

**TITLE**

Motion-Insensitive Determination of  $B_1^+$  Amplitudes Based on the Bloch-Siegert Shift in Single Voxels of Moving Organs Including the Human Heart

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**RUNNING TITLE**

$B_1^+$  Amplitudes Based on Bloch-Siegert Shift in Moving Organs Including the Human Heart

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**WORD COUNT:** 4129 words of body text

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Supported by the Swiss National Science Foundation #31003A-132935 and #310030-149779 to CB)

## ABSTRACT

**Purpose:** To reliably determine the amplitude of the transmit radiofrequency ( $B_1^+$ ) field in moving organs like the liver and heart, where most current techniques are usually not feasible.

**Methods:**  $B_1^+$  field measurement based on the Bloch-Siegert (BS) shift induced by a pair of Fermi pulses in a double-triggered modified Point RESolved Spectroscopy (PRESS) sequence with motion-compensated crusher gradients has been developed. Performance of the sequence was tested in moving phantoms and in muscle, liver, and heart of 6 healthy volunteers each, using different arrangements of transmit/receive coils.

**Results:**  $B_1^+$  determination in a moving phantom was almost independent of type and amplitude of the motion and agreed well with theory. In vivo, repeated measurements led to very small coefficients of variance if the amplitude of the Fermi pulse was chosen above an appropriate level (CV in muscle 0.6 %, liver 1.6 %, heart 2.3 % with moderate amplitude of the Fermi pulses and 1.2 % with stronger Fermi pulses).

**Conclusion:** The proposed sequence shows a very robust determination of  $B_1^+$  in a single voxel even under challenging conditions (transmission with a surface coil or measurements in the heart without breath-hold).

## KEYWORDS

$B_1^+$  determination

Radiofrequency fields

Bloch-Siegert shift

Moving organs

Single voxel

PRESS

## INTRODUCTION

In MRI and MRS, determining the amplitude of the  $B_1^+$  field is crucial to adjust the transmitter gain which leads to the desired rotation of the magnetization by the applied radiofrequency (RF) pulses. While the adjustment of the mostly homogeneous RF field of a body coil is standard in commercially available scanners, the problem of  $B_1^+$  field determination produced by surface coils and particularly in moving organs is far from being solved. Generally speaking, magnetic resonance based  $B_1$  mapping methods can be divided into two groups as magnitude-based and phase-based methods (1,2).

The so-called Double Angle Methods (DAM) use the ratio of two images obtained with two different flip angles to calculate the  $B_1^+$  field (3). In order to avoid the  $T_1$  dependence, it is essential to use long repetition times (TR), which results in increased acquisition times. For that reason, time-efficient versions of DAM have been developed which include sequences that reset magnetization at the end of each acquisition, known as Saturated Double Angle Method (SDAM), or sequences that are combinations of DAM with fast imaging methods (4-6). Actual Flip Angle Imaging uses two identical RF pulses that are followed by different delays. While the method is quite insensitive to  $T_1$  (7) and results in accurate  $B_1^+$  field determinations, this method has some disadvantages which include sensitivity to flow, low signal-to-noise ratios (SNR) due to quick excitations, and possible inaccurate results because of insufficient spoiling of transverse magnetization (7-9). To overcome some of these limitations, derivatives of the method have been developed which include methods with better gradient and RF spoiling schemes (9) and methods that do not use the short TR approximation (8,10). Besides these, there are methods based on multi-pulse sequences like the three-pulse sequence which uses the ratio of stimulated echo and spin echo images (11). Along with other methods (12-18), magnitude-based methods comprise a larger part of  $B_1$  mapping methods.

Phase-based methods cover those employing a pair of composite RF pulses (19), those using non-selective  $180^\circ_x$  and  $90^\circ_y$  pulses which act as  $2\alpha$  and  $\alpha$  pulses (20), and those which take advantage of adiabatic pulses (21-23). Their disadvantages include sensitivity to main field ( $B_0$ ) inhomogeneities (7), requirement for long repetition times, the need for a  $B_0$  map for off-resonance effects (20), and limitations due to specific absorption rate (SAR), respectively. A robust phase-based method for  $B_1$  mapping uses the Bloch-Siegert (BS) shift phenomenon, which is the shift in the resonance frequency of nuclei when an off-resonance RF pulse is applied following excitation (1,24). This method was shown to be insensitive to TR,  $T_1$

