared conducting xerogels. This is the first work reporting TTF-based metallo-gels and opens up an avenue for the construction of new low-bandgap materials with various functional properties and novel nanostructures. It is



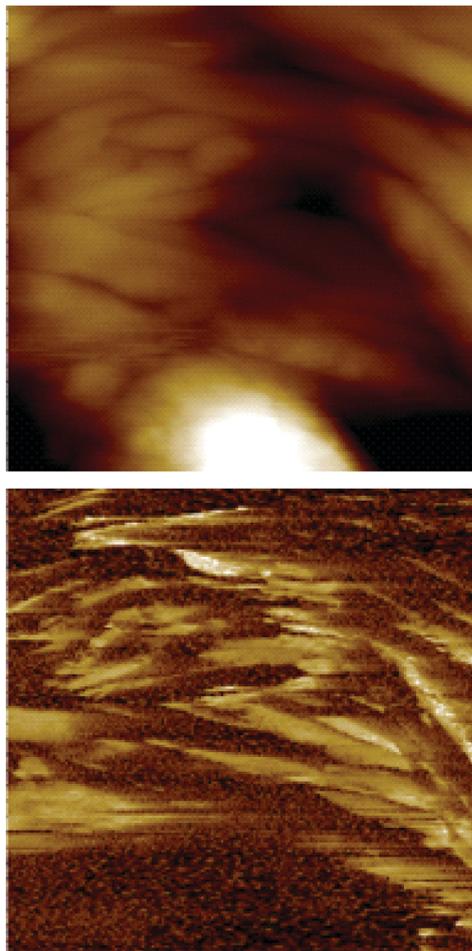


Fig. 4 Topography (top) and current (bottom, bias voltage 1.1 V) images from CS-AFM of the same region (3×3 microns) of a CuCl_2 xerogel of **1** on highly oriented pyrolytic graphite. The brighter regions in the current map correspond to higher conductivity. The independence of the responses is highlighted by the insulating bump in the topographic image.

envisaged that such intensely coloured metallopolymers represent a promising class of materials for photovoltaic applications.

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