1-Loop Matching of gauge invariant dim-6 operators for $B$ decays

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Physics beyond the Standard Model, realized above the electroweak scale, can be incorporated in a model independent way in the Wilson coefficients of higher dimensional gauge invariant operators. In these proceedings we review the matching of the $SU(3)_C \times SU(2)_L \times U(1)_Y$ gauge invariant dimension-six operators on the effective Hamiltonian governing $b \to s$ and $b \to c$ transitions, including the leading 1-loop effects [1].

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1. Introduction

Despite numerous confirmations of its validity, the Standard Model (SM) of particle physics is thought to be only an effective theory valid up to a new physics scale $\Lambda$, where additional dynamical degrees of freedom enter. The SM effective theory (SMET) Lagrangian can be written in the following form \[ (1.1) \]:

$$ L_{SM} = L_{SM}^{(4)} + \frac{1}{\Lambda} C_{VV}^{(5)} Q_{VV} + \frac{1}{\Lambda^2} \sum_{k} C_{k}^{(6)} Q_{k}^{(6)} + \mathcal{O} \left( \frac{1}{\Lambda^3} \right). $$

In this approach, physics beyond the SM is encoded in a model independent way in the Wilson coefficients of the higher dimensional operators $Q_k$. For $B$ physics, the Wilson coefficients $C_{k}^{(6)}$, multiplying the dimension-six operators are relevant, while the dimension-five Weinberg operator only provides neutrino mass terms after electroweak (EW) symmetry breaking \[ 4 \].

In order to compare the prediction of a NP model to $B$ physics observables, the following steps have to be performed.

1. Running of the Wilson coefficients $C_{k}^{(6)}$ from the matching scale $\Lambda$ to the electroweak (EW) symmetry breaking scale $\mu_W$ \[ 5 \].
2. EW symmetry breaking is performed and the SMET Lagrangian is matched onto the effective Hamiltonian governing $B$ physics.\footnote{For the corresponding calculation in the lepton sector see Refs. \[ 6, 7 \].}
3. Renormalization group equations can be used to perform the evolution of the Wilson coefficients from the electroweak scale down to the $B$ scale $\mu_b$ (see for example \[ 8, 9 \]).

This procedure is depicted in Fig. 1. The requirement of gauge invariance reduces the number of operators in $B$ physics \[ 10 \] and correlates charged with neutral currents (see for example \[ 11 \]).

Above the EW symmetry breaking scale the gauge invariant dimension-six operators are given in the interaction basis, since the mass basis is not defined above $\mu_W$. After the EW symmetry breaking, the fermions are rotated into the mass eigenstates by diagonalizing their mass matrices, which affects the Wilson coefficients. All rotation matrices appearing in the operators can be absorbed by a redefinition of the Wilson coefficients, with the exception of the misalignment between the left-handed up-quark and down-quark rotations, i.e. the Cabibbo-Kobayashi-Maskawa matrix.

In Ref. \[ 1 \] we performed the matching of the SMEFT onto the effective Hamiltonian of $B$ physics by integrating out all heavy degrees of freedom (compared to $\mu_b$), i.e. the top quark, the $W$ and $Z$ bosons and the Higgs field. The full tree-level matching has been computed for $b \to s$ and $b \to c$ transitions. In addition, 1-loop contributions have been performed, which involve dimension-six operators that do not enter the matching at tree level.

2. 1-loop matching

Operators with a top quark do not contribute to the $b \to s$ and $b \to c$ transitions at the tree-level, since the top is not contained in the $B$ physics Hamiltonian. The 1-loop matching contributions due to dimension-six operators containing right-handed top quarks can be divided in the following six classes:
1. 4-fermion operators to 4-fermion operators ($\Delta B = \Delta S = 1$).

2. 4-fermion operators to 4-fermion operators ($\Delta B = \Delta S = 2$).

3. 4-fermion operators to $\mathcal{O}_7$ and $\mathcal{O}_8$.

4. Right-handed $Z$ couplings to $\mathcal{O}_9$, $\mathcal{O}_{10}$ and $\mathcal{O}_{3-6}$.

5. Right-handed $W$ couplings to $\mathcal{O}_7$ and $\mathcal{O}_8$.

6. Magnetic operators to $\mathcal{O}_7$, $\mathcal{O}_8$, $\mathcal{O}_9$, $\mathcal{O}_{10}$ and $\mathcal{O}_{3-6}^q$.

As an example we consider the dimension-six operator $Q_{\phi ud} = (\tilde{\phi}^\dagger iD_\mu \phi)(\bar{u}_i \gamma^\mu P_R d_j)$, which couples the $W$-boson to right-handed quarks. This anomalous $W - t - b$ coupling induces a non-zero contribution to the magnetic operators $\mathcal{O}_7$, $\mathcal{O}_8$. The magnetic operators which are contained in the $\Delta B = \Delta S = 1$ effective Hamiltonian read:

$$\mathcal{O}_7 = \frac{e}{16\pi^2} m_b (\bar{s} \sigma^{\mu\nu} P_R b) F_{\mu\nu}, \quad \mathcal{O}_8 = \frac{g_s}{16\pi^2} m_b (\bar{s} \sigma^{\mu\nu} T^A b) G_{\mu\nu}^A.$$  \hspace{1cm} (2.1)

The matching contributions from $Q_{\phi ud}$ are given by (in agreement with [12, 13]):

$$C_7 = \frac{m_t}{m_b} \frac{\gamma^2}{\Lambda^2} E_{\phi ud}(x_t) C^{33}_{\phi ud} V_{ts}, \quad C_8 = \frac{m_t}{m_b} \frac{\gamma^2}{\Lambda^2} E_{\phi ud}(x_t) C^{33}_{\phi ud} V_{ts},$$  \hspace{1cm} (2.2)

where the dimensionless $x_t = m_t^2/M_W^2$-functions are defined in Ref. [1].
3. Conclusions

We presented the complete tree-level matching coefficients for $b \to s$ and $b \to c$ transitions including lepton flavor violating operators. 27 out of the 59 gauge invariant dimension-six operators contribute to the tree-level matching. Another 14 operators enter the 1-loop matching. They involve 4-fermion operators, electromagnetic and chromomagnetic dipole operators as well as operators involving Higgs and quark fields. Once the running from the EW scale down to the $B$ meson scale will be performed, our results can be used to perform systematic tests on the sensitivity of $B$ physics observables on the dimension-six operators.

References


[7] Giovanni Marco Pruna and Adrian Signer. The $\mu \to e\gamma$ decay in a systematic effective field theory approach with dimension 6 operators. *JHEP*, 10:014, 2014.