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3 relaxation effects, chemical shift,  $B_0$  inhomogeneities, and magnetization transfer (1,2). For  
4 spectroscopy applications, voxel-based  $B_1^+$  field calculations based on a PRESS sequence can  
5 be designed such that they receive data exactly from the same volume as the subsequent  
6 signal acquisition (25).  
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10 Most of the methods are not feasible for moving organs and to the best of our knowledge, the  
11 only successful attempts of  $B_1^+$  mapping in the heart were performed using a magnitude-based  
12 SDAM with spiral readouts (5,26,27), providing whole heart coverage. The examinations  
13 were performed in breath-hold during 16-20 heartbeats resulting in acquisitions as long as  
14 25 s, which makes the method inherently susceptible to irregular heartbeats and poses  
15 additional problems in patients who have difficulties to hold their breath (5).  
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21 In this work, we aimed to develop and test a method that would provide a single-voxel  $B_1^+$   
22 amplitude measurement in the heart within a single heartbeat, such that the method could be  
23 combined with navigator triggering for the free breathing patient and ECG triggering for the  
24 heart. BS shift-based  $B_1^+$  sensitization introduced in a PRESS sequence with motion-  
25 compensated crusher gradients offered the best approach to generate such a sequence. Even if  
26 an additional acquisition without BS effect is necessary to measure the reference phase, the  
27 two required heartbeats can be separated by an arbitrary time period and double-triggering  
28 without breath-hold is possible. Since some spectroscopy applications require the use of a  
29 transmit surface coil, we wanted to test the suggested sequence with different coil  
30 arrangements, including body, head, or surface coil transmission combined with either head or  
31 surface coil reception. The determination of  $B_1^+$  amplitude during surface coil transmission is  
32 particularly important and also demanding since the field falls off rapidly with increasing  
33 distance from the coil. An additional challenge was the fact that no "gold standard" for the  $B_1^+$   
34 determination in the heart is available, which could be used for a comparison with the newly  
35 developed sequence. Therefore, the study was designed such that the methodological  
36 complexity was gradually increased (no motion to various types of motion; homogeneous  
37 transmission with a volume coil to surface coil excitation; experiments in a well defined  
38 phantom to measurements in muscle, liver, and finally heart in vivo).  
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## METHODS

Figure 1 illustrates the modified Point RESolved Spectroscopy (PRESS) sequence with motion-compensated crusher gradients (28,29) and two 8-ms-long Fermi pulses. This PRESS sequence can be double-triggered (navigator-triggered prospective acquisition correction [PACE] and ECG-triggering) and has been implemented on a 3T MR-system (VERIO, SIEMENS, Erlangen, Germany). Fermi pulses were symmetrically applied around the water resonance frequency, selectable at various offset frequencies including  $\pm 2$  kHz,  $\pm 4$  kHz, and  $\pm 8$  kHz. According to the notation of reference (1), the Bloch-Siegert shift  $\Phi_{BS}$ , i.e. the phase difference between two scans with and without Fermi pulse can be calculated as:

$$\phi_{BS} = (B_{1,peak})^2 * K_{BS} \quad [1]$$

$$K_{BS} = \int_0^T \frac{(\gamma B_{1,normalized}(t))^2}{2\omega_{RF}(t)} dt \quad [2]$$

where T is the pulse duration and  $\gamma$  the gyromagnetic ratio.  $K_{BS}$  is a pulse specific parameter which is proportional to the area under the (normalized) pulse envelope and inversely proportional to the off-resonance frequency ( $\omega_{RF}$ ) for the Fermi pulses (1). Both Fermi pulses were programmed to match their  $K_{BS}$  value to 74.01 [rad/G<sup>2</sup>] for the  $\pm 4$  kHz offset case, which was the value used in the original BS shift-based  $B_1^+$  mapping method (1). For a pulse with constant offset frequency as used here,  $\omega_{RF}$  can be extracted from  $K_{BS}$ , leading to:

$$\phi_{BS} = (B_{1,peak})^2 * \frac{1}{\omega_{RF}} * k_{Fermi} \quad [3]$$

$$k_{Fermi} = \int_0^T \frac{(\gamma B_{1,normalized}(t))^2}{2} dt \quad [4]$$

$$B_{1,peak} = \sqrt{\frac{\phi_{BS} * \omega_{RF}}{k_{Fermi}}} \quad [5]$$

with  $k_{Fermi} = 29.60$  [rad\*Hz/ $\mu$ T<sup>2</sup>]. In the current implementation, the Fermi pulses are applied twice such that  $k_2 * k_{Fermi} = 59.20$  [rad\*Hz/ $\mu$ T<sup>2</sup>].

Crusher gradients with two lobes in the standard PRESS sequence were converted into motion-compensated crusher gradients having four lobes such that the net zeroth, first, and second moments would be equal to zero. For each two-lobe-part, the net area obviously needs to be non-zero to act as a crusher for unwanted excitations. The latter two-lobe-part was

