A mononuclear Fe(II) complex involving a tetrathiafulvalene-based ligand exhibits thermal spin-crossover (around 143 K) with pronounced hysteresis behaviour (48 K). The chromophoric and π-extended ligand allows Near-Infrared (NIR) sensitization for the light-induced excited spin-state trapping (LIESST) with $T_{\text{LIESST}} = 90$ K.

Molecular magnetism is a very active field where both chemist and physicist communities are looking for systems with high-density memory capabilities. In this context, Single Molecule Magnets (SMM) 3 and Spin Crossover (SCO) materials 2 are intensively studied because of their potential applications. 3 SCO behaviour has been known since the 1930s. 4 In this respect, the Fe(II) ion is the most popular magnetic centre owing to its large variation of the spin value ($S = 2 ightarrow S = 0$), the largest possible for 3d metal centres along the high-spin (HS) to low-spin (LS) transition. A crucial parameter for device applications is the presence of magnetic bistability that is driven mainly by the cooperativity induced by intermolecular interactions in the crystal lattice. Thus, the switching of spin states can be induced by changing the size of the coordinated ligands, counter ions and/or the molecules of crystallization. 5 The magnetic bistability can also be reached under kinetic conditions by light irradiation at low temperatures (light-induced excited spin state trapping, LIESST). 6 All these aspects have already been studied on several derivatives of [Fe(H$_2$Bpz$_2$)$_2$(X)] complexes (H$_2$Bpz$_2$ = dihydrobis(1-pyrazolyl)borate, X = 2,2′-bipyridine, 1,10-phenanthroline, dipyrdo-[3,2-a:2′,3′-c]phenazine (dppz)). 5,7 Moreover, the SCO behaviour of such derivatives has recently been observed upon deposition on a gold surface and in thin films. 8 Herein, we propose to extend the π-system of the dppz ligand through fusion with a tetrathiafulvalene (TTF) derivative in order (i) to combine redox activity with SCO behaviour, 9 (ii) to increase the cooperativity via enhancement of π-stacking interactions, thus the hysteretic behaviour around the transition temperature, 10 and (ii) to induce strong intra-ligand charge transfer (ILCT) transitions reaching the NIR spectral region for LIESST investigation. Sensitization of the LIESST effect through ILCT excitations could present some advantages such as (i) the ILCT bands allow lower-energy excitations than the MLCT and (ii) they can be easily modulated in energy through chemical changes of the push-pull ligands and (iii) the absorption intensity is higher than forbidden MLCT transitions.

In this article, we report the first TTF-based complex of Fe(II) displaying both thermal and photoinduced spin crossover. One should mention here that Oshio et al. 11 previously proposed an iron(II) complex containing a TTF moiety that shows thermal spin-crossover but without a photomagnetism study. 11 [Fe(H$_2$Bpz$_2$)$_2$(L)]($\text{CH}_2\text{Cl}_2$)$_2$, called 1($\text{CH}_2\text{Cl}_2$)$_2$, was elaborated from the in situ formation of the Fe(H$_2$Bpz$_2$)$_2$ precursor followed by coordination with the TTF-fused dppz ligand (L). 12 The photophysical properties of the related [Fe(phen)$_3$(L)]$^+$ complex have been reported. 13 Compound 1($\text{CH}_2\text{Cl}_2$)$_2$ crystallises in the triclinic space group P1 (no. 2). 6 The ORTEP drawing of 1 is shown in Fig. 1.

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**Note:** The ORTEP view of the structure of 1($\text{CH}_2\text{Cl}_2$)$_2$. Atomic displacement ellipsoids are drawn at the 50% probability level. Hydrogen atoms are omitted for clarity.
TTF donors. Upon coordination to Fe(II), the dppz-centred dication potentials are comparable to a range of functionalized TTF derivatives, as well as those at higher energies were in good agreement with the calculated ones (Table S2†) and best fit (full black line) \( R = 0.9991 \). Colour code: red, ILCT; blue, LLCT (Ligand to Ligand Charge Transfer); purple IL and black IA. Representation of the molecular orbitals involved in the low-energy ILCT.

The cyclic voltammogram of 1 shows three single-electron redox waves at \(-1.11 \) V, 0.71 V and 1.06 V corresponding to the formation of the dppz radical anion, the radical cation and the dication of the TTF fragment, respectively (Fig. S1†). These oxidation potentials are comparable to a range of functionalized TTF donors. Upon coordination to Fe(u), the dppz-centred reduction wave is anodically shifted from \(-1.23 \) V (ref. 12) for the free L to \(-1.11 \) V for 1 while the TTF-centred oxidation waves remain unchanged. These electrochemical properties look promising for controlling magnetic properties by changing oxidation states.