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3 implemented so that it would be applied after the Fermi pulse to remove also the unwanted  
4 signals that could be caused by the Fermi pulse.  
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7 Following validation by in vitro experiments using a spherical phantom at rest and with  
8 motion, the sequence was applied on skeletal muscle and liver in one group of 6 healthy  
9 volunteers (3 males, 3 females, age  $24.7 \pm 2.5y$ , BMI  $23.2 \pm 3.0 \text{ kg/m}^2$ ) and in the heart of  
10 another group of 6 healthy subjects (5 males, 1 female, age  $30.2 \pm 3.8y$ , BMI  
11  $24.4 \pm 1.8 \text{ kg/m}^2$ ). Informed consent was obtained from all volunteers and the study was  
12 approved by the Institutional Ethical Committee. Different types of coils were used to  
13 evaluate possible coil specific problems. In vitro experiments were performed using a  
14 transmit/receive (Tx/Rx) volume coil ( $^1\text{H}/^{31}\text{P}$  head coil, RAPID Biomedical, Rimpar,  
15 Germany) and a Tx/Rx surface coil ( $^1\text{H}/^{31}\text{P}$ , RAPID Biomedical, Rimpar, Germany). While  
16 skeletal muscle and liver were measured with the same surface coil used for the phantom  
17 experiments, body coil transmission was combined with reception using spine matrix coil and  
18 body array coil for the heart measurements.  
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21 A phantom (diameter 17 cm, "braino phantom", General Electric, Milwaukee) was measured  
22 at rest and during various types of motion with different speeds and amplitudes generated by a  
23 home-built device located inside the scanner room. The displacement in the positive direction  
24 was controlled by a servo-motor while the phantom was pulled back to its initial position by a  
25 spring force. Two different speeds and amplitudes were used to test motion dependence and  
26 were adjusted to resemble the breathing motion determined by a navigator sequence as much  
27 as possible. Single-shot measurements were taken from a voxel size of  $20 \times 20 \times 20 \text{ mm}$  with an  
28 echo time of 64.2 ms and a repetition time of 1500 ms. Fermi pulses were applied at offset  
29 frequencies of  $\pm 2 \text{ kHz}$ ,  $\pm 4 \text{ kHz}$ , and  $\pm 8 \text{ kHz}$  with varying voltages in order to validate the  
30 expected behavior of the BS shift-based phase change in the phantom using a volume coil.  
31 Surface coil measurements were performed in the phantom by applying the pulses only at  
32  $\pm 2 \text{ kHz}$  offset frequencies as in the in vivo measurements.  
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35 Measurement parameters for the resting skeletal muscle (m.quadriceps femoris) included a  
36 voxel size of  $20 \times 20 \times 20 \text{ mm}$ , an echo time of 64.2 ms, and a repetition time of 1500 ms like the  
37 phantom examinations. The sequence was applied with PACE-triggering in liver due to  
38 respiratory motion and also in breath-hold for comparison. Voxel sizes in the liver varied  
39 between  $20 \times 30 \times 25 \text{ mm}$  and  $40 \times 40 \times 40 \text{ mm}$ . Echo time used for the liver was 64.2 ms while the  
40 repetition time was either 1500 ms in breath-hold or determined by the trigger signal.  
41 Measurements in the heart were taken in PACE/ECG double triggering mode as described in  
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3 Ith et al. (30) to freeze respiratory and cardiac motion with a voxel size of 12x25x20mm  
4 positioned in the cardiac septum and a repetition time defined by the respiration rate. ECG  
5 trigger delay (typically 250 ms) was adjusted such that the sequence started in the late  
6 contraction/beginning relaxation phase. In order to test the variation of the measured BS shift,  
7 16 repetitions of single-shot measurements were acquired in the heart with a scan time of  
8 approximately 2 minutes.  
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13 Except for the heart, DICOM files of single coil receptions were analyzed using MATLAB  
14 (The MathWorks, Inc.). Due to the array coil arrangement used in the heart and thus inherent  
15 phase adjustment between the single coils in the DICOM file, which made a determination of  
16 the BS shift impossible, raw data from the scanner had to be analyzed using a modified  
17 version of the MATLAB program provided by Maolin Qiu at Yale University (SIEMENS  
18 IDEA forum). The  $B_1^+$  amplitude was calculated from the phase difference of two single-shot  
19 acquisitions with and without the application of Fermi pulses according to the relation given  
20 in Eq. [1]. The phase was calculated from the first point in the time domain signal for each  
21 scan, since the first point corresponds to the echo top and contains most of the signal  
22 information.  
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31  $B_1^+$  maps from a work-in-progress  $B_1^+$  mapping package (3P-MAP, based on the stimulated vs.  
32 spin echo responses for a 3-pulse sequence, duration 17 s, SIEMENS, Erlangen, Germany)  
33 were obtained for comparison in all cases except for the heart where the  $B_1^+$  determination  
34 was not feasible because the 3P-MAP sequence did not have PACE/ECG double triggering  
35 options. The scanner provides a reference voltage (REF) that corresponds to a 90 degree  
36 excitation by a rectangular 500  $\mu$ s pulse, i.e., to a  $B_1^+$  field of 11.75  $\mu$ T. The determination of  
37 REF is reliable as long as homogeneous excitation fields (e.g. body coil) are used; however,  
38 with surface coil excitation, it often underestimates the required RF power since the algorithm  
39 uses the high intensity signal from the regions close to the coil.  
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47 To describe the repeatability, coefficients of variance were calculated. Agreement between the  
48 two  $B_1^+$  methods was calculated based on the Lin's Concordance Correlation Coefficient  
49 (MATLAB, The MathWorks, Inc.).  
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## RESULTS

In order to compare the behavior of the BS shift effect in the phantom with theory, Fermi pulses were applied using four different voltages (120, 180, 240, 300 V with the transmitter reference voltage being 174 V) and three different frequency offsets. Fig.2 shows these stationary and moving phantom results for the induced phase change with respect to the reference scan. Black curves, representing the fitted data for the non-moving phantom, are almost perfectly aligned with the measured values of the non-moving phantom. The single BS shift-based  $B_1^+$  measurements of the non-moving phantom (red crosses, 15 single-shot repetitions) are indistinguishable. When the motion was started with very fast speed and large amplitude, the measured phase changes became slightly (average CV of 4.9%) scattered (blue crosses) without systematic under- or overestimation of the theoretical curve for the non-moving phantom. BS shift-based  $B_1^+$  measurements determined the  $B_1^+$  amplitude as  $0.065 \mu\text{T/V}$  (SD  $0.001 \mu\text{T/V}$ ) for the stationary case (black curve) and  $0.067 \mu\text{T/V}$  (SD  $0.002 \mu\text{T/V}$ , curve not shown) for the moving case. 3P-MAP was repeated three times for the stationary case giving a  $B_1^+$  amplitude of  $0.066 \mu\text{T/V}$  (SD  $0.002 \mu\text{T/V}$ ) in the voxel in all repetitions. The scanner adjustment corresponds to the  $B_1^+$  amplitude as  $0.068 \mu\text{T/V}$ ; however, determined over the whole 17 cm-diameter phantom.

A comparison between the BS shift-based  $B_1^+$  measurement and the 3P-MAP in the phantom using the double-tuned surface coil for different types of motion is displayed in Fig.3. Average values and error bars, representing the standard deviations, are shown for 10 single-shot measurements for the BS shift-based  $B_1^+$  measurement and 10 repetitions of the 17-second-long 3P-MAP sequence. It is evident that with increasing speeds and amplitudes, the standard deviation of the 3P-MAP increased. In addition, a systematic overestimation of the  $B_1^+$  field was observed even with slow motion. In contrast, the BS shift-based  $B_1^+$  measurement continued to produce consistent results. For the fast motion with large amplitude, standard deviation of the 10 single-shot measurements for the BS shift-based  $B_1^+$  measurement for the calculated flip angle was 0.0384 radians.

The resting muscle was used to compare the 3P-MAP with the newly developed BS shift-based sequence in an organ without motion but with a very inhomogeneous  $B_1^+$  profile of the transmitting surface coil (explaining the large differences between subjects). Figure 4 shows the comparison between the single-shot measurements of the BS shift-based  $B_1^+$  measurement and the repetitions of the 3P-MAP sequence. Lin's Concordance Correlation Coefficient was



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3 calculated as 0.98 for the two methods with a calculated mean flip angle difference of 0.0227  
4 radians. A test of repeatability of the BS shift-based sequence in 15 single-shot repetitions  
5 each in 6 volunteers resulted in an average coefficient of variation (CV) of 0.6%.  
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9 The next in vivo tests were performed in the liver in order to compare 3P-MAP with BS shift-  
10 based  $B_1^+$  measurements in a slowly moving organ. Since motion during the application of the  
11 gradients leads to phase shifts, BS shift-based  $B_1^+$  measurements were applied both in breath-  
12 hold and with triggering to test the efficiency of the motion-compensation. Figure 5  
13 demonstrates the results for the triggered BS shift-based  $B_1^+$  measurement (16 single-shot  
14 measurements), the breath-hold BS shift-based  $B_1^+$  measurement (7 single-shot  
15 measurements), and the 3P-MAP sequence (3 repetitions of breath-hold measurements). Lin's  
16 Concordance Correlation Coefficient for the triggered BS and breath-hold BS measurements  
17 was 0.92 and the mean  $B_1^+$  amplitude difference was 0.069  $\mu\text{T}$  (SD 0.090  $\mu\text{T}$ ). While the two  
18 BS-based methods agree very well (Bland-Altman plot Fig.5B), the 3P-MAP shows  
19 significantly higher  $B_1^+$  amplitudes than those measured by the BS shift-based method (Bland-  
20 Altman plots in Fig.5C and 5D). The same behavior has been observed in the moving  
21 phantom where the BS-based method was not influenced by the motion (Fig.3). Mean  $B_1^+$   
22 amplitude difference between the triggered BS shift-based measurements and 3P-MAP  
23 measurements was 1.33  $\mu\text{T}$  (SD 0.14  $\mu\text{T}$ ). A test of repeatability of the triggered BS shift-  
24 based sequence using surface coil transmission was performed in the liver of 6 volunteers in  
25 16 single-shot repetitions each and resulted in an average CV of 1.6%.  
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38 The heart with two independent motions of the respiratory and cardiac action represents the  
39 most challenging organ for motion-sensitive sequences. Figure 7 illustrates the changes in the  
40 repetitions of the single-shot measurements when the Fermi pulses were applied at different  
41 voltages relative to the transmitter reference voltage for all volunteers. Similar to the muscle  
42 and liver results, single-shot repetitions were stable with an average CV of 2.3 % for the  $B_1^+$   
43 amplitude change due to the BS shift for 16 single-shot repetitions in 6 volunteers when the  
44 Fermi pulses were applied at the transmitter reference voltage. Above the reference voltage,  
45 the variations were even lower (1.2 %, measured in 5 subjects only) while the results were not  
46 optimal when the Fermi pulse voltages were below the REF voltage (CV of 7.5 %).  $B_1^+$   
47 amplitudes (generated by the body coil) were calculated as 10.52  $\mu\text{T}$ , 9.33  $\mu\text{T}$ , 10.43  $\mu\text{T}$ ,  
48 10.07  $\mu\text{T}$ , 13.55  $\mu\text{T}$ , and 9.85  $\mu\text{T}$  in different subjects. The relative difference between the  
49 subject with the strongest and weakest  $B_1^+$  in the ventricular septum (9.33  $\mu\text{T}$  for a male  
50 1.93 m tall and weight of 100 kg, vs. 13.55  $\mu\text{T}$  for a female 1.57 m tall and weight of 57 kg) is  
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3 37% (difference/average) and would correspond to excitation angles of 62 degrees vs. 90  
4 degrees if the RF field would not be adjusted to the value in the voxel.  
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## DISCUSSION AND CONCLUSIONS

The suggested PACE/ECG-triggered PRESS sequence reveals reliable measurements of  $B_1^+$  with single-shot acquisitions (applied twice - once with BS shift and once without). It is almost immune against motion related perturbations and can also be applied in the heart. The sequence has been validated in vitro and in vivo in various organs and with different coil arrangements.