The UV-visible absorption of 1 was studied in CH2Cl2 solution (Fig. 3) and in the solid state (Fig. S2†). TD-DFT calculations were performed to rationalize the absorption spectrum (Fig. S3, S4 and Table S2†), however considering the Fe(u) ion in its LS (\( S = 0 \)) state (see the ESI† for computational details). The lowest energy bands were attributed to the ILCT transition process are so severe that crystals crack and do not recrystallize when heated up to 120 °C (Fig. S5†). After the initial loss of the wetting solvent of a fresh sample, the observed mass loss (14%) conforms to a full desolvation process (theo, 15%). After desolvation, the magnetic behaviour of 1 is radically different since it shows a HS to LS transition centred at 143 K with a wide thermal hysteresis (48 K) which can be cycled (Fig. S6†). The transition occurs at \( T_\text{C1} = 119 \) K in the cooling mode and \( T_\text{C1} = 167 \) K in the warming mode. Both curves have been recorded at a rate of 0.3 K min\(^{-1}\). The transition occurs in the same temperature range as in the parent \( \text{H}_2\text{Bpz}_2 \) ligand based compounds, but with larger thermal hysteresis. It is expected that the loss of solvent from the crystal structure increases the overlap between TTF moieties (Fig. 2), leading to the extension of the interaction between the iron(u) centres. Unfortunately, the structural modifications upon the desolvation process are so severe that crystals crack and do not diffract anymore.

Starting from the desolvated form (1), irradiation at 780 nm (12 820 cm\(^{-1}\), 5 mW) for 5 min.

Fig. 3 UV-visible absorption spectrum of 1 in CH2Cl2 solution (c = 2 \times 10^{-5} \text{ mol L}^{-1}) with respective Gaussian deconvolutions (dashed lines) and best fit (full black line) \( R = 0.9991 \). Colour code: red, ILCT; blue, LLCT (Ligand to Ligand Charge Transfer); purple IL and black IA. Representation of the molecular orbitals involved in the low-energy ILCT.

Fig. 2 Crystal packing of 1 (HS) in 1/CH2Cl22 highlighting the “head-to-tail” dimers.
90 K at a sweep rate of 0.3 K min⁻¹. This value is much higher than the one (44 K) reported for [Fe(H₂Bpz₂)₂phen].²⁶ Irradiation at 530 nm (18 870 cm⁻¹) and 660 nm (15 150 cm⁻¹) produces the same quantitative LIESST effect while the irradiation at 404 nm (24 750 cm⁻¹), outside the HOMO → LUMO band (Fig. S2†), is less efficient. Excitation in the low-energy singlet ILCT band (780 nm) might induce fast relaxation to the triplet T₁ level of Fe²⁺ and progressive population of the quintet T₃ state of the HS Fe²⁺ ion. Even if we are aware that the sensitization process is controversial,¹⁷ a hypothetical proposed mechanism for the LIESST effect is given in Fig. S7.†

The spectacular difference in magnetic behaviour of 1·(CH₂Cl₂)₂ and 1 calls for further investigation. Electronic structure calculations have been conducted on the periodic crystal structure using the density functional theory (DFT) framework as implemented in the VASP code.¹⁸ The Perdew–Burke–Ernzerhof form generalized gradient approximation was used for the exchange and correlation functional.¹⁹ Wave functions were expanded using a plane-wave basis set with a cutoff energy of 400 eV. Core electrons were treated within the projector augmented wave method.²⁰ The simulation cell contains two molecules of 1 and, in the case of the solvated structure, three molecules of CH₂Cl₂ and two of H₂O. The latter structure was chosen to account for (i) the presence of two CH₂Cl₂ molecules and (ii) the disorder created at the two half-occupied sites in the crystal structure of 1·(CH₂Cl₂)₂, that are mimicked by two H₂O molecules. Given the size of the cell and the nature of the system, Γ-point calculations turn out to give satisfactory accuracy. Geometries were optimized until the forces were smaller than 0.01 eV Å⁻¹. Forces were corrected for the missing van der Waals interactions using the Grimme scheme.²¹ It is known that usual functionals fail to take into account the strongly correlated character of iron d-electrons. Therefore, a Hubbard-like term is added to the Kohn–Sham Hamiltonian,²² controlled by an effective term Ueff.²³ Several propositions have been formulated regarding the appropriate value for Ueff in the case of SCO compounds ranging from 1.55 to 2.65 eV.²⁴ As we did not intend to reach a quantitative agreement, but rather a qualitative description of the physical processes, we used a value Ueff = 2.00 eV. We checked that a value of 3.00 eV was not affecting dramatically the geometry of the systems (variations are less than 0.5%). For both solvated and non-solvated compounds, the HS (S = 2) and LS (S = 0) solutions have been considered. In the case of the HS state with solvents, the resulting optimized cell is in good agreement with the experimental structure (Table S3†). In particular, the optimized HS molecule shows deviations of less than 3% in bond lengths (Table S4†).