To begin with, the fitted curves of the measurements in the phantom confirm the theoretically expected behavior with the quadratic phase increase with the Fermi voltage as well as the linear phase increase with the offset frequency (see Fig.4). Measurements during motion are almost identical to the measurements without motion and the variations are still very small, proving the robustness of the sequence against motion in vitro. In addition to the compliance with theory, the results agree well with the 3P-MAP results when the 3P-MAP is applied in the phantom at rest.

The variation among the repetitions of the single-shot measurements in motionless organs like the quadriceps muscle is very small when the BS shift-based sequence is applied. Comparison of the sequences in the muscle and the high concordance correlation coefficient indicate the high level of agreement between the two methods in an immobile organ where the 3P-MAP can be used as a "gold-standard" for the newly developed sequence. Since only the 3P-MAP deteriorates with motion while the BS-based  $B_1^+$  values remain stable, we conclude that the values measured by the BS-shift method are accurate, even if there is no longer a "gold-standard" to compare with.

Single-shot measurements in the heart are highly reproducible with small variances as long as the nominal Fermi voltage is above a threshold value like the transmitter reference voltage determined by the scanner. Lower Fermi pulse voltages resulted in larger noise of the phase shift and thus of the  $B_1^+$  value. This observation is important since the specific absorption rate (SAR) of the Fermi pulse is relatively high and imaging implementations of the BS technique tend to choose smaller Fermi pulse amplitudes; however, for MR spectroscopy sequences with long TR and few acquisitions, this does not apply. Since 3P-MAP does not work in heart due to the lack of double-triggering, agreement between BS shift-based sequence and 3P-MAP cannot be shown.

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3 The original BS shift-based  $B_1^+$  mapping sequences, which were modifications of spin echo  
4 and gradient echo sequences, were applied in the abdomen as a moving organ without breath-  
5 hold (1). Since the sequence was not motion-compensated, artifacts in the phase-encoding  
6 direction were reported even though they were stated not to affect the results significantly (1).  
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8 In the heart, due to motion complexity, this method would be insufficient without motion-  
9 compensated gradients. Besides the successful tests for  $B_1^+$  mapping in the heart using a  
10 magnitude-based SDAM with spiral readouts (5,26,27), we were able to determine the single-  
11 voxel  $B_1^+$  amplitude in the heart for the first time using a phase-based  $B_1^+$  measurement  
12 method without the need of a breath-hold. It will be highly beneficial to single-voxel MRS  
13 measurements in the heart or in any organ especially when surface coils are used since the  
14 adjustment routine from the scanner seemed to be insufficient for the flip angle calibration.  
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16 The 37%  $B_1^+$  value difference between two subjects illustrates the importance of a localized  
17 determination of the RF field strength in the heart. Without adjusting the RF power for the  
18 heart, the flip-angles for the two subjects would differ as 90 degrees to 62 degrees resulting in  
19 a considerable reduction of signal quality, e.g. in a PRESS selected spectrum. Based on an  
20 earlier study on heart lipids without a localized determination of  $B_1^+$  (30), we assume that  
21 insufficient RF adjustment is responsible for a considerable fraction of low quality spectra.  
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23 The direct consequences of incorrect  $B_1^+$  adjustment for spectral quality and accuracy have not  
24 been shown experimentally in this study; however, there are many publications which show  
25 the crucial influence of this adjustment e.g. on the pulse profile and thus for the shape of the  
26 voxel (25,27), signal contamination from neighboring tissues, and also for overall signal  
27 magnitude.

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30 Phase dependent methods have to solve the problem of phase wrapping, in particular in  
31 imaging mode where large differences of regional  $B_1^+$  may occur. In the proposed single voxel  
32 sequence, even if phase wrapping occurred between the acquisitions with vs. without Fermi  
33 pulse, unwrapping was straightforward if the Fermi pulse strength was chosen appropriately,  
34 i.e. approximately at the REF level such that the effect was large enough to generate sufficient  
35 effects, yet small enough to be below  $2\pi$ .

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38 An inherent strength of the Bloch-Siegert effect is its independence of the signal amplitude  
39 which makes it particularly suited for triggered acquisitions where the TR can change  
40 depending on the breathing and/or cardiac cycle. Under these conditions, amplitude-based  
41 methods need special corrections; however, Bloch-Siegert shift based methods are immune  
42 against variations in amplitude and hence TR.

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3 Another strength of phase-based methods is the fact that the first point in the FID (or central  
4 part of the echo) contains the full phase information; i.e. the sequence is inherently short  
5 which is beneficial in moving organs since it is not necessary to acquire the full acquisition  
6 window.  
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10 The major limitation of the current method is its single-voxel nature, giving accurate values  
11 for a specific location, yet not for the complete volume. For single-voxel MR spectroscopy,  
12 this limitation is not a problem; however, if data acquisition on a whole volume is required,  
13 for instance to correct for absolute quantities of a CSI experiment or a FID acquisition, this  
14 limitation would be severe. 2D-Spiral readout methods could be combined with the sequence;  
15 however, at the price of a longer acquisition window.  
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19 Because it was aimed as a proof of principle, we did not try to shorten the sequence duration  
20 for the experiments further. Nevertheless, removing the second Fermi pulse from the  
21 sequence and instead of the reference scan without the BS shift, applying the only Fermi pulse  
22 with the negative frequency offset is a first step towards shortening the sequence. Fermi  
23 pulses were applied with positive and negative offsets in the same scan in order to avoid  
24 problems that might occur between different scans. Combining the crusher lobes with the  
25 refocusing lobe of the slice-selection gradient and combining all crusher gradients with fewer  
26 lobes after recalculating the moments further shorten the sequence. By performing these  
27 modifications, it was already possible to reduce the echo time from 64.2 ms to 42.6 ms. An  
28 additional potential for a shorter sequence is reducing the duration of the Fermi pulses which  
29 are relatively long.  
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33 Another limitation of the sequence is the fact that the slice selection gradients and the  
34 refocusing gradients were only balanced for zeroth moments. However, these gradients were  
35 considerably smaller than the crusher gradients and their possible effects on the phase due to  
36 motion would be minimal.  
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40 Using surface coils for transmission poses specific problems due to the large range of  $B_1^+$   
41 strengths and the possibility of hotspots close to the coil. Nonetheless, some spectroscopy  
42 applications (in particular the combination with non-proton nuclei) require surface coils where  
43 also the proton-part is integrated (e.g. for decoupling). The determination of  $B_1^+$  is particularly  
44 crucial for such an arrangement; however, a single voxel sequence can only provide the field  
45 strength in one part of the tissue.  
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3 SAR issues due to the long Fermi pulses should not be a major problem thanks to the small  
4 number of only 2 scans in a final implementation and since the acquisition and the reference  
5 scan can be separated by an arbitrarily long period. The much larger number of scans in the  
6 current study served only to test the repeatability and will not be necessary in a final  
7 implementation where two scans are sufficient.  
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11 In conclusion, results suggest that the new BS shift-based PACE/ECG-triggered modified  
12 PRESS sequence is a robust and fast method for  $B_1^+$  amplitude determination in a single voxel  
13 within moving organs.  
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## ACKNOWLEDGEMENTS

The authors thank SIEMENS (Erlangen, Germany) for providing the work-in-progress packages for  $B_1^+$  mapping and navigator-triggered prospective acquisition correction (PACE); and Karin Zwygart for organizing volunteers and data storage.

For Peer Review



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## FIGURE CAPTIONS

Figure 1:

Diagram of the proposed sequence with motion-compensated crusher gradients (cyan) and 8 ms long Fermi pulses (yellow) following the 180°-pulses (red). Fermi pulses are applied at symmetric offset frequencies (+ and -) with respect to the water resonance frequency. The sequence can be preceded by a navigator-triggered prospective acquisition correction trigger (PACE, not shown) and can also be ECG-triggered.

Figure 2:

Phase changes as a function of Fermi pulse voltage and frequency offset for a phantom without motion (blue crosses, indistinguishable for the 15 repetitions each) and with motion (red crosses showing 15 repetitions each) in a volume transmit/receive coil. The black line represents the fit of the measurements without motion with the theoretical behavior according to Eq. [3] (scaling of the curve with  $0.1245 \text{ Hz} \cdot \text{rad}/\text{V}^2$ ), resulting in a  $B_1^+$  amplitude of  $0.0648 \mu\text{T}/\text{V}$ . The fit of the measurements with motion resulted in a  $B_1^+$  amplitude of  $0.0673 \mu\text{T}/\text{V}$ . As a comparison, the 3P-MAP method repeated three times for the stationary case gave a  $B_1^+$  amplitude of  $0.0660 \mu\text{T}/\text{V}$  for the selected voxel, while the scanner adjustment for the whole object determined a value of  $0.0675 \mu\text{T}/\text{V}$ .

Figure 3:

Influence of motion (speed and amplitude) on the  $B_1^+$  field determination by the BS shift-based  $B_1^+$  acquisition (blue) and 3P-MAP sequence (red), using surface coil transmission and reception in vitro (10 repetitions each). Without motion, the two methods agree very well while increasing motion (in amplitude and speed) deteriorates the results of the 3P-MAP sequence resulting in large variations and systematically overestimated values. The BS shift-based acquisition shows only a small increase of the variation between the 10 repetitions without a systematic bias. Error bars represent the standard deviation.

Figure 4:

Comparison of  $B_1^+$  field amplitudes determined by the BS shift-based  $B_1^+$  sequence and 3P-MAP in human quadriceps muscle. Different locations for the voxels in the inhomogeneous  $B_1^+$  field of a surface coil transmission mainly explain the variations between subjects. In addition, different coil loading could add to the variations since a fixed transmitter voltage was used for all subjects.

Figure 5

5A) Comparison of the  $B_1^+$  measurements in liver with 3P-MAP (black crosses, 3 repetitions), triggered BS shift-based measurements (red crosses, 16 repetitions) and BS shift-based results in breath-hold (blue crosses, 7 repetitions) for 6 subjects. Differences between the subjects are mainly a consequence of the voxel position. A comparison of the three methods is shown by Bland-Altman plots (identical vertical scale) between (B) triggered BS shift-based measurements vs. BS shift-based results in breath-hold, (C) triggered BS shift-based

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3 measurements vs. 3P-MAP; and (D) BS shift-based results in breath-hold vs. 3P-MAP. In  
4 order to obtain the same number of examinations for both methods as required in the Bland-  
5 Altman plots, only the first 7 repetitions shown in (A) were used in (B) and 3 in (C) and (D),  
6 respectively. The comparison of the two BS-based examinations in (B) shows very good  
7 agreement (zero included in the error range which is small) while the comparisons of BS-  
8 based results with 3P-MAP in (C) and (D) show a significant difference (zero outside the  
9 error range of 1.96 standard deviations).  
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13 Figure 6