In order to achieve SCO, a system has to possess specific features. In particular, the enthalpy should favour the LS state since the entropic contributions (electronic and vibrational) will benefit to the HS state. This is characterized by the adiabatic gap ΔEadia = EHS − ELS, i.e. the energy difference between the HS and LS states each considered in their own geometry. As a necessary, but not sufficient condition, ΔEadia should be positive. In the case of solvated 1 (Fig. 5a), ΔEadia = −0.40 eV, which leads to an HS system whatever the temperature, in agreement with the experimental observation. The situation is different with 1. Indeed, in the absence of solvents, the LS state is stabilized and ΔEadia reaches 0.41 eV and SCO behaviour becomes possible.

When considering the optimized geometries of the HS and LS solutions with and without solvents (Fig. S8 and Table S4†), it becomes clear that the variation of the adiabatic gap can be traced back to a competition between Fe-coordination and π–π interactions. In particular, when considering the HS solution with and without solvents (Fig. 5b, c and Table S4†), one can notice that the H₂Bpz₂ ligand opts either for a Fe-coordination which leads to an expected Fe–N bond of 2.15 Å, or without solvent, an interaction with the π-system of L which leads to a much longer bond (2.42 Å). The distortion of the coordination sphere leads to a destabilization of the HS solution and makes SCO possible.

As a final remark, let us note that in the absence of solvent, the Fe–Fe distance goes from 13.48 Å to 13.42 Å when going from HS to LS. It is known that a limited contraction in the Fe–Fe distance along the HS to LS transition favours the electrostatic contribution to cooperativity hence to the opening of a hysteresis loop.²⁵ In addition to the existing π-network, this strongly supports the potential of TTF-based compounds as candidates for SCO systems with large hysteresis.

![Fig. 5 (a) Adiabatic gap ΔEadia = EHS − ELS for 1 with and without solvent molecules (CH₂Cl₂ and H₂O), calculated at the PBE + D3 + U level (Ueff = 2.00 eV). (b), (c) Optimized geometry of the HS state with and without solvent, respectively. Fe, Cl, S, O, N, C, B and H atoms are depicted in orange, light green, yellow, red, blue, brown, dark green, and pink, respectively. Without solvent molecules, the HS geometry shows a competition between Fe-coordination and the π–π interaction between H₂Bpz₂ and the ligand L.](image-url)
To summarize, compound 1 was obtained from the association of the TTF-fused dipyrido[3,2-a:2',3']-phenazine ligand (L) and the metallo-precurser Fe(H2Bpz2)2. This neutral heteroleptic complex presents both thermal and light-induced spin transitions. The π-extended chromophoric ligand L exhibits an intense and energetically low-lying ILCT absorption band, which enables sensitization of LIESST in the NIR spectral region. The analogous NIR antenna effect of this ligand has recently been described for the NIR emission of Yb(III) doped silica nanoparticles.26 The large π-system of L increases as well the interaction between the magnetic centres which leads to one of the largest thermal hysteresis for a mononuclear spin transition complex of Fe(n). The use of a TTF-based ligand provides a promising possibility to have high temperature LIESST effects and large thermal hysteresis loops as well as to use the redox activity of the TTF ligand to switch the magnetic properties.

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Notes and references

\[ \text{[Fe(H\_Bpz2)\_2(L)\_2\_CH\_2\_Cl\_2]} = 16 \text{ mg of FeCl\_2\_4H\_2O (0.04 mmol)} \]

\[ \text{Crystals data for [Fe(H\_Bpz2)\_2(L)\_2\_CH\_2\_Cl\_2]} \]

\[ \text{C} = 42H_{40}N_{12}B_{2}Cl_{4}FeS_{6}, \text{calcd} = 1.467 \text{ g cm}^{-3}, \mu = 0.744 \text{ mm}^{-1}, R_{\text{obs}} = 0.0356, R_{\text{w}} = 0.0547, wR_{\text{298}} = 0.1562 \text{ for 9208 reflections with } I > 2\sigma(I) \text{ and 613 parameters, } S = 1.058. \]

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