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15 Radiofrequency field strength ( $B_1^+$ ) for the 16 repetitions (each with vs. without Fermi pulses)  
16 in the heart of 6 volunteers (body coil transmission, surface coil detection) with various  
17 amplitudes of the Fermi pulse (relative to "REF" = voltage which is determined by the  
18 vendor's adjustment routine). The curves in grey were acquired with Fermi pulses below the  
19 REF voltage and show relatively large variations while the curves in red and blue were  
20 measured at the REF voltage and above, respectively, and show very high reproducibility  
21 within the 16 repetitions.  
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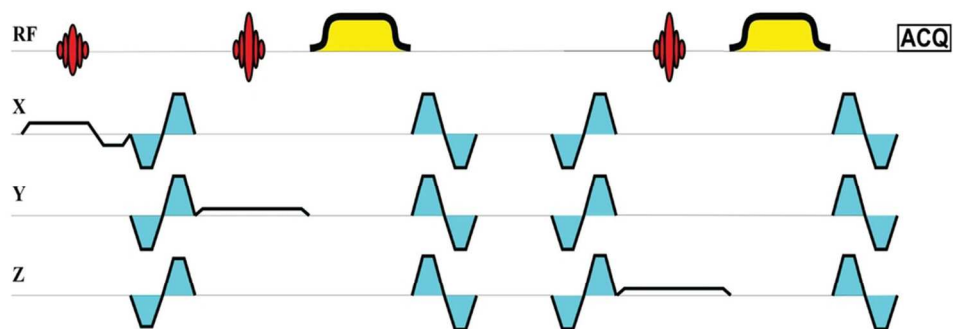


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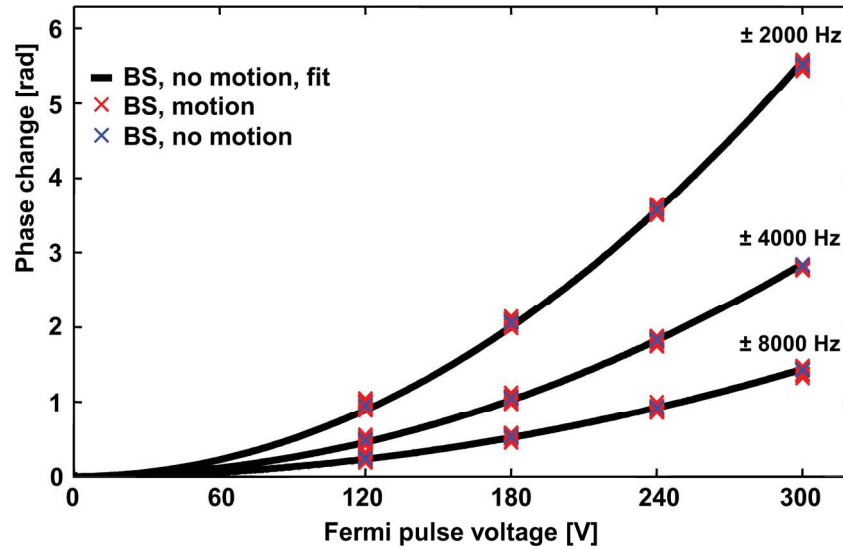


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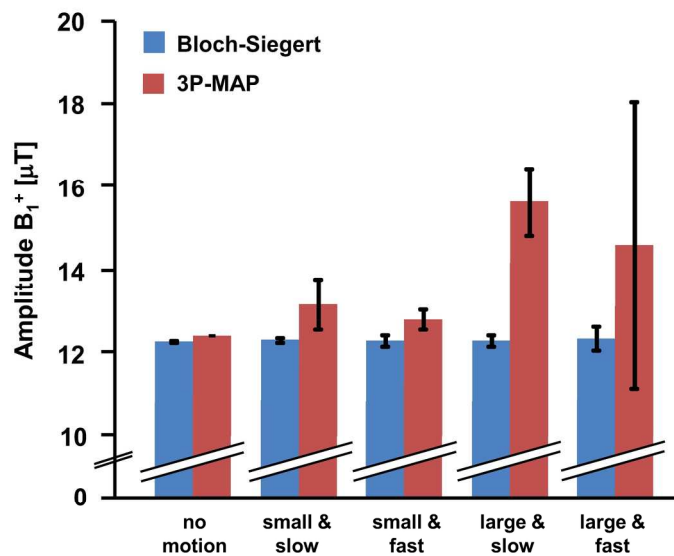


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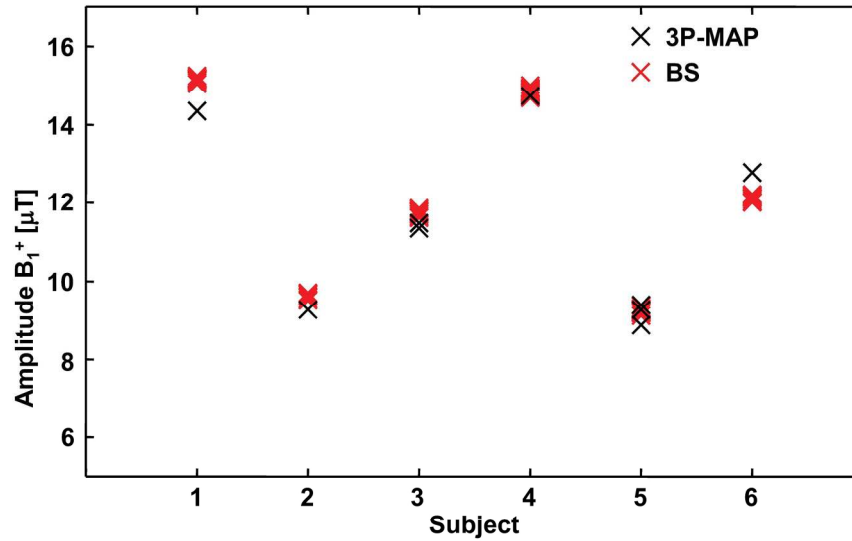


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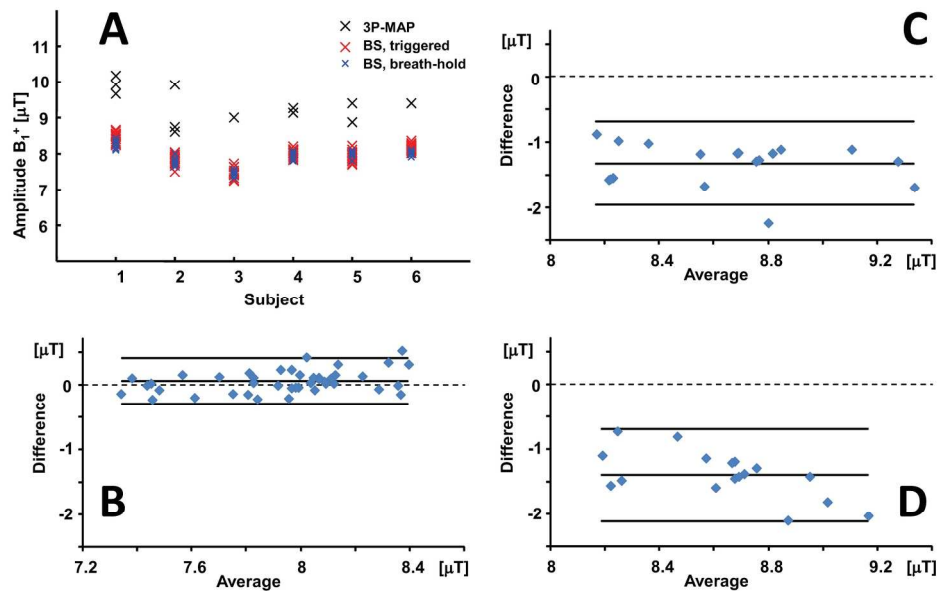


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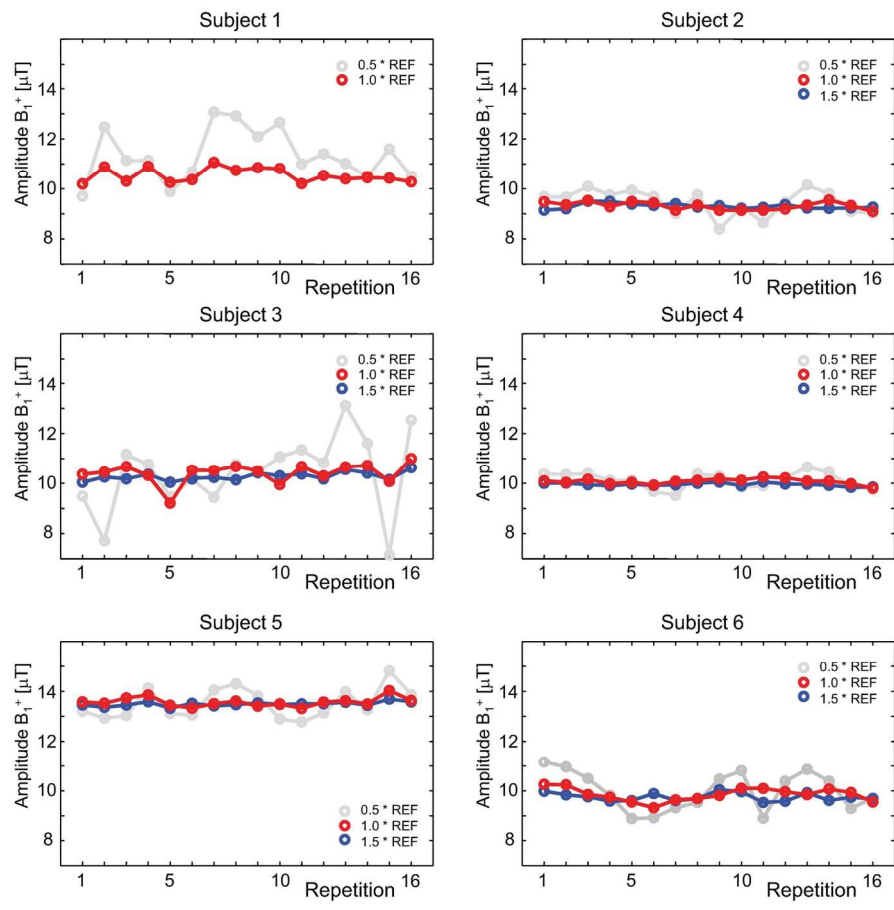


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